

1. Assume that $u_0 = u_0(x_1, x_2)$ is a smooth function on \mathbf{R}^2 which is 1-periodic in each direction, i. e. $u(x_1 + 1, x_2) = u(x_1, x_2) = u(x_1, x_2 + 1)$. Let us consider the problem

$$u_t = \frac{\partial^2 u}{\partial x_1^2} + 3 \frac{\partial^2 u}{\partial x_2^2} + \frac{\partial u}{\partial x_1}, \quad x \in \mathbf{R}^2, t \geq 0 \quad (1)$$

$$u|_{t=0} = u_0. \quad (2)$$

If we express the solution u as

$$u(x, t) = \sum_{k \in \mathbf{Z}^2} c(k, t) e^{2\pi i(k_1 x_1 + k_2 x_2)}, \quad (3)$$

what are the formulae for $c(k, t)$ (assuming we know the Fourier coefficients of u_0)?

Solution: Differentiating (3), we see that the Fourier coefficients of the right-hand side of (1) is given by $A(k)c(k, t)$ with $A(k) = -(2\pi)^2 k_1^2 - 3(2\pi)^2 k_2^2 + 2\pi i k_1$. Hence $c(k, t) = e^{tA(k)} c(k, 0)$, where $c(k, 0) = c_0(k)$, the Fourier coefficients of u_0 .

2. If you have a computer program which can calculate the Fourier coefficient of a functions and also sum a given Fourier series on the square $[0, 1] \times [0, 1]$, how would you use it to solve the problem

$$-\Delta u = f \quad \text{in } \Omega \quad (4)$$

$$u|_{\partial\Omega} = 0 \quad \text{in } \Omega \quad (5)$$

when Ω is the rectangle $\{(x_1, x_2), 0 < x_1 < 2, 0 < x_2 < 3\}$?

Solution: We first need to adapt the problem to the setting of periodic functions. For that we extend f to a function $f_{\text{per}}: \mathbf{R}^2 \rightarrow \mathbf{R}$, which is odd and periodic with period 4 in x_1 , and odd and periodic with period 6 in x_2 , in a similar way as in Problem 2 of Homework 2. In particular, when $2 < x_2 < 4$ and $0 < x_1 < 3$, we set $f_{\text{per}}(x_1, x_2) = -f(4 - x_1, x_2)$; when $0 < x_1 < 2$ and $3 < x_2 < 6$, we set $f_{\text{per}}(x_1, x_2) = -f(x_1, 6 - x_2)$; and when $2 < x_1 < 4$ and $3 < x_2 < 6$ we set $f_{\text{per}}(x_1, x_2) = f(4 - x_1, 6 - x_2)$. It is now enough to find a function u_{per} which is odd and periodic with period 4 in x_1 , and odd and periodic with period 6 in x_2 such that $-\Delta u_{\text{per}} = f_{\text{per}}$. Such a function will solve (4), and the boundary condition $u|_{\partial\Omega} = 0$ will be satisfied due to the symmetries. As our program can only work with periodic functions of period 1 in each direction, we need to change variables to transfer the equation to $[0, 1] \times [0, 1]$. For this we set $F(y_1, y_2) = f_{\text{per}}(4y_1, 6y_2)$ and $U(y_1, y_2) = u_{\text{per}}(4y_1, 6y_2)$. In the y variables, the equation becomes $\frac{\partial^2 U}{16\partial y_1^2} + \frac{\partial^2 U}{36\partial y_2^2} = F$.

We use the program to calculate the Fourier coefficients $\hat{F}(k)$ of F . The Fourier coefficients of U will be given by $\hat{U}(k) = \frac{\hat{F}(k)}{\frac{4\pi^2 k_1^2}{16} + \frac{4\pi^2 k_2^2}{36}}$. We then use

the program to sum the Fourier series for U to obtain the function $U(y_1, y_2)$. The solution u is then given by $u(x_1, x_2) = U(\frac{x_1}{4}, \frac{x_2}{6})$.

3. Let $\Omega \subset \mathbf{R}^3$ be a bounded smooth domain. Assume that $\phi_1, \phi_2, \phi_3 \dots$ are the eigenfunctions of the Laplacian $-\Delta$ in Ω with the zero boundary condition and eigenvalues $\lambda_1 < \lambda_2 \leq \lambda_3 \leq \dots$. (In other words, $-\Delta\phi_j = \lambda_j\phi_j$ and $\phi_j|_{\partial\Omega} = 0$.) Find the solution of the initial-value problem for the Schrödinger equation

$$iu_t + \Delta u - u = 0 \quad \text{in } \Omega \times (0, \infty), \quad (6)$$

$$u|_{\partial\Omega} = 0, \quad (7)$$

$$u(x, 0) = u_0(x) \quad \text{in } \Omega, \quad (8)$$

where $i = \sqrt{-1}$ and $u_0(x) = \sum_{k=1}^n c_k \phi_k(x)$.

Solution: We seek the solution in the form $u(x, t) = \sum_k c(k, t) \phi_k(x)$. Substituting this expression into the PDE, we obtain $ic(k, t) - (\lambda_k + 1)c(k, t) = 0$ for each k . Solving the ode for each k , we obtain $u(x, t) = \sum_{k=1}^n e^{-i(\lambda_k + 1)t} c_k \phi_k(x)$.

4. Let $\Omega \subset \mathbf{R}^3$ be a smooth bounded domain and let g be a smooth function on $\partial\Omega$. What will be the PDE and the boundary conditions corresponding to the minimization of the functional

$$J(u) = \int_{\Omega} \frac{1}{2} (|\nabla u|^2 + u^2) dx - \int_{\partial\Omega} gu dx \quad (9)$$

over all smooth functions on Ω ?

Solution: If the minimum is attained at u , then the derivative of the function $\varepsilon \rightarrow J(u + \varepsilon\varphi)$ at $\varepsilon = 0$ must vanish for each smooth function φ on Ω . The derivative is given by $\int_{\Omega} (\nabla u \nabla \varphi + u\varphi) dx - \int_{\partial\Omega} g\varphi dx$. Integrating by parts in the first integral, we see that the last expression can also be written as $\int_{\Omega} (-\Delta u + u)\varphi dx + \int_{\partial\Omega} \left(\frac{\partial u}{\partial n} - g\right)\varphi dx$. This can only vanish for all smooth φ if $-\Delta u + u = 0$ in Ω and $\frac{\partial u}{\partial n} = g$ at $\partial\Omega$.

5. Find the solution $u(x, t)$ of the equation

$$u_t + \frac{\partial u}{\partial x_1} - 2 \frac{\partial u}{\partial x_2} + 5u = 0 \quad \text{in } \mathbf{R}^3 \times (-\infty, \infty) \quad (10)$$

with $u(x, 0) = u_0(x)$.

Solution: The solution of the equation $u_t + \frac{\partial u}{\partial x_1} - 2 \frac{\partial u}{\partial x_2} = 0$ corresponds to the transport of the original function with speed $(1, -2, 0)$, and hence is given by $u_0(x_1 - t, x_2 + 2t, x_3)$. The solution of $u_t + 5u = 0$ corresponds to $u_0(x)e^{-5t}$. The two "processes" are going on at the same time, but they commute. Hence the solution of (10) is given by $u(x, t) = u_0(x_1 - t, x_2 + 2t, x_3)e^{-5t}$.

6. Let $\Omega \subset \mathbf{R}^3$ be the complement of the ball of radius 3, i. e. $\Omega = \{x \in \mathbf{R}^3, |x| > 3\}$. Find the Green's function of the domain Ω for the equation $\Delta u = f$ with the boundary conditions $u|_{\partial\Omega} = 0$ and $u(x) \rightarrow 0$ as $|x| \rightarrow \infty$.

Solution: The Green's function of the ball of radius R is given by $G(x, y) = -\frac{1}{4\pi|x-y|} + \frac{R}{4\pi|y||x-y^*|}$, with $y^* = \frac{R^2}{|y|^2}y$, and we can think of it as the field of the unit charge at y minus the field of a "fictitious" charge $R/|y|$ at y^* . For the ball we think of y being inside the ball, i. e. $|y| < R$. When the domain is the outside of the ball, we can just exchange the role of y and y^* and normalize the charge outside of the ball to 1. The resulting formula is again $G(x, y) = -\frac{1}{4\pi|x-y|} + \frac{R}{4\pi|y||x-y^*|}$, with $y^* = \frac{R^2}{|y|^2}y$, except that this time y is outside of the ball.