## ENTROPY BOUNDS AND ISOPERIMETRY

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## Preface

During the second half of the XXth century one could observe an interesting development in the infinite dimensional analysis slowly moving it from the "experimental science" based on interesting examples to a brink of becoming something more than that. In particular during its last decade we have experienced a considerable progress in the domain of coercive inequalities and isoperimetry; as a starting point to explore the related literature we suggest [A-B-C-F-G-M-R-S], [G-Z] and [L5], as they contain an extensive overview and a comprehensive list of references. For the purposes of this work we would like to recall two important inequalities: the so called Logarithmic Sobolev Inequality of $[\mathbf{G} \mathbf{1}]$ and Isoperimetric Functional Inequality introduced in [B2]. The first one in plain language is simply an estimate of the relative entropy in terms of the Fisher entropy. Alternatively, from the perspective of application to control the ergodicity of Markov semigroup generated by the Dirichlet operator, we could think of the relative density as given by a square of a function divided by a normalisation constant and in this setting the upper bound in question is given by the Dirichlet form as follows

$$
\begin{equation*}
\operatorname{Ent}\left(|f|^{2}\right) \leq c \int \sum_{i}\left|\nabla_{i} f\right|^{2} d \mu \tag{2}
\end{equation*}
$$

The Isoperimetric Functional Inequality is a bound on the value of an isoperimetric function $\mathcal{I}$ computed at the point equal to the expectation of a positive bounded by 1 function $f$ by an expectation of the length of a vector which first component equals to the composition of $\mathcal{I}$ with $f$ and the other components are given by the components of the gradient of $f$ scaled by a constant

$$
\begin{equation*}
\mathcal{I}(\mu f) \leq \int \sqrt{\mathcal{I}(f)^{2}+c \sum_{i}\left|\nabla_{i} f\right|^{2}} d \mu \tag{2}
\end{equation*}
$$

By an appropriate approximation of a characteristic function the right hand side converges to the surface measure of the set whereas the limit of the left hand side is simply equal to the value of the isoperimetric function at the volume of the set in question; this provides justification for the name of the inequality. The first key point is that both of these inequalities have tensorisation property, (i.e. if they hold in two measure spaces, then the product of the measures also satisfies them). The second important common point is that both of them involve the same $l_{2}$-norm on the tangent space. Naturally given a variety of probability measures one may like to ask for an optimal information; that is for inequalities reflecting possibly precisely properties of a given measure. It is certainly well known that in finite dimensions the isoperimetric functions may be different for different families of the measures. In our work we would like to emphasise the role of the metric on the
tangent space showing in particular that nonequivalent in infinite dimensions $l_{q^{-}}$ metrics are associated to qualitatively different families of measures.

In brief the results and organisation of the paper is as follows. The first part is devoted to the study of the following relative entropy bounds

$$
\begin{equation*}
\operatorname{Ent}\left(|f|^{q}\right) \leq c \int \sum_{i}\left|\nabla_{i} f\right|^{q} d \mu \tag{q}
\end{equation*}
$$

with $q \in(1,2)$. In Section 1, after a brief introduction, we indicate that $\mathrm{LS}_{q}$ is intimately related to a faster decay of the tails of the corresponding distributions and on an abstract level via Herbst arguments one can obtain a suitable exponential bounds. Section 2 describes the $q$-Poincaré inequality- the first abstract consequence of our coercive inequality- as well as a certain improvement on the lower bound of the spectral gap in finite dimension. In Section 3 we show that similarly as in $[\mathbf{B}-\mathbf{G 1}]$ for $\mathrm{LS}_{2}$ also when in the extended region of the parameter $q$ one can think of $\mathrm{LS}_{q}$ as a Poincaré inequality with respect to a suitable Orlicz space. This section contains also an analog of the so called Rothaus lemma (on improvements of an inequality involving additionally $L_{q}$ norm of the function on its right hand side provided $q$-Poincaré is known). Sections 4 and 5 specialise to the probability measures on the real line providing essentially optimal bounds for the constants in the $q$-Poincaré and $q$-Log Sobolev inequality. The ideas involved here are based on utillisation of Hardy inequalities and the Orlicz norms. These results naturally extend the ones of $[\mathbf{B}-\mathbf{G 1}]$. An easy application of these results shows that for example measures with density $\left(\exp -x^{p}\right) / Z$, with $p \in(2, \infty)$ being the Hölder dual to $q$, satisfy $\mathrm{LS}_{q}$. Hence by the product property we recover the similar result of [B-L] which proof was essentially based on the Prekopa-Leindler inequality involving strong concavity properties of the log of the density. Additionally, in Section 5, we provide a class of examples of measures on the real line with smooth densities for which $\mathrm{LS}_{q}$ is true although the conditions of $[\mathbf{B}-\mathbf{L}]$ are infinitely violated; this class of examples goes in the spirit of $[\mathbf{G}-\mathbf{R}]$, ([A-B-C-F-G-M-R-S] $]$. We come back to exponential bounds in Section 6 providing its more optimal version, generalising the one of [B-G1] and allowing unbounded gradients of the random variables. In this section we indicate also that such bounds together with the Rothaus lemma can be used to show $\mathrm{LS}_{q}$ in finite dimensions for measures with highly singular densities. Section 7 contains a proof of the $\mathrm{LS}_{q}$ for Gibbs measures with superGaussian tails, (that is such that distribution of the natural coordinate functions decay faster than Gaussian's), associated to a local specification (i.e. an a'priori given family of regular conditional expectations with respect to complements of some finite dimensional sigma algebras) which satisfies a strong mixing condition. This result is obtained by a suitable adaptation of the technology developed in the past for $\mathrm{LS}_{2}$ and nicely confirms the robustness of the original ideas. Naturally one can be interested what, if any, consequences one can derive from $\mathrm{LS}_{q}$ for some family of Markov semigroups. Some preliminary results on this subject the reader can find in Section 8. In particular, confining to arbitrary but finite dimensions, we show that the semigroup generated by the ordinary Dirichlet operator is ultracontractive. Additionally we present there also some ergodicity estimates for the semigroups generated by the nonlinear Laplacians, (in particular obtaining an algebraic entropy decay which in the limit $q \rightarrow 2$ recovers the well known exponential
decay of entropy associated with the $\mathrm{LS}_{2}$ ). Sections 9-13 concern the high dimensional asymptotics of the isoperimetric problem for convex bodies with the uniform measure and its dependence on the underlying metric. The key result provides a lower bound for so called Sobolev constant in dependence on the intrinsic characteristic of the convex body involving inner metric and the dimension of the space. The proof utilises a great localisation technique of $[\mathbf{K}-\mathbf{L}-\mathbf{S}]$ which for convex bodies allows to reduce the problem to a suitable one dimensional problem. As a corollary of the estimates, by employing arguments based on the coarea formula, in Section 14 one obtains the Classical Sobolev inequality. This precise asymptotics together with some appropriate limiting arguments allow us to recover in Section 15 the $\mathrm{LS}_{q}$ inequality. Thus in the sense we have proved the equivalence of the $\mathrm{LS}_{q}$ and the infinite dimensional asymptotics of the isoperimetric problem. Finally in Section 16 we consider the following functional isoperimetric inequality $\mathrm{IFI}_{q}$ introduced in [Z1]

$$
\begin{equation*}
\mathcal{U}(\mu f) \leq \int \sqrt[q]{\mathcal{U}(f)^{q}+c \sum_{i}\left|\nabla_{i} f\right|^{q}} d \mu \tag{2}
\end{equation*}
$$

As shown in $[\mathbf{Z 1}]$ it has a product property. Moreover, by a suitable choice of the isoperimetric function $\mathcal{U}$, in natural way it reflects known isoperimetric properties of the finite dimensional restrictions of the measures with super Gaussian tails. In this section we show the implication $\mathrm{IFI}_{q} \Longrightarrow \mathrm{LS}_{q}$ generalizing the known case $q=2$. (We remark that in the case $q=2$ the converse inequality was also proved in [Fo]; though at this point of time the semigroup technique does not seem to admit the necessary generalisation.)

These are the results, but at this point it may be proper to indicate where all that could be going how that could be useful. Certainly further deeper understanding the relation between the metric properties and isoperitmetry as well as coercive and functional isoperimetric inequalities would be useful and interesting. While in the above we have mentioned only the super Gaussian tails, we should mention that there exist also some results for the measures with sub Gaussian tails (that is with decay between exponential and the Gaussian). In particular an interesting inequality which has a product property was proved in $[\mathbf{L}-\mathbf{O}]$ providing a bound for a difference between the square of $L_{2}$ and the square of $L_{r}$ norm (for any $r \in[1,2)$ ) in terms of the Dirichlet form. As showed by the authors such bound allows to recover the correct exponential bounds; (for an earlier work involving generalized Poincaré inequalities in some other context see $[\mathbf{B e}]$ and for the product property in a more general setting see $[\mathbf{L 3}]$, $[\mathbf{L i e}])$. Later, it was demonstrated in the epi$\log$ to $[\mathbf{G}-\mathbf{Z}]$ that their inequality extend to a large class of Gibbs measures with sub Gaussian tails. Alternative results involving the same class of measures and bounds of entropy by higher order differential expressions were proven in $[\mathbf{Z 2}]$ following natural link with classical Sobolev inequalities for probability measures of [Ros]. Understanding properly all interconnections as well as the isoperimetry in this case remains still a challegenge. One should hope that plentiful of interesting things may follow this research in terms of understanding the new classes of infinite dimensional equations involving nonlinear as well as differential operators of higher order.

## CHAPTER 1

## Introduction and notations

Let $(\Omega, \mu)$ denote a probability space, and assume there is an operator $\Gamma$ defined on some algebra $\mathcal{A}$ of measurable functions on $\Omega$ with the following properties:

1) for any $f \in \mathcal{A}, \Gamma(f)$ is a non-negative measurable function on $\Omega$;
2) for all $f \in \mathcal{A}$ and smooth functions $u$ on $\mathbf{R}, u(f) \in \mathcal{A}$ and $\Gamma(u(f))=$ $\left|u^{\prime}(f)\right| \Gamma(f)$.
Introduce the entropy functional:

$$
\operatorname{Ent}(g)=\int g \log g d \mu-\int g d \mu \log \int g d \mu
$$

It is well defined for all measurable $g \geq 0$ and then it is finite if and only if so is the integral $\int g \log (1+g) d \mu$.

One says that $(\Omega, \mu, \Gamma)$ satisfies a logarithmic Sobolev inequality, for short LSI with constant $c \geq 0$ if, for all $f \in \mathcal{A}$,

$$
\begin{equation*}
\operatorname{Ent}\left(|f|^{2}\right) \leq c \int \Gamma(f)^{2} d \mu \tag{1.0.1}
\end{equation*}
$$

The definition is similar to the one considered in [A-M-S].
For example, when $\Omega$ is a metric space with metric $d$, the "modulus of the gradient" comes naturally via the identity

$$
\Gamma(f)(x)=|\nabla f(x)|=\limsup _{d(x, y) \rightarrow 0^{+}} \frac{|f(x)-f(y)|}{d(x, y)}
$$

with the convention that $\Gamma(f)(x)=0$ at isolated points $x$ in $\Omega$. In this case, we may define $\Gamma$ on the class $\mathcal{A}$ of all Lipschitz functions $f$, i.e., such that $\|f\|_{\text {Lip }}<\infty$, or for a larger class of all locally Lipschitz functions (i.e., Lipschitz in a neighbourhood of any point).

Moreover, in case of the Euclidean space $\Omega=\mathbf{R}^{n}$ with the usual Euclidean metric $d(x, y)=|x-y|$, we clearly have $\Gamma(f)(x)=|\nabla f(x)|$ at each point $x \in \mathbf{R}^{n}$ where $f$ is differentiable and has gradient $\nabla f(x)$. Since locally Lipschitz functions are differentiable almost everywhere (with respect to Lebesgue measure), we then arrive in (1.1) at the usual definition of logarithmic Sobolev inequalities.

More generally, given a number $p \in(1,+\infty]$, we may equip $\mathbf{R}^{n}$ with $\ell^{p}$-metric $d(x, y)=\|x-y\|_{p}=\left(\sum_{i=1}^{n}\left|x_{i}-y_{i}\right|^{p}\right)^{1 / p}$, and then we obtain another important modulus of gradient,

$$
\Gamma(f)(x)=|\nabla f(x)|_{q}=\left(\sum_{i=1}^{n}\left|\frac{\partial}{\partial x_{i}} f(x)\right|^{q}\right)^{1 / q},
$$

where $q=\frac{p}{p-1}$ is conjugate to $p$.

Actually, already the one-dimensional case contains a number of open problems. One of questions of interest is to determine whether or not, a given probability measure $\mu$ satisfies LSI with some finite $c$. In the case of the real line $\Omega=\mathbf{R}$, the answer is known [B-G1], and we recall the obtained characterization.

Let $F(x)=\mu((-\infty, x]), x \in \mathbf{R}$, denote the distribution function of $\mu$, and let $p$ be the density of the absolutely continuous part of $\mu$ with respect to Lebesgue measure. Let $m$ denote a median of $\mu$. Set

$$
\begin{gathered}
D_{0}=\sup _{x<m}\left(F(x) \log \frac{1}{F(x)}\right) \int_{x}^{m} \frac{1}{p(t)} d t \\
D_{1}=\sup _{x>m}\left((1-F(x)) \log \frac{1}{1-F(x)}\right) \int_{m}^{x} \frac{1}{p(t)} d t
\end{gathered}
$$

defining $D_{0}$ and $D_{1}$ to be zero in case $\mu((-\infty, m))=0$ or $\mu((m,+\infty))=0$, respectively.

Theorem 1.1. The measure $\mu$ satisfies the log-Sobolev inequality (1.1) with some constant $c$ if and only if $D_{0}+D_{1}<+\infty$. In this case, the optimal value of $c$ satisfies

$$
K_{0}\left(D_{0}+D_{1}\right) \leq c \leq K_{1}\left(D_{0}+D_{1}\right),
$$

where $K_{0}$ and $K_{1}$ are certain absolute positive constants.
Due to the tensorization property, the log-Sobolev inequality (1.1) can be extended to product spaces without any loss in constant. In high-dimensional spaces it can be used to recover various concentration phenomena, dependently on which classes $\mathcal{A}$ and gradients $\Gamma$ are involved in (1.1). On the other hand, it is a powerful tool in the study of hypercontractivity of semi-groups (cf. [G1], [G2]). However, more specific and delicate properties of certain measures suggest to consider somewhat different forms of log-Sobolev inequalities. In particular, inequalities

$$
\begin{equation*}
\operatorname{Ent}\left(|f|^{q}\right) \leq c \int \Gamma(f)^{q} d \mu \tag{1.0.2}
\end{equation*}
$$

with (necessarily) some $q \in[1,2]$, may serve as a certain sharpening of (1.1), describing more accurately the behaviour of distributions of Lipschitz functions. For example, as shown by O. Rothaus $[\mathrm{R}]$, the particular case $q=1$ reduces (1.2) on a Riemannian manifold to a certain isoperimetric inequality.

Definition 1.2. For short, in what follows (1.2) is called $\mathrm{LS}_{q}$-inequality and if a measure $\mu$ satisfies it with a given coefficient $c \in(0, \infty)$ we denote that by $\mu \in \mathrm{LS}_{q}(c)$.

Let us recall an argument which relates $\mathrm{LS}_{q}$-inequality to deviations of "Lipschitz" functions from their means and therefore to the concentration of measure. The idea goes back to an unpublished letter by I. Herbst (of mid 70s), and only in late 80s it was reanimated and developed in the works [A-M-S] and [L2,3]). Let $1<q \leq 2$ and assume $g$ is a bounded function in $\mathcal{A}$ such that $\int g d \mu=0$ and $\Gamma(g) \leq 1 \mu$-a.e. We may apply (1.2) to the function $f=e^{\lambda g / q}, \lambda \geq 0$ : due to the axiom 2), it belongs to $\mathcal{A}$ and, moreover, $\Gamma(f)^{q} \leq\left(\frac{\lambda}{q}\right)^{q} e^{\lambda g}$. Write the Laplace transform on $g$ as $\int e^{\lambda g} d \mu=e^{\lambda v(\lambda)}$ for some function $v$ which is clearly smooth in $\lambda>0$ and satisfies $v(0+)=0$ (since $g$ has $\mu$-mean zero). Thus, the right hand side of (1.2)
does not exceed $c\left(\frac{\lambda}{q}\right)^{q} e^{\lambda v(\lambda)}$. On the other hand, $\operatorname{Ent}\left(e^{\lambda g}\right)=\lambda^{2} v^{\prime}(\lambda) \int e^{\lambda g} d \mu$, so $\lambda^{2} v^{\prime}(\lambda) \leq c\left(\frac{\lambda}{q}\right)^{q}$ and

$$
v(\lambda)=\int_{0}^{\lambda} v^{\prime}(t) d t \leq \frac{c}{q^{q}} \int_{0}^{\lambda} t^{q-2} d t=\frac{c}{q^{q}(q-1)} \lambda^{q-1}
$$

Thus, we have obtained the following bound on the Laplace transform:

$$
\begin{equation*}
\int e^{\lambda g} d \mu \leq \exp \left\{\frac{c}{q^{q}(q-1)} \lambda^{q}\right\}, \quad \lambda>0 \tag{1.0.3}
\end{equation*}
$$

Now, given a number $h>0$, by Chebyshev's inequality and (1.3), we get

$$
\mu\{g \geq h\} \leq e^{-\lambda h} \int e^{\lambda g} d \mu \leq \exp \left\{\lambda h-\frac{c}{q^{q}(q-1)} \lambda^{q}\right\}, \quad \lambda>0 .
$$

Optimization over all $\lambda$ leads to the one-sided estimate $\mu\{g \geq h\} \leq \exp \left\{-\frac{(q-1)^{p}}{c^{p-1}} h^{p}\right\}$, where $p$ is conjugate to $q$. Applying it to $-g$, we obtain a similar estimate. Hence, we arrive at the following important conclusion:

Theorem 1.3. Under $\mathrm{LS}_{q}, 1<q \leq 2$, for every bounded function $g \in \mathcal{A}$ with mean $a=\int g d \mu$, such that $\Gamma(g) \leq 1 \mu$-a.e.,

$$
\begin{equation*}
\mu\{|g-a| \geq h\} \leq 2 \exp \left\{-\frac{(q-1)^{p}}{c^{p-1}} h^{p}\right\}, \quad h>0 \tag{1.0.4}
\end{equation*}
$$

where $p=\frac{q}{q-1}$.
The main feature of the inequality (1.4) is that, in case $q<2$, it indicates a much stronger decay of the tails of $g$ in comparison with the classical "Gaussian" case $q=2$. Of course, in the reasonable situation such as an abstract probability metric space, the assumption on boundedness of $g$ can be removed (via usual approximations). Thus, (1.4) remains to hold for unbounded Lipschitz $g$, as well. Moreover, application of (1.4) to functions of the form $\operatorname{dist}(A, x)$ leads to certain concentration inequalities often complementing or illustrating the so-called concentration of measure phenomenon (cf. e.g. [L4,5]). As for (1.3), one may wonder whether or not it is possible to replace the Lipschitz condition by a more flexible property of the modulus of gradient; we will concern this question in section 6 .
$\mathrm{LS}_{q}$ inequalities (1.2), together with deviation inequality (1.4), are known to hold for probability measures on $\mathbf{R}^{n}$ with sufficiently "convex" densities, cf. [B-L]. Here we will mainly specialize on the question of characterization of $\mu$ on the real line $\mathbf{R}$ which thus satisfy the inequality

$$
\operatorname{Ent}\left(|f|^{q}\right) \leq c \int_{-\infty}^{+\infty}\left|f^{\prime}(x)\right|^{q} d \mu(x)
$$

in the class of all smooth functions $f$, or equivalently, locally Lipschitz $f$, up to some constant $c$.

The proof of Theorem 1.1 given in [B-G1] uses a particular $L_{2}$-case of result of M. Artola, G. Talenti and G. Tomaselli on Hardy-type inequalities with weights on the real line (cf. [Mu], [M2]). Here we have to use it in the full volume involving $L_{q}$ spaces. A non-trivial point is that the general case requires a more careful analysis
of possible behaviour of the entropy functional Ent $\left(|f|^{q}\right)$ and of what is connected with it,

$$
\begin{equation*}
\mathcal{L}_{q}(f)=\sup _{a \in \mathbf{R}} \operatorname{Ent}\left(|f+a|^{q}\right) . \tag{1.0.5}
\end{equation*}
$$

Indeed, the log-Sobolev inequality (1.1) may formally be strengthened as

$$
\mathcal{L}_{q}(f) \leq c \int \Gamma(f)^{q} d \mu
$$

Note that, when $f$ is bounded, $\operatorname{Ent}\left(|f+a|^{q}\right)$ is of order $a^{q-2} \operatorname{Var}(f)$, for large values of $a$ (where $\operatorname{Var}(f)$ denotes variance of $f$ with respect to $\mu$ ). So, in case $q>2$, $\mathcal{L}_{q}(f)=+\infty$ and thus (1.2) can never hold. In the case $1 \leq q<2$, we have Ent $\left(|f+a|^{q}\right) \rightarrow 0$, as $|a| \rightarrow+\infty$, and so we cannot derive from (1.2) a Poincarétype inequality similarly to the case $q=2$. This is a typical example illustrating some features of these more general inequalities.

## CHAPTER 2

## Poincaré-type inequalities

Here, in the setting of a probability metric space $(\Omega, d, \mu)$, we deduce from (1.2) a corresponding Poincaré-type inequality:

Theorem 2.1. Under $\mathrm{LS}_{q}$-inequality with constant $c$, for any locally Lipschitz function $f$ on $\Omega$, we have the following Poincaré type inequality called later on $\mathrm{SG}_{q}$-inequality

$$
\begin{equation*}
\int\left|f-\int f d \mu\right|^{q} d \mu \leq \frac{4 c}{\log 2} \int|\nabla f|^{q} d \mu \tag{2.0.1}
\end{equation*}
$$

More precisely, as soon as $|\nabla f| \in L_{q}(\mu)$, we have $f \in L_{q}(\mu)$, and (2.1) holds true.
Recall that $|\nabla f(x)|=\lim \sup _{d(x, y) \rightarrow 0^{+}} \frac{|f(x)-f(y)|}{d(x, y)}$. In particular, this definition is applied on the real line, so $f$ does not need to be differentiable everywhere (correspondingly, we understand partial derivatives $\frac{\partial f(x, y)}{\partial y}$ for functions defined on the product space $\Omega \times \mathbf{R}$ ).

For the proof of (2.1), we need a general
Lemma 2.2. For any non-negative measurable function $g$ on $\Omega$,

$$
\begin{equation*}
\operatorname{Ent}(g) \geq-\log \mu\{g>0\} \int g d \mu \tag{2.0.2}
\end{equation*}
$$

Proof. The distribution of $g$ under $\mu$ represents a probability measure on the positive half-axis $[0,+\infty)$ and may be written as $(1-\alpha) \delta_{0}+\alpha \nu$ where $\alpha=\mu\{g>0\}$ and where $\nu$ is concentrated on $(0,+\infty)$. In terms of a random variable, say $\xi$, having the distribution $\nu$, we thus may write

$$
\begin{gathered}
\int g d \mu=\int_{0}^{+\infty} x d\left((1-\alpha) \delta_{0}+\alpha \nu\right)=\alpha \mathbf{E} \xi \\
\int g \log g d \mu=\int_{0}^{+\infty} x \log x d\left((1-\alpha) \delta_{0}+\alpha \nu\right)=\alpha \mathbf{E} \xi \log \xi
\end{gathered}
$$

so,

$$
\begin{aligned}
\operatorname{Ent}(g) & =\alpha \mathbf{E} \xi \log \xi-(\alpha \mathbf{E} \xi) \log (\alpha \mathbf{E} \xi) \\
& =\alpha \operatorname{Ent}(\xi)-\alpha \log \alpha \mathbf{E} \xi \geq-\log \alpha(\alpha \mathbf{E} \xi)
\end{aligned}
$$

Proof of Theorem 2.1. Let $f$ be a locally Lipschitz function from $L_{q}(\Omega, \mu)$. Without loss of generality, assume that zero is a median of $f$, i.e., $\mu\{f>0\} \leq \frac{1}{2}$ and $\mu\{f<0\} \leq \frac{1}{2}$. First assume that $\mu\{f=0\}=0$.

Set $f^{+}=\max (f, 0), f^{-}=\max (-f, 0)$. These functions are also locally Lipschitz, and we can apply to them the log-Sobolev inequality (1.2). The moduli
of gradients of $f^{+}$and $f^{-}$are respectively vanishing on the open sets $\{f<0\}$, $\{f>0\}$ and coincide with $|\nabla f|$ on $\{f>0\},\{f<0\}$. Since $\mu\{f=0\}=0$, the further application of (2.2) to $g=\left(f^{+}\right)^{q}$ and $g=\left(f^{-}\right)^{q}$ gives respectively

$$
\begin{aligned}
& c \int_{\{f>0\}}|\nabla f|^{q} d \mu \geq \operatorname{Ent}\left(\left(f^{+}\right)^{q}\right) \geq \log 2 \int_{\{f>0\}}\left(f^{+}\right)^{q} d \mu \\
& c \int_{\{f<0\}}|\nabla f|^{q} d \mu \geq \operatorname{Ent}\left(\left(f^{-}\right)^{q}\right) \geq \log 2 \int_{\{f<0\}}\left(f^{-}\right)^{q} d \mu
\end{aligned}
$$

Summing these yields

$$
c \int|\nabla f|^{q} d \mu \geq \log 2 \int|f|^{q} d \mu=\log 2\|f\|_{q}^{q}
$$

But $\|f-\mathbf{E} f\|_{q} \leq\|f\|_{q}+\|f\|_{1} \leq 2\|f\|_{q}$, so $\|f-\mathbf{E} f\|_{q}^{q} \leq 2^{q}\|f\|_{q}^{q} \leq 4\|f\|_{q}^{q}$, and this is what we need.

Note that, with a similar argument, we could proceed without the assumption $\mu\{f=0\}=0$ and gain the double constant in (2.1) by using the estimates $\left|\nabla f^{+}\right| \leq$ $|\nabla f|,\left|\nabla f^{-}\right| \leq|\nabla f|$ which hold true on the whole space $\Omega$. There is another simple way to avoid this assumption which however saves the same constant.

Namely, let $\xi$ be a random variable with mean zero and whose distribution, say, $\nu$ satisfies on the real line $\operatorname{LS}_{q}(1.2)$ with constant $c_{q}$. For example, the measure $\nu$ may have density of the form const $e^{-|x|^{p}}$ where $p$ is conjugate to $q$. Hence, the distribution $\nu_{\varepsilon}$ of the random variable $\varepsilon \xi$ satisfies $\mathrm{LS}_{q}$ with constant $c_{q} \varepsilon^{q}$. By the tensorization property of entropy, the log-Sobolev inequality (1.2) implies

$$
\begin{aligned}
\operatorname{Ent}_{\mu \otimes \nu_{\varepsilon}}\left(|g(x, y)|^{q}\right) \leq & c \int_{R} \int_{\Omega}\left|\nabla_{x} g(x, y)\right|^{q} d \mu(x) d \nu_{\varepsilon}(y)+ \\
& c_{q} \varepsilon^{q} \int_{R} \int_{\Omega}\left|\partial_{y} g(x, y)\right|^{q} d \mu(x) d \nu_{\varepsilon}(y)
\end{aligned}
$$

where $g$ is now locally Lipschitz on the product space $\Omega \times \mathbf{R}$ equipped with the product measure $\mu \otimes \nu_{\varepsilon}$. Thus, the product probability space ( $\Omega \times \mathbf{R}, \mu \otimes \nu_{\varepsilon}$ ) satisfies $\mathrm{LS}_{q}$ with respect to the (non-canonical) gradient functional

$$
\Gamma_{\varepsilon} g(x, y)=\left(c\left|\nabla_{x} g(x, y)\right|^{q}+c_{q} \varepsilon^{q}\left|\partial_{y} g(x, y)\right|^{q}\right)^{1 / q}
$$

Now, given a locally Lipschitz function $f$ from $L_{q}(\Omega, \mu)$, the function $g(x, y)=$ $f(x)+y$ is also locally Lipschitz, belongs to $L_{q}\left(\Omega \times \mathbf{R}, \mu \otimes \nu_{\varepsilon}\right)$, and is equidistributed, under $\mu \otimes \nu_{\varepsilon}$, with $f+\varepsilon \xi$ where $\xi$ is viewed as being independent of $f$. In particular, $g$ has a smooth positive density on the whole line, so the median $m$ of $g$ is uniquely determined and, moreover, $\mu \otimes \nu_{\varepsilon}\{g=m\}=0$. The functional $\Gamma_{\varepsilon}$ enjoys many properties of the modulus of gradient, at least in part concerning the function $g$. Namely, $\Gamma_{\varepsilon} g^{+}=\Gamma_{\varepsilon} g$ on $\{g>0\}, \Gamma_{\varepsilon} g^{+}=0$ on $\{g<0\}$, and similarly for $g^{-}$. Therefore, one can repeat the previous argument and thus arrive, according to (2.1) and the first step, at

$$
\begin{aligned}
\mathbf{E} \int\left|(f(x)+\varepsilon \xi)-\int f d \mu\right|^{q} d \mu(x) & \leq \frac{4}{\log 2} \int\left(\Gamma_{\varepsilon} g\right)^{q} d \mu \otimes \nu_{\varepsilon} \\
& =\frac{4}{\log 2} \int_{R} \int_{\Omega}\left(c\left|\nabla_{x} f(x)\right|^{q}+c_{q} \varepsilon^{q}\right) d \mu(x) d \nu_{\varepsilon}(y) \\
& =\frac{4}{\log 2}\left(c \int_{\Omega}|\nabla f|^{q} d \mu+c_{q} \varepsilon^{q}\right)
\end{aligned}
$$

It remains to let $\varepsilon \rightarrow 0$. Proposition 2.1 is now proved.
It will also be useful to note that any Poincaré-type inequality in $L^{q}$ space with $q \in[1,2)$,

$$
\begin{equation*}
\int\left|f-\int f d \mu\right|^{q} d \mu \leq C_{q} \int|\nabla f|^{q} d \mu \tag{2.0.3}
\end{equation*}
$$

holding true in the class of all locally Lipschitz functions $f$ on $\Omega$, is stronger than the usual spectral gap,

$$
\begin{equation*}
\int\left|f-\int f d \mu\right|^{2} d \mu \leq C_{2} \int|\nabla f|^{2} d \mu \tag{2.0.4}
\end{equation*}
$$

Namely, we have:
Proposition 2.3. For any $q \in[1,2]$, the optimal constants in (2.3) - (2.4) are connected by

$$
C_{2}^{1 / 2} \leq 6 C_{q}^{1 / q}
$$

Remark 2.4. Since we are still dealing here with the abstract probability metric space $(\Omega, d, \mu)$, the modulus of gradient has unique sense and should not be mixed in case, for example, $\Omega=\mathbf{R}^{n}$ or $\Omega=\mathbf{R}^{\infty}$ with other natural gradients such as $\|\nabla f\|_{q}$ in (2.3).

Proof of Proposition 2.3. Let $1 \leq q<2$. First we transform (2.3) to an analogue with the median in the place of the mean. Let $f$ be a locally Lipschitz function in $L^{q}(\mu)$ such that its $\mu$-mean is zero. Thus, by (2.3),

$$
\begin{equation*}
\int|f|^{q} d \mu \leq C_{q} \int|\nabla f|^{q} d \mu \tag{2.0.5}
\end{equation*}
$$

Denote by $m=m(f)$ its median so that $\mu\{f \leq m\} \geq \frac{1}{2}$ and $\mu\{f \geq m\} \geq \frac{1}{2}$. Assume for definiteness that $m>0$. Since by Chebyshev's inequality $\mu\{f \geq m\} \leq$ $\|f\|_{q}^{q} / m^{q}$, we may conclude that $m^{q} \leq 2\|f\|_{q}^{q}$. Thus, in general,

$$
|m| \leq 2^{1 / q}\|f\|_{q}
$$

Hence, $\|f-m\|_{q} \leq\left(1+2^{1 / q}\right)\|f\|_{q}$, and by (2.5),

$$
\begin{equation*}
\int|f-m(f)|^{q} d \mu \leq\left(1+2^{1 / q}\right)^{q} C_{q} \int|\nabla f|^{q} d \mu \tag{2.0.6}
\end{equation*}
$$

This inequality is translation invariant, so it holds true for all locally Lipschitz functions in $L^{q}(\mu)$. Now, take such a function and assume moreover that $f$ is square integrable, and that $m(f)=0$. Assume that $\mu\{f=0\}=0$ and consider the locally Lipschitz functions $f^{+}=\max (f, 0), f^{-}=\max (-f, 0)$. As in the proof of Proposition 2.1, let us note once more that the moduli of gradients of $f^{+}$and $f^{-}$are respectively vanishing on the open sets $\{f<0\},\{f>0\}$ and coincide with $|\nabla f|$ on $\{f>0\},\{f<0\}$. Since $m\left(f^{+}\right)=m\left(f^{-}\right)=0$, the application of (2.6) to $\left(f^{+}\right)^{2 / q}$ and $\left(f^{-}\right)^{2 / q}$ gives respectively

$$
\begin{aligned}
& \int_{\{f>0\}}|f|^{2} d \mu \leq\left(1+2^{1 / q}\right)^{q} C_{q}\left(\frac{2}{q}\right)^{q} \int_{\{f>0\}}|f|^{2-q}|\nabla f|^{q} d \mu \\
& \int_{\{f<0\}}|f|^{2} d \mu \leq\left(1+2^{1 / q}\right)^{q} C_{q}\left(\frac{2}{q}\right)^{q} \int_{\{f<0\}}|f|^{2-q}|\nabla f|^{q} d \mu
\end{aligned}
$$

Summing these yields

$$
\begin{equation*}
\int|f|^{2} d \mu \leq C_{q}\left(\frac{2\left(1+2^{1 / q}\right)}{q}\right)^{q} \int|f|^{2-q}|\nabla f|^{q} d \mu \tag{2.0.7}
\end{equation*}
$$

Since the conjugate of $\frac{2}{q}$ is $\frac{2}{2-q}$, the latter integral can be estimated with the help of Hölder's inequality by $\left\||f|^{2-q}\right\|_{2 /(2-q)}\left\||\nabla f|^{q}\right\|_{2 / q}=\|f\|_{2}^{2-q}\|\nabla f\|_{2}^{q}$. Hence, (2.7) implies

$$
\int|f|^{2} d \mu \leq C_{q}^{2 / q}\left(\frac{2\left(1+2^{1 / q}\right)}{q}\right)^{2} \int|\nabla f|^{2} d \mu
$$

Since $\int\left|f-\int f d \mu\right|^{2} d \mu \leq \int|f|^{2} d \mu$, we arrive at

$$
\int\left|f-\int f d \mu\right|^{2} d \mu \leq C_{q}^{2 / q}\left(\frac{2\left(1+2^{1 / q}\right)}{q}\right)^{2} \int|\nabla f|^{2} d \mu
$$

where the assumption that the function $f$ has median zero may already be omitted. The assumption that $\mu\{f=m(f)\}=0$ can be omitted by an argument used in the proof of Proposition 2.1. Thus, we arrived at the claim (2.4) with constant

$$
C_{2}^{1 / 2} \leq \frac{2\left(1+2^{1 / q}\right)}{q} C_{q}^{1 / q} \leq 6 C_{q}^{1 / q}
$$

Combining Propositions 2.1 and 2.3, we may conclude that $\mathrm{LS}_{q}$ with constant $c$ implies the usual Poincaré inequality with constant $C_{2} \leq 36\left(\frac{4 c}{\log 2}\right)^{2 / q}$. This spectral gap can actually be improved. Indeed, when assuming in the proof of Proposition 2.1 that $f$ has median zero, we derived a Poincaré-type inequality (2.1) with a better constant,

$$
\int|f|^{q} d \mu \leq \frac{c}{\log 2} \int|\nabla f|^{q} d \mu
$$

And under the same assumption, in the proof of Proposition 2.3 we deduced from the latter inequality the following one:

$$
\int|f|^{2} d \mu \leq\left(\frac{c}{\log 2}\right)^{2 / q}\left(\frac{2}{q}\right)^{2} \int|\nabla f|^{2} d \mu
$$

The constant is maximized at $q=1$, and this leads to
Corollary 2.5. Under $\mathrm{LS}_{q}$-inequality with constant $c$, for any locally Lipschitz function $f$ on $\Omega$,

$$
\int\left|f-\int f d \mu\right|^{2} d \mu \leq 9 c^{2 / q} \int|\nabla f|^{2} d \mu
$$

## CHAPTER 3

## Entropy and Orlicz spaces

In order to further explore $\mathrm{LS}_{q}$-inequalities, one has to bring the attention to some delicate properties, although yet general, of the entropy functional. In this section, we collect mainly those of them which connect this functional with norm in Orlicz space.

Given a Young function $N: \mathbf{R} \rightarrow[0,+\infty)$, i.e., an even, convex function with $N(0)=0, N(x)>0$ for $x>0$, the Orlicz space $L_{N}=L_{N}(\Omega, \mu)$ consists of all measurable functions $f$ such that

$$
\|f\|_{N}=\inf \left\{\lambda>0: \int N(f / \lambda) d \mu \leq 1\right\}<+\infty
$$

Any Young function $N$ strictly increases on $[0,+\infty)$ so an inverse $N^{-1}:[0,+\infty) \rightarrow$ $[0,+\infty)$ exists. When $N(x)=|x|^{q}(1 \leq q<+\infty), L_{N}$ is the usual Lebesgue space with norm $\|f\|_{q}$. We will consider the Orlicz norms for the two Young functions

$$
N_{q}(x)=|x|^{q} \log \left(1+|x|^{q}\right) \quad \text { and } \quad \Psi(x)=|x| \log (1+|x|) .
$$

Eventually, we will obtain here the following characterization.

Proposition 3.1. Let $1 \leq q \leq 2$. In the class of all locally Lipschitz integrable functions on a probability metric space $(\Omega, d, \mu)$, the $\mathrm{LS}_{q}$-inequality with constant $c$ is equivalent to

$$
\begin{equation*}
\left\|f-\int f d \mu\right\|_{N_{q}}^{q} \leq C \int|\nabla f|^{q} d \mu \tag{3.0.1}
\end{equation*}
$$

Moreover, the optimal constants there satisfy $\frac{1}{7} C \leq c \leq 16 C$.
Thus, up to a numerical constant, $\mathrm{LS}_{q}$ can be expressed as a certain Poincarétype inequality in the Orlicz space $L_{N_{q}}(\Omega, \mu)$.

We start with observations in the setting of a general probability space $(\Omega, \mu)$.
Lemma 3.2. For any function $f \in L_{N_{q}}(\mu), q \geq 1$,

$$
\|f\|_{q}^{q} \leq \frac{5}{4}\|f\|_{N_{q}}^{q} .
$$

Proof. The optimal constant is attained at $f=1$ and thus equals $\Psi^{-1}(1)<\frac{5}{4}$.
Lemma 3.3. For any function $f \in L_{N_{q}}(\mu), q \geq 1$,

$$
\|f\|_{N_{q}}^{q} \leq \operatorname{Ent}\left(|f|^{q}\right)+\|f\|_{q}^{q}
$$

Proof. We may assume $\|f\|_{q}=1$ so that $\operatorname{Ent}\left(|f|^{q}\right)=\int|f|^{q} \log |f|^{q} d \mu$. Using $x \log (1+x) \leq x \log x+1$, valid for all $x \geq 0$, we get

$$
\int N_{q}(f) d \mu \leq \operatorname{Ent}\left(|f|^{q}\right)+1 \equiv \alpha
$$

As follows from the definition, for all $t \in[0,1]$ and $x \in \mathbf{R}, N_{q}(t x) \leq t^{q} N_{q}(x)$. Hence, applying this to $t=\frac{1}{\alpha}$, we have

$$
\int N_{q}\left(\frac{f}{\alpha^{1 / q}}\right) d \mu \leq \frac{1}{\alpha} \int N_{q}(f) d \mu \leq 1
$$

which means that $\|f\|_{N_{q}} \leq \alpha^{1 / q}$.
Lemma 3.4. For any function $f \in L_{N_{q}}(\mu), q \geq 1$,

$$
\operatorname{Ent}\left(|f|^{q}\right) \leq \frac{e+1}{e}\|f\|_{N_{q}}^{q}
$$

Proof. We may now assume that $\|f\|_{N_{q}}=1$ so that $\int|f|^{q} \log \left(1+|f|^{q}\right) d \mu=1$. Put $g=|f|^{q}$. The function $-x \log x$ attains its maximum at $x=1 / e$ equal to $1 / e$. Therefore,

$$
\operatorname{Ent}(g) \leq \int g \log g d \mu+\frac{1}{e} \leq \int g \log (1+g) d \mu+\frac{1}{e}=\frac{e+1}{e}
$$

Remark 3.5. Combining Lemma 3.2 with Lemma 3.4, and recalling Lemma 3.3, we may write a two-sided estimate

$$
\frac{1}{3}\left(\operatorname{Ent}\left(|f|^{q}\right)+\|f\|_{q}^{q}\right) \leq\|f\|_{N_{q}}^{q} \leq \operatorname{Ent}\left(|f|^{q}\right)+\|f\|_{q}^{q}
$$

In the case $q=2$, the proof of Proposition 3.1 is given in [B-G1]. One of the steps of the proof is based on an important observation of O. S. Rothaus $[R]$ who showed, in the setting of an arbitrary probability space $(\Omega, \mu)$, that

$$
\mathcal{L}_{2}(f)=\sup _{a \in \mathbf{R}} \operatorname{Ent}(|f+a|)^{2} \leq \operatorname{Ent}\left(f^{2}\right)+2 \int f^{2} d \mu
$$

whenever $\int f d \mu=0$. We need an appropriate generalization.
Lemma 3.6. For any function $f \in L_{N_{q}}(\Omega, \mu), 1 \leq q \leq 2$, such that $\int f d \mu=0$,

$$
\mathcal{L}_{q}(f) \leq 16\|f\|_{N_{q}}^{q} .
$$

Recall that, according to remark 3.5, $\|f\|_{N_{q}}^{q}$ is comparable with $\operatorname{Ent}\left(|f|^{q}\right)+$ $\|f\|_{q}^{q}$.

Proof of Lemma 3.6. By homogeneity, assume $\|f\|_{N_{q}}=\frac{1}{2}$.
Using Rothaus' estimate, we may write

$$
\begin{aligned}
\operatorname{Ent}\left(|f+a|^{q}\right) & \leq \operatorname{Ent}\left(\left(|f+a|^{\frac{q}{2}}-\int|f+a|^{\frac{q}{2}} d \mu\right)^{2}\right) \\
& +2 \int\left(|f+a|^{\frac{q}{2}}-\int|f+a|^{\frac{q}{2}} d \mu\right)^{2} d \mu
\end{aligned}
$$

The expression under the integral sign in the second term can be written as

$$
\left(\int\left(|f(\omega)+a|^{\frac{q}{2}}-\left|f\left(\omega^{\prime}\right)+a\right|^{\frac{q}{2}}\right) d \mu\left(\omega^{\prime}\right)\right)^{2}
$$

and by Cauchy-Schwarz, it is bounded by

$$
\int\left(|f(\omega)+a|^{\frac{q}{2}}-\left|f\left(\omega^{\prime}\right)+a\right|^{\frac{q}{2}}\right)^{2} d \mu\left(\omega^{\prime}\right)
$$

Since $q \leq 2$, we may apply a simple inequality $\left||x+a|^{\frac{q}{2}}-|y+a|^{\frac{q}{2}}\right|^{2} \leq|x-y|^{q}$, valid for all real $x, y$ and $a$, and estimate the above integral by $\int\left|f(\omega)-f\left(\omega^{\prime}\right)\right|^{q} d \mu\left(\omega^{\prime}\right)$. Thus, the second term does not exceed $2\left\|f(\omega)-f\left(\omega^{\prime}\right)\right\|_{q}^{q} \leq 2 \cdot 2^{q}\|f\|_{q}^{q} \leq \frac{5}{2}$ (according to Lemma 3.2 and by the assumption that $\left.\|f\|_{N_{q}}=\frac{1}{2}\right)$.

A similar argument applies to the first entropy term. Using a general inequality Ent $(g) \leq \int \Psi(g) d \mu+\frac{1}{e}$ and the property that $\Psi$ is a Young function, this term can be bounded by

$$
\begin{aligned}
& \int \Psi( \left.\left(|f(\omega)+a|^{\frac{q}{2}}-\int\left|f\left(\omega^{\prime}\right)+a\right|^{\frac{q}{2}} d \mu\left(\omega^{\prime}\right)\right)^{2}\right) d \mu(\omega)+\frac{1}{e} \\
& \quad \leq \int \Psi\left(\int\left|f(\omega)-f\left(\omega^{\prime}\right)\right|^{q} d \mu\left(\omega^{\prime}\right)\right) d \mu(\omega)+\frac{1}{e} \\
& \quad \leq \iint \Psi\left(\left|f(\omega)-f\left(\omega^{\prime}\right)\right|^{q}\right) d \mu(\omega) d \mu\left(\omega^{\prime}\right)+\frac{1}{e} \\
& \quad=\iint N_{q}\left(f(\omega)-f\left(\omega^{\prime}\right)\right) d \mu(\omega) d \mu\left(\omega^{\prime}\right)+\frac{1}{e}
\end{aligned}
$$

Since $\|f\|_{N_{q}}=\frac{1}{2}$, we have $\left\|f(\omega)-f\left(\omega^{\prime}\right)\right\|_{N_{q}} \leq\|f(\omega)\|_{N_{q}}+\left\|f\left(\omega^{\prime}\right)\right\|_{N_{q}} \leq 1$, so the last double integral is bounded by 1 .

Combining the estimates for the both terms, we get

$$
\operatorname{Ent}\left(|f+a|^{q}\right) \leq\left(1+\frac{1}{e}\right)+\frac{5}{2}=\left(\frac{1}{e}+3.5\right) \cdot 2^{q}\|f\|_{N_{q}}^{q} .
$$

The constant does not exceed 16, and this finishes the proof.
Proof of Proposition 3.1. First assume the Poincaré-type inequality (3.1) holds true with constant $C$. Then, by Lemma 3.6,

$$
\operatorname{Ent}\left(|f|^{q}\right) \leq \mathcal{L}_{q}(f) \leq 16\left\|f-\int f d \mu\right\|_{N_{q}}^{q} \leq 16 C \int|\nabla f|^{q} d \mu
$$

Hence, $\mathrm{LS}_{q}$ holds true with $c \leq 16 C$. Conversely, start with $\mathrm{LS}_{q}$ with constant $c$ and assume $\int f d \mu=0$. By Proposition 2.1, $\|f\|_{q}^{q} \leq \frac{4 c}{\log 2} \int|\nabla f|^{q} d \mu$. Therefore, by Lemma 3.3,

$$
\|f\|_{N_{q}}^{q} \leq \operatorname{Ent}\left(|f|^{q}\right)+\|f\|_{q}^{q} \leq\left(c+\frac{4 c}{\log 2}\right) \int|\nabla f|^{q} d \mu
$$

Hence, $C \leq c+\frac{4 c}{\log 2} \leq 7 c$. Proposition 3.1 is now proved.

## CHAPTER 4

## $\mathbf{L S}_{q}$ and Hardy-type inequalities on the line

When $\Omega=\mathbf{R}$, the Poincaré-type inequality (3.1) in the Orlicz space can further be reduced to a Hardy-type inequality for the Orlicz space norm $\|\cdot\|_{\Psi}$.

Let $m=m(\mu)$ denote a median of the probability measure $\mu$ on $\mathbf{R}$. It is any number such that $\mu\{(-\infty, m]\} \geq \frac{1}{2}$ and $\mu\{[m,+\infty)\} \geq \frac{1}{2}$.

For functions $f$ on $\mathbf{R}$, we denote $f_{0}=f \mathbf{1}_{(-\infty, m]}$ and $f_{1}=f \mathbf{1}_{[m,+\infty)}$.
Proposition 4.1. Let $1 \leq q \leq 2$. Assume that $\mathrm{LS}_{q}$-inequality holds true on the line with constant $c$, i.e., for any smooth function $f$ on $\mathbf{R}$,

$$
\begin{equation*}
\operatorname{Ent}\left(|f|^{q}\right) \leq c \int_{-\infty}^{\infty}\left|f^{\prime}(x)\right|^{q} d \mu(x) \tag{4.0.1}
\end{equation*}
$$

Then, for any smooth function $f$ on $\mathbf{R}$ with $f(m)=0$,

$$
\begin{equation*}
\left\|\left|f_{0}\right|^{q}\right\|_{\Psi}+\left\|\left|f_{1}\right|^{q}\right\|_{\Psi} \leq d \int_{-\infty}^{\infty}\left|f^{\prime}(x)\right|^{q} d \mu(x) \tag{4.0.2}
\end{equation*}
$$

with $d=42 c$. Conversely, (4.2) implies (4.1) with $c=144 d$.
The inequality (4.2) as an equivalent form for logarithmic Sobolev inequalities on the real line will be used in the next section to characterize probability measures $\mu$ satisfying (4.1). The proof of Proposition 4.1 is given in this section.

Clearly, (4.1) as well (4.2) may be extended from the class of all smooth functions to the class of all absolutely continuous functions (with the same condition $f(m)=0$ in the case of (4.2)).

To prove Proposition 4.1, we need an elementary general lemma.
Lemma 4.2. For any function $f \in L_{N_{q}}(\Omega, \mu), q \geq 1$,

$$
\left\|f-\int f d \mu\right\|_{N_{q}} \leq 2\|f\|_{N_{q}} .
$$

Conversely, if $1 \leq q \leq 2$ and $f=0$ on some set $A \subset \Omega$ with $\mu(A) \geq \frac{1}{2}$, we also have

$$
\|f\|_{N_{q}}^{q} \leq 6\left\|f-\int f d \mu\right\|_{N_{q}}^{q}
$$

Proof. As already noted in the proof of Lemma 3.2, the optimal constant in the general inequality $\|f\|_{1} \leq \alpha\|f\|_{N_{q}}$ is attained at $f=1$, so for all $f$ we have $\|f\|_{1}\|1\|_{N_{q}} \leq\|f\|_{N_{q}}$. Therefore,

$$
\left\|f-\int f d \mu\right\|_{N_{q}} \leq\|f\|_{N_{q}}+\left\|\int f d \mu\right\|_{N_{q}}=\|f\|_{N_{q}}+\|f\|_{1}\|1\|_{N_{q}} \leq 2\|f\|_{N_{q}} .
$$

This proves the first inequality. In order to prove the second one, with some absolute constant $\alpha$, we first need to derive an inequality

$$
\left|\int f d \mu\right| \leq \alpha\left\|f-\int f d \mu\right\|_{N_{q}}
$$

valid in the class of all functions $f$ in $L_{N_{q}}$ vanishing on the set $A$. Assuming that $a=\int f d \mu \neq 0$, the above is equivalent to

$$
\int N_{q}\left(\frac{\alpha(f-a)}{a}\right) d \mu \geq 1
$$

But on the set $A$, the function under the integral sign is just the constant $N_{q}(\alpha)$, so it suffices to pick $\alpha$ satisfying $N_{q}(\alpha) \geq 2$. Consequently, we may take $\alpha=$ $N_{q}^{-1}(2)=\left(\Psi^{-1}(2)\right)^{1 / q}$. Now,
$\|f\|_{N_{q}} \leq\left\|f-\int f d \mu\right\|_{N_{q}}+\left\|\int f d \mu\right\|_{N_{q}} \leq\left(1+\left(\frac{\Psi^{-1}(2)}{\Psi^{-1}(1)}\right)^{1 / q}\right)\left\|f-\int f d \mu\right\|_{N_{q}}$.
It remains to note that $\Psi^{-1}(1)>1, \Psi^{-1}(2)<2$, so $\left(1+\left(\frac{\Psi^{-1}(2)}{\Psi^{-1}(1)}\right)^{1 / q}\right)^{q}<$ $\left(1+2^{1 / q}\right)^{q}<6$, for all $q \in[1,2]$.

Proof of Proposition 4.1. First we derive (4.2) from (4.1). By Proposition 3.1, for any smooth function $f$ from $L_{N_{q}}(\mathbf{R}, \mu)$,

$$
\left\|f-\int f d \mu\right\|_{N_{q}}^{q} \leq 7 c \int_{-\infty}^{\infty}\left|f^{\prime}(x)\right|^{q} d \mu(x)
$$

Applying it to $f_{0}$ and $f_{1}$, we get

$$
\begin{aligned}
& \left\|f_{0}-\int f_{0} d \mu\right\|_{N_{q}}^{q} \leq 7 c \int_{(-\infty, m)}\left|f^{\prime}(x)\right|^{q} d \mu(x) \\
& \left\|f_{1}-\int f_{1} d \mu\right\|_{N_{q}}^{q} \leq 7 c \int_{(m,+\infty)}\left|f^{\prime}(x)\right|^{q} d \mu(x)
\end{aligned}
$$

Applying the second inequality of Lemma 4.2 and the general identity $\|g\|_{N_{q}}^{q}=$ $\left\||g|^{q}\right\|_{\Psi}$, we get

$$
\begin{aligned}
& \left\|\left|f_{0}\right|^{q}\right\|_{\Psi} \leq 42 c \int_{(-\infty, m)}\left|f^{\prime}(x)\right|^{q} d \mu(x) \\
& \left\|\left|f_{1}\right|^{q}\right\|_{\Psi} \leq 42 c \int_{(m,+\infty)}\left|f^{\prime}(x)\right|^{q} d \mu(x)
\end{aligned}
$$

Adding these inequalities, we obtain (4.2) with $d=42 c$.
To derive (4.1) from (4.2), first assume that $f(m)=0$. Since $f=f_{0}+f_{1}$,

$$
\left\|f-\int f d \mu\right\|_{N_{q}} \leq\left\|f_{0}-\int f_{0} d \mu\right\|_{N_{q}}+\left\|f_{1}-\int f_{1} d \mu\right\|_{N_{q}}
$$

On the other hand, by Lemma 3.2, $\left|\int f_{0} d \mu\right| \leq\left(\frac{5}{4}\right)^{1 / q}\left\|f_{0}\right\|_{N_{q}}$ and similarly for $f_{1}$. Hence,

$$
\begin{aligned}
\left\|f-\int f d \mu\right\|_{N_{q}}^{q} & \leq\left(1+\left(\frac{5}{4}\right)^{1 / q}\right)^{q}\left(\left\|f_{0}\right\|_{N_{q}}+\left\|f_{1}\right\|_{N_{q}}\right)^{q} \\
& \leq q\left(1+\left(\frac{5}{4}\right)^{1 / q}\right)^{q}\left(\left\|f_{0}\right\|_{N_{q}}^{q}+\left\|f_{1}\right\|_{N_{q}}^{q}\right) \\
& \leq 9\left(\left\|f_{0}\right\|_{N_{q}}^{q}+\left\|f_{1}\right\|_{N_{q}}^{q}\right)=9\left(\left\|\left|f_{0}\right|^{q}\right\|_{\Psi}+\left\|\left|f_{1}\right|^{q}\right\|_{\Psi}\right)
\end{aligned}
$$

By the assumption (4.2) applied to $f_{0}$ and $f_{1}$, we thus get

$$
\left\|f-\int f d \mu\right\|_{N_{q}}^{q} \leq 9 d \int_{-\infty}^{+\infty}\left|f^{\prime}(x)\right|^{q} d \mu(x)
$$

This inequality is invariant under translations $f \rightarrow f+$ const, so it holds without the condition $f(m)=0$. At last, by Proposition 3.1, it implies that

$$
\operatorname{Ent}\left(|f|^{q}\right) \leq 16 \cdot 9 d \int_{-\infty}^{+\infty}\left|f^{\prime}(x)\right|^{q} d \mu(x)
$$

and Proposition 4.1 follows.

## CHAPTER 5

## Probability measures satisfying $\mathbf{L S}_{q}$-inequalities on the real line

Using Proposition 4.1, we will give in this section a direct characterization of probability measures $\mu$ on the real line $\mathbf{R}$ satisfying the logarithmic Sobolev inequality

$$
\begin{equation*}
\operatorname{Ent}\left(|f|^{q}\right) \leq c \int_{-\infty}^{+\infty}\left|f^{\prime}(x)\right|^{q} d \mu(x) \tag{5.0.1}
\end{equation*}
$$

with some (finite) constant $c$ for all smooth functions $f$ on $\mathbf{R}$. This part is very similar to the case $q=2$ studied in [B-G1]: the basic tool is the following theorem due to M. Artola, G. Talenti and G. Tomaselli (cf. [Mu]) on the optimal constant $A_{q}=A_{q}(\nu, \lambda)$ in the Hardy-type inequality with weights

$$
\int_{0}^{+\infty}|f(x)|^{q} d \nu(x) \leq A_{q} \int_{0}^{+\infty}\left|f^{\prime}(x)\right|^{q} d \lambda(x)
$$

Here $f$ is supposed to be an arbitrary smooth function on $[0,+\infty)$ such that $f(0)=$ 0 , and $\nu$ and $\lambda$ are (non-negative) Borel measures on $[0,+\infty)$. Denote by $p_{\lambda}=p_{\lambda}(x)$ the absolutely continuous component of $\lambda$ with respect to Lebesgue measure, and define the constant $B_{q}=B_{q}(\nu, \lambda)$ as follows:

$$
B_{q}(\nu, \lambda)=\sup _{x>0} \nu([x,+\infty))\left(\int_{0}^{x} \frac{d t}{p_{\lambda}(t)^{1 /(q-1)}}\right)^{q-1}
$$

Theorem 5.1. ([Mu], [M2]) For any $q \in[1,+\infty), B_{q} \leq A_{q} \leq \frac{q^{q}}{(q-1)^{q-1}} B_{q}$.
The quantity $\frac{q^{q}}{(q-1)^{q-1}}$ represents an increasing function in $q$, so, in the range of interest $1 \leq q \leq 2$, it is bounded by 4 .

Theorem 5.1 has the following natural generalization. Consider a Borel measure $\nu$ on $[0,+\infty)$ and a Banach space $(X,\|\cdot\|)$ of Borel measurable functions on $[0,+\infty)$ (with usual factorization with respect to measure $\nu$ ) such that

1) $f \leq|g| \nu$-a.e., $g \in X$ implies $f \in X$ and $\|f\| \leq\|g\|$, for all Borel measurable functions $f$;
2) any pointwise non-decreasing sequence $f_{n}$ of non-negative functions in $X$ converging pointwise to a function $f \in X$ satisfies $\left\|f_{n}\right\| \rightarrow\|f\|$.

By property 1), $X$ is an ideal Banach space, and property 2) is called order semi continuity of the norm (cf. [K-A]). For an ideal Banach space $X$, the last
property is equivalent to a representation of the norm in $X$ in the form

$$
\begin{equation*}
\|f\|=\sup _{g \in \mathcal{G}} \int_{0}^{+\infty}|f(x)| g(x) d \nu(x) \tag{5.0.2}
\end{equation*}
$$

for some family $\mathcal{G}$ of non-negative Borel measurable functions $g$ on $[0,+\infty)$. This statement holds in the setting of an abstract probability space ( $\Omega, \nu$ ) ([K-A], p.190).

For these Banach spaces $X$, one immediately obtains by Theorem 5.1:
Corollary 5.2. Let $A_{q}=A_{q}(X, \lambda)$ be the optimal constant in the inequality

$$
\begin{equation*}
\left\||f|^{q}\right\| \leq A_{q} \int_{0}^{+\infty}\left|f^{\prime}(x)\right|^{q} d \lambda(x) \tag{5.0.3}
\end{equation*}
$$

where $f \in X$ is an arbitrary smooth function such that $f(0)=0$. Then, $B_{q} \leq A_{q} \leq$ $\frac{q^{q}}{(q-1)^{q-1}} B_{q}$, where

$$
B_{q}=B_{q}(X, \lambda)=\sup _{x>0}\left\|\mathbf{1}_{[x,+\infty)}\right\|\left(\int_{0}^{x} \frac{d t}{p_{\lambda}(t)^{1 /(q-1)}}\right)^{q-1}
$$

Indeed, the measures $\nu_{g}(d x)=g(x) \nu(d x)$ satisfy $B_{q}\left(\nu_{g}, \lambda\right) \leq A_{q}\left(\nu_{g}, \lambda\right) \leq$ $4 B_{q}\left(\nu_{g}, \lambda\right)$. Using the definitions (5.2) and (5.3), we get

$$
A_{q}(X, \lambda)=\sup _{g \in \mathcal{G}} A_{q}\left(\nu_{g}, \lambda\right), \quad B_{q}(X, \lambda)=\sup _{g \in \mathcal{G}} B_{q}\left(\nu_{g}, \lambda\right),
$$

hence, $B_{q}(X, \lambda) \leq A_{q}(X, \lambda) \leq \frac{q^{q}}{(q-1)^{q-1}} B_{q}(X, \lambda)$.
In particular, one may apply Corollary 5.2 to the Orlicz space $X=L_{\Psi}(\nu)$ which of course satisfies the properties 1 ) and 2) above. Recall that $\Psi(x)=|x| \log (1+|x|)$. For indicator functions, we get by definition of the Orlicz norm

$$
\left\|\mathbf{1}_{[x,+\infty)}\right\|_{\Psi}=\frac{1}{\Psi^{-1}\left(\frac{1}{\nu([x,+\infty))}\right)}
$$

where $\Psi^{-1}$ denotes the inverse function. Consequently, the optimal constant $A_{q}$ in (5.3) for the norm $\|\cdot\|=\|\cdot\|_{\Psi}$ and for the range $1 \leq q \leq 2$ can be estimated as follows:

$$
\begin{aligned}
& \sup _{x>0} \frac{1}{\Psi^{-1}\left(\frac{1}{\nu([x,+\infty))}\right)}\left(\int_{0}^{x} \frac{d t}{p_{\lambda}(t)^{1 /(q-1)}}\right)^{q-1} \leq A_{q} \\
\leq & 4 \sup _{x>0} \frac{1}{\Psi^{-1}\left(\frac{1}{\nu([x,+\infty))}\right)}\left(\int_{0}^{x} \frac{d t}{p_{\lambda}(t)^{1 /(q-1)}}\right)^{q-1} .
\end{aligned}
$$

In case $\nu((0,+\infty)) \leq \frac{1}{2}$, the marginal sides of these inequalities can be simplified by noting that, for any $x>0$, the number $t=\frac{1}{\nu((x,+\infty))}$ satisfies $t \geq \frac{1}{2}$. But in this range, we have $\frac{t}{2 \log t} \leq \Psi^{-1}(t) \leq \frac{2 t}{\log t}$, and so

$$
\frac{1}{2} \sup _{x>0} \nu([x,+\infty)) \log \frac{1}{\nu([x,+\infty))}\left(\int_{0}^{x} \frac{d t}{p_{\lambda}(t)^{1 /(q-1)}}\right)^{q-1} \leq A_{q}
$$

$$
\begin{equation*}
8 \sup _{x>0} \nu([x,+\infty)) \log \frac{1}{\nu([x,+\infty))}\left(\int_{0}^{x} \frac{d t}{p_{\lambda}(t)^{1 /(q-1)}}\right)^{q-1} \tag{5.0.4}
\end{equation*}
$$

Using Proposition 4.1, we are ready to formulate and to prove the main result of this section.

Let $\mu$ be a Borel probability measure on $\mathbf{R}$ with distribution function $F(x)=$ $\mu((-\infty, x])$, and density function $p=p(x), x \in \mathbf{R}$, for its absolutely continuous part with respect to Lebesgue measure. Denote by $m$ a median of $\mu$. Define

$$
\begin{gathered}
D_{0}(q)=\sup _{x<m} F(x) \log \frac{1}{F(x)}\left(\int_{x}^{m} \frac{d t}{p(t)^{1 /(q-1)}}\right)^{q-1} \\
D_{1}(q)=\sup _{x>m}(1-F(x)) \log \frac{1}{1-F(x)}\left(\int_{m}^{x} \frac{d t}{p(t)^{1 /(q-1)}}\right)^{q-1} .
\end{gathered}
$$

defining $D_{0}(q)$ and $D_{1}(q)$ to be zero in case $\mu((-\infty, m))=0$ or $\mu((m,+\infty))=0$, respectively.

Theorem 5.3. Let $1 \leq q \leq 2$. For some positive absolute constants $K_{0}$ and $K_{1}$, the optimal value of $c$ in the logarithmic Sobolev inequality (5.1) satisfies

$$
K_{0}\left(D_{0}(q)+D_{1}(q)\right) \leq c \leq K_{1}\left(D_{0}(q)+D_{1}(q)\right)
$$

Proof. Without loss of generality, let $m=0$. The inequality (4.2) may be divided into the two inequalities:

$$
\begin{align*}
& \left\|\left|f_{0}\right|^{q}\right\|_{\Psi} \leq d_{0} \int_{-\infty}^{0}\left|f_{0}^{\prime}(x)\right|^{q} d \mu(x)  \tag{5.0.5}\\
& \left\|\left|f_{1}\right|^{q}\right\|_{\Psi} \leq d_{1} \int_{0}^{+\infty}\left|f_{1}^{\prime}(x)\right|^{q} d \mu(x) \tag{5.0.6}
\end{align*}
$$

where $f_{0}$ and $f_{1}$ are arbitrary smooth functions defined on $(-\infty, 0]$ and $[0,+\infty)$, respectively, with $f_{0}(0)=f_{1}(0)=0$. More precisely, the optimal constants in (5.5)-(5.6) are connected with the optimal constant $d$ in (4.2) by

$$
d=\max \left\{d_{0}, d_{1}\right\}
$$

Now, according to (5.4) with $\nu=\lambda$ being the restriction of $\mu$ to $[0,+\infty)$, we have, for the optimal constant $d_{1}$,

$$
\begin{equation*}
\frac{1}{2} D_{1} \leq d_{1} \leq 8 D_{1} \tag{5.0.7}
\end{equation*}
$$

Here we also used an obvious fact that the supremum in the definition of $D_{1}$ will not change if one replaces the expression $\mu([x,+\infty))$ by $\mu((x,+\infty))$. Similarly, for the optimal constant $d_{0}$, we have

$$
\begin{equation*}
\frac{1}{2} D_{0} \leq d_{0} \leq 8 D_{0} \tag{5.0.8}
\end{equation*}
$$

By Proposition 4.1, (5.1) implies both (5.5) and (5.6) with $d_{i}=42 c$. Hence, by (5.7)-(5.8), we get

$$
\frac{1}{2} \max \left(D_{0}, D_{1}\right) \leq 42 c
$$

Therefore, $D_{0}+D_{1} \leq 168 c$ which implies Theorem 5.3 with $K_{0}=1 / 168$. On the other hand, again by Proposition $4.2, c \leq 144 d=144 \max \left(d_{0}, d_{1}\right)$, so, $c \leq$ $144 \cdot 8 \max \left(D_{0}, D_{1}\right) \leq 1152\left(D_{0}+D_{1}\right)$. Thus, one may choose $K_{1}=1152$ which proves Theorem 5.3.

A natural question arising in connection with Theorem 5.3 is how to determine whether or not $D_{0}(q)+D_{1}(q)$ is finite in terms of simple conditions on the density $p$ of the measure $\mu$. In the case $q=2$ this question was studied in [G-R], ([A-B-C-F-G-M-R-S]), and here we follow similar arguments to settle the general case.

Assume the density has the form $p(x)=e^{-U(x)}$ for a twice continuously differentiable function $U: \mathbf{R} \rightarrow \mathbf{R}$ (for short, $U \in C^{2}(\mathbf{R})$ ). We need the following elementary lemma appearing in [A-B-C-F-G-M-R-S], see pp. 107-109, as a corollary of a more general Proposition 6.4.1 (with further reference to [V-P]).

Lemma 5.4. Given a function $U \in C^{2}(\mathbf{R})$ such that $U^{\prime}(x)>0$ in $x \geq M$ for some real $M$, and $\frac{U^{\prime \prime}(x)}{\left(U^{\prime}(x)\right)^{2}} \rightarrow 0$, as $x \rightarrow+\infty$, we have, for all $m \in \mathbf{R}$,

$$
\lim _{x \rightarrow+\infty}\left(U^{\prime}(x) e^{-U(x)} \int_{m}^{x} e^{U(t)} d t\right)=1, \quad \lim _{x \rightarrow+\infty}\left(U^{\prime}(x) e^{U(x)} \int_{m}^{x} e^{-U(t)} d t\right)=1
$$

Applying this lemma to $p(x)=e^{-U(x)}$, we obtain that, as $x \rightarrow+\infty$ :

1) the function $1-F(x)=\int_{m}^{x} e^{-U(t)} d t$ is equivalent to $\frac{e^{-U(x)}}{U^{\prime}(x)}$;
2) the function $\log \frac{1}{1-F(x)}$ is equivalent to $U(x)+\log U^{\prime}(x)$;
3) $\left(\int_{m}^{x} p(t)^{-1 /(q-1)} d t\right)^{q-1}=\left(\int_{m}^{x} e^{U(t) /(q-1)} d t\right)^{q-1}$ is equivalent to $\frac{(q-1)^{q-1} e^{U(x)}}{\left(U^{\prime}(x)\right)^{q-1}}$.

Collecting these asymptotics together, we obtain that the function

$$
(1-F(x)) \log \frac{1}{1-F(x)}\left(\int_{m}^{x} \frac{d t}{p(t)^{1 /(q-1)}}\right)^{q-1}
$$

appearing in the definition of $D_{1}(q)$, is equivalent to $(q-1)^{q-1} \frac{U(x)+\log U^{\prime}(x)}{U^{\prime}(x)^{q}}$. We can now summarize (also taking into account an analogous description of the asymptotics near the point $-\infty$ ):

Theorem 5.5. Let $\mu$ be an absolutely continuous probability measure on $\mathbf{R}$ with density $p(x)=e^{-U(x)}$ where $U \in C^{2}(\mathbf{R})$. Assume that
a) $U^{\prime}(x)>0$ if $x \geq M$ and $U^{\prime}(x)<0$ in $x \leq-M$, for some real $M$;
b) $\frac{U^{\prime \prime}(x)}{U^{\prime}(x)^{2}} \rightarrow 0$, as $|x| \rightarrow \infty$.

Then, given $q \in(1,2]$, the measure $\mu$ satisfies the logarithmic Sobolev inequality (5.1) with some finite $c$ if and only if

$$
\begin{equation*}
\limsup _{|x| \rightarrow \infty} \frac{U(x)+\log \left|U^{\prime}(x)\right|}{\left|U^{\prime}(x)\right|^{q}}<+\infty . \tag{5.0.9}
\end{equation*}
$$

As an immediate consequence, we get:
Corollary 5.6. Given $q \in(1,2]$, the probability measure $\mu$ on $\mathbf{R}$ with density $p(x)=$ const $e^{-|x|^{p}}$ satisfies (5.1) if and only if $p \geq \frac{q}{q-1}$.

The measure $\mu$, with $U(x)=|x|^{p}+C$, is log-concave, and moreover it has convexity power $p$ (if $p \geq 2$ ). So, Corollary 5.6 can be proved by a different method, for example, on the basis of Brunn-Minkowski's inequality (cf. [B-L]).

However, Theorem 5.5 allows one to produce a variety of measures which satisfy the logarithmic Sobolev inequality (5.1) but which are not log-concave. Here is one possible example (in a spirit of [A-B-C-F-G-M-R-S]).

Example 5.7. Let $U$ be an even function of class $C^{2}(\mathbf{R})$ of the form

$$
U(x)=|x|^{p}+\alpha(x) \cos x+C
$$

such that $\alpha(x)=|x|^{p-1-\delta}$ for $|x|>1$, where $p>2$ and $\delta \in(0,1)$ are fixed parameters. For $x>1$ we have:

$$
U^{\prime}(x)=p x^{p-1}+(p-1-\delta) x^{p-2-\delta} \cos x-x^{p-1-\delta} \sin x,
$$

hence $U^{\prime}(x)$ is equivalent to $p x^{p-1}$, as $x \rightarrow+\infty$ (the first term dominates). Next,

$$
\begin{aligned}
U^{\prime \prime}(x)= & p(p-1) x^{p-2}+(p-1-\delta)(p-2-\delta) x^{p-3-\delta} \cos x \\
& -2(p-1-\delta) x^{p-2-\delta} \sin x-x^{p-1-\delta} \cos x
\end{aligned}
$$

Now, the last term is of most importance, or, more precisely, we have

$$
\frac{U^{\prime \prime}(x)}{x^{p-1-\delta}}=-\cos x+O\left(x^{-(1-\delta)}\right), \quad \text { as } x \rightarrow+\infty .
$$

Therefore, $U^{\prime \prime}(x)$ oscillates and attains any values between $-\infty$ to $+\infty$ as the variable $x$ varies along the positive half-axis. Thus, roughly speaking, the corresponding measure $\mu$ is very far from being log-concave. On the other hand,

$$
\frac{U^{\prime \prime}(x)}{U^{\prime}(x)^{2}}=O\left(x^{-(p-1+\delta)}\right)
$$

and the property (5.9) is true whenever $p \geq \frac{q}{q-1}$. Hence, for such $p$ the measure $\mu$ satisfies the logarithmic Sobolev inequality (5.1).

Remark 5.8. On the basis of Theorem 5.1 and Lemma 5.4 one can easily study the Poincaré-type inequality on the real line,

$$
\begin{equation*}
\int_{-\infty}^{+\infty}\left|f-\int_{-\infty}^{+\infty} f d \mu\right|^{q} d \mu \leq c_{q} \int_{-\infty}^{+\infty}\left|f^{\prime}(x)\right|^{q} d \mu(x) \tag{5.0.10}
\end{equation*}
$$

where $q>1$ is a parameter, and where $f$ is an arbitrary $\mu$-integrable absolutely continuous function on $\mathbf{R}$ (now, there is no need to restrict ourselves to the bound $q \leq 2$ ). Under the assumptions a) and b) of Theorem 5.5, we arrive at the following conclusion: the measure $\mu$ satisfies (5.10) with some finite $c_{q}$ if and only if

$$
\limsup _{|x| \rightarrow \infty} \frac{1}{\left|U^{\prime}(x)\right|^{q}}<+\infty
$$

Although it resembles (5.9), this property does not depend on $q$ (!) Therefore, all Poincaré-type inequalities (5.10) are equivalent to each other up to constants $c_{q}$ (of course, under the above regularity assumptions on the density $p$ ).

## CHAPTER 6

## Exponential integrability and perturbation of measures

Let us return to the abstract situation and to Theorem 1.3 considered in the introductory section. Thus, assume we are given the triple $(\Omega, \mu, \Gamma)$ satisfying $\mathrm{LS}_{q^{-}}$ inequality

$$
\begin{equation*}
\operatorname{Ent}\left(|f|^{q}\right) \leq c \int \Gamma(f)^{q} d \mu \tag{6.0.1}
\end{equation*}
$$

for every function $f \in \mathcal{A}$ (i.e., from the domain of $\Gamma$ ). As before, here $q$ is a given parameter, $1<q \leq 2$. We are going to generalize Theorem 1.3 by removing the boundedness assumption on the modulus of gradient.

Theorem 6.1. Under (6.1), for every bounded function $g \in A$ with $\mu$-mean zero,

$$
\begin{equation*}
\int e^{g} d \mu \leq \int e^{c p q^{-q} \Gamma(g)^{q}} d \mu \tag{6.0.2}
\end{equation*}
$$

where $p=\frac{q}{q-1}$ is conjugate to $q$.
As in Theorem 1.3, if we are dealing with probability metric spaces satisfying (6.1), the estimate (6.2) easily extends to non-bounded integrable, locally Lipschitz $g$. Thus, if $\Gamma(g)^{q}$ has finite exponential moments, then so does $g$. Of course, this property cannot be reached on the basis of Theorem 1.3 or the bound (1.3) on the Laplace transform,

$$
\begin{equation*}
\int e^{\lambda g} d \mu \leq \exp \left\{\frac{c}{q^{q}(q-1)} \lambda^{q}\right\}, \quad \lambda>0 \tag{6.0.3}
\end{equation*}
$$

Note that, if $\Gamma(g) \leq 1$ as in (6.3), the estimate (6.2) yields somewhat different:

$$
\int e^{\lambda g} d \mu \leq \exp \left\{\frac{c p}{q^{q}} \lambda^{q}\right\}, \quad \lambda>0
$$

Here the constant in the exponent $\frac{p}{q^{q}}=\frac{q}{q^{q}(q-1)}$ is $q \leq 2$ times worse than that in (6.3). Hence, Theorem 1.3 still has a small advantage over Theorem 6.1 in case $\Gamma(g) \leq 1$.

Proof of Theorem 6.1. The argument is similar to the one given in [B-G1] for the particular case $q=2$. Given a number $\alpha>0$, define $\beta>0$ by

$$
\begin{equation*}
\int e^{\alpha \Gamma(g)^{q}} d \mu=e^{\beta} . \tag{6.0.4}
\end{equation*}
$$

We use the following well-known representation for the entropy functional: for every measurable, bounded function $w \geq 0$ on $\Omega$,

$$
\begin{equation*}
\operatorname{Ent}(w)=\sup _{\int e^{h} d \mu \leq 1} \int w h d \mu \tag{6.0.5}
\end{equation*}
$$

Here the supremum may be taken over all measurable bounded function $h$ on $\Omega$ such that $\int e^{h} d \mu \leq 1$.

Let $w=e^{g}$ and let $h=\alpha \Gamma(g)^{q}-\beta$ so that, by (6.4), $\int e^{h} d \mu=1$. Hence, by (6.5), we get

$$
\int e^{g}\left(\alpha \Gamma(g)^{q}-\beta\right) d \mu \leq \operatorname{Ent}\left(e^{g}\right)
$$

On the other hand, by the log-Sobolev inequality (6.1) with $f=e^{g / q}$,

$$
\begin{equation*}
\operatorname{Ent}\left(e^{g}\right) \leq \frac{c}{q^{q}} \int \Gamma(g)^{q} e^{g} d \mu \tag{6.0.6}
\end{equation*}
$$

Combining the both estimates gives

$$
\left(\alpha-\frac{c}{q^{q}}\right) \int \Gamma(g)^{q} e^{g} d \mu \leq \beta \int e^{g} d \mu .
$$

Assuming $\alpha>\frac{c}{q^{q}}$, the left hand side can further be estimated from below by applying once more (6.6). Then we get

$$
\operatorname{Ent}\left(e^{g}\right) \leq C \beta \int e^{g} d \mu, \quad C=\frac{1}{\frac{\alpha q^{q}}{c}-1}
$$

Now, given $\lambda \in[0,1]$, let us write this inequality for the function $\lambda g$ in the place of $g$. Then,

$$
\begin{equation*}
\operatorname{Ent}\left(e^{\lambda g}\right) \leq C \beta(\lambda) \int e^{\lambda g} d \mu, \quad C=\frac{1}{\frac{\alpha q^{q}}{c}-1} \tag{6.0.7}
\end{equation*}
$$

where $\beta(\lambda)$ should be defined according to (6.4) as

$$
\int e^{\alpha \lambda^{q} \Gamma(g)^{q}} d \mu=e^{\beta(\lambda)}
$$

Since $\lambda \in[0,1]$, by Hölder's inequality, $\int e^{\alpha \lambda^{q} \Gamma(g)^{q}} d \mu \leq\left(\int e^{\alpha \Gamma(g)^{q}} d \mu\right)^{\lambda^{q}}$, so $\beta(\lambda) \leq$ $\beta \lambda^{q}$, where recall that $\beta(1)=\beta$. Therefore, from (6.7),

$$
\begin{equation*}
\operatorname{Ent}\left(e^{\lambda g}\right) \leq C \beta \lambda^{q} \int e^{\lambda g} d \mu, \quad C=\frac{1}{\frac{\alpha q^{q}}{c}-1} \tag{6.0.8}
\end{equation*}
$$

Now we can proceed as in the proof of Theorem 1.3. Write the Laplace transform on $g$ as $\int e^{\lambda g} d \mu=e^{\lambda v(\lambda)}$. The function $v$ is smooth in $\lambda>0$ and satisfies $v(0+)=$ 0 in view of the assumption that $g$ has $\mu$-mean zero. Recall that Ent $\left(e^{\lambda g}\right)=$ $\lambda^{2} v^{\prime}(\lambda) \int e^{\lambda g} d \mu$. Hence, by (6.8), $\lambda^{2} v^{\prime}(\lambda) \leq C \beta \lambda^{q}$ and

$$
v(\lambda)=\int_{0}^{\lambda} v^{\prime}(t) d t \leq C \beta \int_{0}^{\lambda} t^{q-2} d t=\frac{C \beta}{q-1} \lambda^{q-1} .
$$

Thus, we have obtained the estimate $\int e^{\lambda g} d \mu \leq \exp \left\{\frac{C \beta}{q-1} \lambda^{q}\right\}$ which for $\lambda=1$ becomes (recalling the definition of $\beta$ and $C$ )

$$
\begin{equation*}
\int e^{g} d \mu \leq e^{\frac{C \beta}{q-1}}=\left(\int e^{\alpha \Gamma(g)^{q}} d \mu\right)^{\frac{C}{q-1}}=\left(\int e^{\alpha \Gamma(g)^{q}} d \mu\right)^{\frac{1}{(q-1)\left(\frac{\alpha q q}{c}-1\right)}} \tag{6.0.9}
\end{equation*}
$$

It remains to equalize

$$
(q-1)\left(\frac{\alpha q^{q}}{c}-1\right)=1
$$

from which we easily find that $\alpha=\frac{c p}{q^{q}}$. For this value (6.9) coincides with (6.2), and this proves Theorem 6.1.

Next we briefly investigate which perturbation of measures preserve $\mathrm{LS}_{q}$-inequality. This will be useful for further purposes.

First of all we mention that following the idea $[\mathbf{H}-\mathbf{S}]$ in the general setting one has

Proposition 6.2. Suppose $\mu \in \operatorname{LS}_{q}(c)$ and let $d \mu_{\rho} \equiv \rho d \mu$ for some strictly positive bounded density $\rho$. Then $\mu_{\rho} \in \operatorname{LS}_{q}\left(c^{\prime}\right)$ with $c^{\prime} \equiv c \exp \{\sup (\log \rho)-\inf (\log \rho)\}$.

For the proof it is sufficient to notice that the nonlinear expression in $\mu_{\rho}$ can be achieved as an infimum of expectation of nonnegative function as follows

$$
\mu_{\rho} f^{q} \log \frac{f^{q}}{\mu f^{q}}=\inf _{t>0} \mu_{\rho}\left(f^{q} \log f^{q}-f^{q} \log t-f^{q}+t\right)
$$

Now we can remove the density $\rho$ and use $\mu \in \mathrm{LS}_{q}(c)$ and finally replace the expectation $\mu|\nabla f|_{q}^{q}$ by the one we need. The total price for these operations equals to the exponential of $\sup \log \rho-\inf \log \rho$. (We mention that in certain situation when more structure is available Proposition 6.2 can be strengthened, see e.g. [G-Z]. Naturally similar arguments works directly for the related Poincaré inequality $\mathrm{SG}_{q}$.)

The second perturbation formula applies to the situation in which we have a local gradient and involves unbounded and possibly locally singular perturbations to which Theorem 5.5 does not apply. This includes for example perturbations where

$$
\log \rho=\sum_{j} a_{j} V\left(\omega-b_{j}\right)+\text { const }
$$

with a singular function $V$, (singularity may be possibly growing with the dimension of the space), and certain infinite sequences of points $a_{j}$ and $b_{j}$.

Proposition 6.3. Suppose $\mu \in \operatorname{LS}_{q}(c)$ and let $d \mu_{\rho} \equiv \rho d \mu$ for a positive density $\rho$. If

$$
\mu_{\rho}\left(\exp \left\{\frac{2^{q-1} q^{-q} c}{\varepsilon}|\nabla \log \rho|_{q}^{q}\right\}\right)<\infty
$$

for some constant $\varepsilon \in(0,1)$ and $\mu_{\rho}$ satisfies $\mathrm{SG}_{q}$ with a constant $m^{\prime} \in(0, \infty)$, then $\mu_{\rho} \in \mathrm{LS}_{q}\left(c^{\prime}\right)$ for some $c^{\prime} \in(0, \infty)$.

Proof. Inserting $\rho^{1 / q} f$, defined with bounded smooth cylinder function $f$ normalized so that $\mu_{\rho}|f|^{q}=1$, into $\mathrm{LS}_{q}(c)$ for the measure $\mu$ after simple arguments (utilizing the Leibnitz rule for the local gradient) one arrives at the following bound

$$
\begin{equation*}
\mu_{\rho}|f|^{q} \log |f|^{q} \leq 2^{q-1} c \mu_{\rho}|\nabla f|_{q}^{q}+2^{q-1} q^{-q} c \mu_{\rho}\left(|f|^{q}|\nabla \log \rho|_{q}^{q}\right) \tag{6.0.10}
\end{equation*}
$$

Next utilizing Young inequality $a b \leq a \log a+2 e^{b}$, (as e.g. in [L6] ), we bound the second term on the right hand side of (6.0.10) as follows

$$
2^{q-1} q^{-q} c \mu_{\rho}\left(|f|^{q}|\nabla \log \rho|_{q}^{q}\right) \leq \varepsilon \mu_{\rho}|f|^{q} \log |f|^{q}+2 \mu_{\rho}\left(\exp \left\{\frac{2^{q-1} q^{-q} c}{\varepsilon}|\nabla \log \rho|_{q}^{q}\right\}\right)
$$

with arbitrary $\varepsilon \in(0,1)$. Inserting this into (6.0.10), after simple transformations we arrive at
(6.0.11) $\quad \mu_{\rho}|f|^{q} \log |f|^{q} \leq \frac{2^{q-1} c}{1-\varepsilon} \mu_{\rho}|\nabla f|_{q}^{q}+\frac{2}{1-\varepsilon} \mu_{\rho}\left(\exp \left\{\frac{2^{q-1} q^{-q} c}{\varepsilon}|\nabla \log \rho|_{q}^{q}\right\}\right)$

To complete the proof it is sufficient to use the adaptation of Rothaus lemma proven in Section 3 together with the assumed $\mathrm{SG}_{q}$ for the perturbed measure. Finally, after this stage we can relax our initial assumption $\mu_{\rho}|f|^{q}=1$.

## CHAPTER 7

## $\mathbf{L S} \mathbf{q}_{q}$-inequalities for Gibbs measures with Super Gaussian Tails.

In this section we show that $\mathrm{LS}_{q}$ is satisfied for a large class of nontrivial infinite dimensional measure. Our arguments generalize the ones developed to show $\mathrm{LS}_{2}$ and to keep the size of this section reasonable we concentrate only on the aspects which require modification. (We refer to $[\mathbf{G}-\mathbf{Z}]$ for a comprehensive description of $\mathrm{LS}_{2}$ case.)

For the present purposes it is convenient to define the infinite dimensional measures by introducing an a'priori given family of its regular conditional expectations which are given in terms of compatible family of probability kernels involving essentially finite dimensional integration as follows. For $\omega \in \mathbf{R}^{\mathbb{Z}^{d}}$ and a finite set $\Lambda \subset \subset \mathbb{Z}^{d}$ we define

$$
\begin{equation*}
E_{\Lambda}^{\omega}(f) \equiv \delta_{\omega}\left(\frac{\int f e^{-U_{\Lambda}} d \nu_{r, \Lambda}}{\int e^{-U_{\Lambda}} d \nu_{r, \Lambda}}\right) \tag{7.0.1}
\end{equation*}
$$

where $\nu_{r, \Lambda} \equiv \otimes_{i \in \Lambda} \nu_{r}\left(d x_{i}\right)$ and

$$
\begin{equation*}
U_{\Lambda} \equiv \sum_{X \cap \Lambda \neq \emptyset} \Phi_{X} \tag{7.0.2}
\end{equation*}
$$

where $\Phi_{X}(\omega) \equiv \psi_{X}\left(\omega_{X}\right)$ for a twice continuously differentiable function $\psi_{X}$ on $\mathbf{R}^{X}$ and $\omega_{X} \equiv\left(\omega_{i}\right)_{i \in X} ; \delta_{\omega}$ denotes the point mass concentrated at $\omega \in \mathbf{R}^{\mathbb{Z}^{d}}$. Such family called a local specification can be defined for any measure $\nu_{r}(d y) \equiv$ $\alpha_{r}^{-1} e^{-|y|^{r}} d y, r \in(1, \infty)$, on $\mathbf{R}$ and the interaction $\Phi \equiv\left\{\Phi_{X}\right\}_{X \subset \subset \mathbb{Z}^{d}}$ such that for any $\Lambda \subset \subset \mathbb{Z}^{d}$ one has

$$
\begin{equation*}
0<\int e^{-U_{\Lambda}} d \nu_{r, \Lambda}<\infty \tag{7.0.3}
\end{equation*}
$$

For simplicity we assume that

$$
\|\Phi\| \equiv \sup _{i \in \mathbf{Z}^{d}} \sum_{X \subset \subset \mathbf{Z}^{d} X \ni i,|X|>1}|X| \cdot\left\|\Phi_{X}\right\|_{\infty}<\infty
$$

although as one can see our results hold true in greater generality. By the very definition of the local specification $E_{\Lambda}^{\omega}(f)$ is measurable with respect the smallest $\sigma$-algebra generated by the coordinate functions $\left\{\omega_{i}: i \in \mathbb{Z}^{d} \backslash \Lambda\right\}$ and one has the following compatibility condition satisfied for any $\widetilde{\Lambda} \subset \Lambda$

$$
E_{\Lambda}^{\omega}\left(E_{\tilde{\Lambda}}(f)\right)=E_{\Lambda}^{\omega}(f)
$$

for any bounded measurable function $f$.

In this section we restrict ourselves to the case of super - Gaussian tails, that is we assume that $r \in(2, \infty)$. Under this condition simple perturbation arguments (mentioned in the previous section) show that $E_{\Lambda}$ inherits the $\mathrm{LS}_{q}$ property with some finite constant, that is in shorthand notation

$$
E_{\Lambda} \in \operatorname{LS}_{q}\left(C_{\Lambda}\right)
$$

with some finite constant $C_{\Lambda} \in(0, \infty)$.
Later on we assume that local specification satisfies a mixing condition as the one assumed in the proof of the $\mathrm{LS}_{2}$ and which can be formulated (in a strong form) as follows

## Strong Mixing Condition

$$
\left|E_{\Lambda}\left(\left(f-E_{\Lambda} f\right)\left(g-E_{\Lambda} g\right)\right)\right| \leq A(\nabla f, \nabla g) e^{-M \operatorname{dist}\left(\Lambda_{f}, \Lambda_{g}\right)}
$$

where $M \in(0, \infty)$ and $A(\nabla f, \nabla g)$ are constants independent of a finite set $\Lambda \subset \subset$ $\mathbb{Z}^{d}$ and $\operatorname{dist}\left(\Lambda_{f}, \Lambda_{g}\right)$ is the Euclidean distance between smallest sets $\Lambda_{f}$ and $\Lambda_{g}$ such that $f$ and $g$ are measurable with respect to smallest $\sigma$-algebra generated by coordinate functions $\left\{\omega_{i}: i \in \Lambda_{f}\right\}$ and $\left\{\omega_{i}: i \in \Lambda_{g}\right\}$, respectively; (see e.g. [G-Z]).

We mention that for example such condition is satisfied under the Dobrushin uniqueness ([G-Z]). Let $\mu$ be a Gibbs measure for the local specification $\left\{E_{\Lambda}\right\}_{\Lambda \subset \subset \mathbb{Z}^{d}}$ that is we have

$$
\mu E_{\Lambda} f=\mu f
$$

for any bounded measurable function $f$; here and later on we use the notation $\mu f \equiv \int f d \mu$.

We will prove that under suitable conditions on the local specification, the Gibbs measure satisfies $\mathrm{LS}_{q}$. To show this, we note first that under our general assumption this inequality is true for any kernel $E_{\Lambda}$, (one apply here similar perturbation lemma as used in case of $\mathrm{LS}_{2}$, see e.g. [G-Z]). The idea of the proof is to use a suitable telescopic representation of relative entropy associated to an appropriate sequence of transition matrices (or conditional expectations) as follows. Given a sequence of transition matrices, $E_{n}, n \in \mathbb{N}$ such that

$$
\begin{equation*}
\mu E_{n} F=\mu F \tag{7.0.4}
\end{equation*}
$$

for a probability measure $\mu$, we set $f_{n} \equiv\left(E_{n} f_{n-1}^{q}\right)^{\frac{1}{q}}$, with $f_{0} \equiv f$ for a nonnegative function $f$. Assuming that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} f_{n}=\mu f \tag{7.0.5}
\end{equation*}
$$

we note that

$$
\begin{equation*}
\mu f^{q} \log \frac{f^{q}}{\mu f^{q}}=\sum_{n \in \mathbb{N}} \mu\left(E_{n} f_{n-1}^{q} \log \frac{f_{n-1}^{q}}{E_{n} f_{n-1}^{q}}\right) \tag{7.0.6}
\end{equation*}
$$

Hence, if with some constant $C \in(0, \infty)$ independent of $n \in \mathbb{N}$ one has

$$
\begin{equation*}
E_{n} f_{n-1}^{q} \log \frac{f_{n-1}^{q}}{E_{n} f_{n-1}^{q}} \leq C E_{n}\left|\nabla_{n} f_{n-1}\right|^{q} \tag{7.0.7}
\end{equation*}
$$

(with a gradient $\nabla_{n}$ associated to the integration variables of $E_{n}$ ), using (7.0.6) and (7.0.4), we obtain

$$
\begin{equation*}
\mu f^{q} \log \frac{f^{q}}{\mu f^{q}} \leq C \sum_{n \in \mathbb{N}} \mu\left|\nabla_{n} f_{n-1}\right|^{q} \tag{7.0.8}
\end{equation*}
$$

Since $f_{n-1} \equiv\left(E_{n-1} f_{n-2}^{q}\right)^{\frac{1}{q}}$, if the probability measures $E_{n-1}$ do not depend on the $n$-th variables (as in the product case), by Minkowski inequality one gets

$$
\begin{equation*}
\left|\nabla_{n} f_{n-1}\right|^{q} \leq E_{n-1}\left|\nabla_{n} f_{n-2}\right|^{q} \tag{7.0.9}
\end{equation*}
$$

and by induction, (if $E_{k}, k \leq n-1$, are independent of the differentiation variable in question), one arrives at

$$
\begin{equation*}
\left|\nabla_{n} f_{n-1}\right|^{q} \leq E_{n-1}\left|\nabla_{n} f_{n-2}\right|^{q} \leq E_{n-1} \ldots E_{1}\left|\nabla_{n} f\right|^{q} \tag{7.0.10}
\end{equation*}
$$

Thus combining (7.0.8) - (7.0.10), we arrive at
Theorem 7.1. (Product property of $\mathrm{LS}_{q}$ ) If $\mu \equiv \otimes_{n \in \mathbb{N}} \mu_{n}$ with $\mu \in \mathrm{LS}_{q}(C)$, then

$$
\begin{equation*}
\mu f^{q} \log \frac{f^{q}}{\mu f^{q}} \leq C \sum_{n \in \mathbb{N}} \mu\left|\nabla_{n} f\right|^{q} \tag{7.0.11}
\end{equation*}
$$

If $E_{n-1}$ depends on the $n$-th variable, (as we would need to have to be able to include the Gibbs and PCA measures), we use similar arguments as described in [G-Z]. One introduces $E_{n-1}$ defined as an appropriate product and/or convolution of probability kernels $E_{\Delta+j}$ associated to suitable translations of a finite cube $\Delta$. To estimate $\left|\nabla_{n} f_{n-1}\right|^{q}$ it is necessary to get an estimate on $\left|\nabla_{i}\left(E_{\Lambda} F^{q}\right)^{\frac{1}{q}}\right|$ defined for a finite set $\Lambda \subset \mathbb{Z}^{d}$. One notes that for $i \notin \Lambda$, we have

$$
\begin{equation*}
\left|\nabla_{i}\left(E_{\Lambda} F^{q}\right)^{\frac{1}{q}}\right|=\left|\frac{1}{q}\left(E_{\Lambda} F^{q}\right)^{\frac{1}{q}-1} \nabla_{i}\left(E_{\Lambda} F^{q}\right)\right| \tag{7.0.12}
\end{equation*}
$$

Using the definition of $E_{\Lambda}$ as a perturbation of a product measure we can compute the derivative on the right hand side as follows

$$
\begin{equation*}
\left|\nabla_{i}\left(E_{\Lambda} F^{q}\right)\right|=\left|E_{\Lambda}\left(F^{q}\left(-\nabla_{i} U_{\Lambda}+E_{\Lambda} \nabla_{i} U_{\Lambda}\right)\right)+q E_{\Lambda}\left(F^{q-1} \nabla_{i} F\right)\right| \tag{7.0.13}
\end{equation*}
$$

Since the covariance vanishes if one of the involved functions is a constant, setting $F_{s} \equiv s F+(1-s) E_{\Lambda} F$, we get

$$
\begin{aligned}
\mid E_{\Lambda}\left(F^{q}( \right. & \left.-\nabla_{i} U_{\Lambda}+E_{\Lambda} \nabla_{i} U_{\Lambda}\right)\left|=\left|\int_{0}^{1} d s \frac{d}{d s} E_{\Lambda}\left(F_{s}^{q}\left(-\nabla_{i} U_{\Lambda}+E_{\Lambda} \nabla_{i} U_{\Lambda}\right)\right)\right|\right. \\
& \leq q\left(E_{\Lambda} F_{s}^{q}\left|-\nabla_{i} U_{\Lambda}+E_{\Lambda} \nabla_{i} U_{\Lambda}\right|^{p}\right)^{\frac{1}{p}}\left(E_{\Lambda}\left|F-E_{\Lambda} F\right|^{q}\right)^{\frac{1}{q}} \\
& \leq q \operatorname{var}_{\Lambda}\left(\nabla_{i} U_{\Lambda}\right) \sup _{s \in[0,1]}\left(E_{\Lambda} F_{s}^{q}\right)^{\frac{1}{p}}\left(E_{\Lambda}\left|F-E_{\Lambda} F\right|^{q}\right)^{\frac{1}{q}}
\end{aligned}
$$

Now using the $\mathbb{L}_{q}$ spectral gap inequality for the measure $E_{\Lambda}$, (which follows from $\mathrm{LS}_{q}$ by general arguments of Proposition 2.1), we arrive at (7.0.14)

$$
\mid E_{\Lambda}\left(F^{q}\left(-\nabla_{i} U_{\Lambda}+E_{\Lambda} \nabla_{i} U_{\Lambda}\right) \mid \leq\right.
$$

$$
\leq q\left(E_{\Lambda} F^{q}\right)^{1-\frac{1}{q}}\left[2\left(\frac{4 C_{\Lambda}}{\log 2}\right)^{\frac{1}{q}} \operatorname{var}_{\Lambda}\left(\nabla_{i} U_{\Lambda}\right) \sum_{j \in \Lambda}\left(E_{\Lambda}\left|\nabla_{j} F\right|^{q}\right)^{\frac{1}{q}}\right]
$$

Combining (7.0.12) - (7.0.14), via simple arguments (similar to the ones used in [G-Z] for the case of $\mathrm{LS}_{2}$ ), one arrives at a bound of the following form

$$
\begin{equation*}
\left|\nabla_{i}\left(E_{\Lambda} F^{q}\right)^{\frac{1}{q}}\right| \leq \left\lvert\,\left(E_{\Lambda}\left|\nabla_{i} F\right|^{q}\right)^{\frac{1}{q}}+C^{\prime} \operatorname{var}_{\Lambda}\left(\nabla_{i} U_{\Lambda}\right) \sum_{j \in \Lambda}\left(E_{\Lambda}\left|\nabla_{j} F\right|^{q}\right)^{\frac{1}{q}}\right. \tag{7.0.15}
\end{equation*}
$$

with some positive constant $C^{\prime}$ and $\operatorname{var}_{\Lambda}\left(\nabla_{i} U_{\Lambda}\right)$ denoting the total variation of $\nabla_{i} U_{\Lambda}$ with respect to the coordinates indexed by points in the finite set $\Lambda$. Note that in case $i \in \Lambda$, the derivative on the right hand side vanishes by definition of our kernel $E_{\Lambda}$. If $\operatorname{var}_{\Lambda}\left(\nabla_{i} U_{\Lambda}\right)$ is sufficiently small, the inequality (7.0.15) is sufficient to complete the proof of $\mathrm{LS}_{q}$ for the Gibbs measure $\mu$. However if the local specification satisfies the strong mixing property, this smallness requirement can be essentially weakened.

To take the full advantage of mixing, one may follow the same inductive idea of conditioning as used in the case of $\mathrm{LS}_{2}$ inequalities (see e.g. [G-Z]) to show appropriate relations for the present case. That is given a finite set $\Lambda$ we choose a subset $\widetilde{\Lambda} \subset \Lambda$ and then using the compatibility property of kernels we note first that

$$
\left|\nabla_{i}\left(E_{\Lambda} F^{q}\right)^{\frac{1}{q}}\right|=\left|\nabla_{i}\left(E_{\Lambda} E_{\widetilde{\Lambda}} F^{q}\right)^{\frac{1}{q}}\right| \equiv\left|\nabla_{i}\left(E_{\Lambda} \widetilde{F}^{q}\right)^{\frac{1}{q}}\right|
$$

Next we apply the above arguments with the function $\widetilde{F}$ where by a proper choice of the set $\widetilde{\Lambda}$, (so that $\widetilde{F}$ is localised possibly far from $\nabla_{i} U_{\Lambda}$ ), we can take the advantage of the mixing condition. Finite iteration of this procedure result with the improved sweeping out relations of the form

## (Sweeping out relations)

$$
\begin{equation*}
\left|\nabla_{i}\left(E_{\Lambda} F^{q}\right)^{\frac{1}{q}}\right| \leq \left\lvert\,\left(E_{\Lambda}\left|\nabla_{i} F\right|^{q}\right)^{\frac{1}{q}}+\sum_{j \in \Lambda} C_{i j}\left(E_{\Lambda}\left|\nabla_{j} F\right|^{q}\right)^{\frac{1}{q}}\right. \tag{7.0.16}
\end{equation*}
$$

where $C_{i j} \in[0, \infty)$ decay in a suitable way with the growth of $\operatorname{dist}(i, j)$ (see [G-Z] for details).

Using the sweeping out relations one shows the following contraction bound involving the transitions matrices

$$
\begin{equation*}
\left|\nabla\left(E_{n-1} g^{q}\right)^{\frac{1}{q}}\right|_{q}^{q} \leq \lambda E_{n-1}|\nabla g|_{q}^{q} \tag{7.0.17}
\end{equation*}
$$

with a constant $\lambda \in(0,1)$ independent of $n \in \mathbf{N}$, where $|\nabla g|_{q}^{q} \equiv \sum_{j}\left|\nabla_{j} g\right|^{q}$. With $g=f_{n-2}$ and the help of simple induction, this yields

$$
\begin{equation*}
\left|\nabla\left(E_{n-1} f_{n-2}\right)^{\frac{1}{q}}\right|_{q}^{q} \leq \lambda^{n-1} E_{n-1}|\nabla f|_{q}^{q} \tag{7.0.18}
\end{equation*}
$$

Application of (7.0.18) together with (7.0.11) implies the following $\mathrm{LS}_{q}$ inequality

$$
\mu f^{q} \log \frac{f^{q}}{\mu f^{q}} \leq \frac{C_{q}}{1-\lambda} \mu|\nabla f|_{q}^{q}
$$

This completes the proof of the following result.

Theorem 7.2. ( $\mathrm{LS}_{q}$ for Gibbs measures) Let $\mu$ be a Gibbs measure corresponding to the local specification $\left\{E_{\Lambda}\right\}_{\Lambda \subset \subset \mathbb{Z}^{d}}$ defined by (7.0.1)-(7.0.3) with the reference product measure with super-Gaussian tails. If the local specification satisfies the mixing condition, then there is a constant $c_{q} \in(0, \infty)$ such that
$\left(\mathrm{LS}_{q}(\mu)\right)$

$$
\operatorname{Ent}_{\mu}(f) \leq c_{q} \cdot \mu\left|\nabla f^{\frac{1}{q}}\right|_{q}^{q}
$$

for all nonnegative functions $f$ for which the right hand side is finite.

Remark 7.3. We also remark that essentially the same arguments work also in the case of PCA measures. The case of perturbations by unbounded potentials can also be included, though they require lengthy and technically slightly more involved description.

The above theorem generalises a result of $[\mathbf{B}-\mathbf{L}]$ where a corresponding result was proven under a very strong log-concavity assumptions of the finite dimensional measures. That is, for a probability measure having a density $\exp (-V) / Z$ with respect to a finite dimensional Lebesgue measure with

$$
V(x)+V(y)-2 V\left(\frac{x+y}{2}\right) \geq \mathrm{const}\|x-y\|^{p}
$$

for any $x$ and $y$, such probability measure satisfies $\mathrm{LS}_{q}$ (with dual norm on the tangent space).

We remark that in the considered case of local specification with super Gaussian tails the following Generalised Nash inequality was proved in $[\mathbf{Z 2}]$ for the corresponding Gibbs measures following natural adaptation of classical Sobolev inequality for probability measures by $[\mathbf{R o s}]$

$$
\begin{equation*}
\mu f^{2} \log \frac{f^{2}}{\mu f^{2}} \leq C \frac{2^{q}}{q^{q}} \sum_{j}\left(\mu\left|\nabla_{j} f\right|^{2}\right)^{\frac{q}{2}}\left(\mu(f-\mu f)^{2}\right)^{\frac{2-q}{2}} \tag{q}
\end{equation*}
$$

It was proved there that such inequality produces the correct exponential bounds for Lipschitz function.

The following result explains the relation of $\mathrm{GN}_{q}$ inequality and $\mathrm{LS}_{q}$
Theorem 7.4. Suppose $r \in[2, \infty)$ and $\frac{1}{r}+\frac{1}{q}=1$. If $\mathrm{LS}_{q}$ is satisfied with a coefficient $C \in(0, \infty)$ and the following spectral gap inequality holds

$$
\begin{equation*}
\mu(f-\mu f)^{2} \leq m^{-1} \mu|\nabla f|_{2}^{2} \tag{2}
\end{equation*}
$$

with some constant $m \in(0, \infty)$, then the Generalised Nash inequality $\left(G N_{q}\right)$ is true with some constant $C \in(0, \infty)$ for all nonnegative functions $f$ for which the right hand side is finite.

Remark 7.5. Note that, although $\mathrm{SG}_{2}$ follows from $\mathrm{SG}_{q}$ in finite dimension (and the last is an abstract consequence of $\mathrm{LS}_{q}$ ), in infinite dimensional setting we need to be more careful as $l_{q}$ and $l_{2}$ norms are not equivalent.

We also remark that if dimension of the underlying space is finite, then $\mathrm{GN}_{q}$, $q \in(1,2)$, implies $\mathrm{LS}_{2}$. But in infinite dimensional situation this is in general not true.

Proof of Theorem 7.4. Suppose $\operatorname{LS}_{q}$ is satisfied with $q \in(1,2)$. Then substituting $f^{\frac{2}{q}}$ one gets

$$
\begin{equation*}
\mu f^{2} \log \frac{f^{2}}{\mu f^{2}} \leq C \mu\left|\nabla f^{\frac{2}{q}}\right|_{q}^{q}=C \mu \sum_{j}\left|\nabla_{j} f^{\frac{2}{q}}\right|^{q} \tag{7.0.19}
\end{equation*}
$$

We have

$$
\mu \sum_{j}\left|\nabla_{j} f^{\frac{2}{q}}\right|^{q}=\frac{2^{q}}{q^{q}} \mu\left(f^{2-q}\left(\sum_{j}\left|\nabla_{j} f\right|^{q}\right)\right)
$$

Since $q \leq 2$ and using Hölder inequality, we get
(7.0.20) $\mu \sum_{j}\left|\nabla_{j} f^{\frac{2}{q}}\right|^{q} \leq \frac{2^{q}}{q^{q}}\left(\mu|\nabla f|_{q}^{2}\right)^{\frac{q}{2}}\left(\mu f^{2}\right)^{\frac{2-q}{2}} \leq \frac{2^{q}}{q^{q}} \sum_{j}\left(\mu\left|\nabla_{j} f\right|^{2}\right)^{\frac{q}{2}}\left(\mu f^{2}\right)^{\frac{2-q}{2}}$

Next note that

$$
\begin{equation*}
\mu f^{2} \log \frac{f^{2}}{\mu f^{2}} \leq \mu(f-\mu f)^{2} \log \frac{(f-\mu f)^{2}}{\mu(f-\mu f)^{2}}+2 \mu(f-\mu f)^{2} \tag{7.0.21}
\end{equation*}
$$

and, if the Poincaré inequality is true, we have

$$
\begin{align*}
\mu(f-\mu f)^{2} & =\left(\mu(f-\mu f)^{2}\right)^{\frac{q}{2}}\left(\mu(f-\mu f)^{2}\right)^{\frac{2-q}{2}}  \tag{7.0.22}\\
& \leq\left(m^{-1} \mu|\nabla f|_{2}^{2}\right)^{\frac{q}{2}}\left(\mu(f-\mu f)^{2}\right)^{\frac{2-q}{2}}
\end{align*}
$$

Thus applying (7.0.19) - (7.0.20) with $|f-\mu f|$ instead of $f$ to estimate the first term on the right hand side of (7.0.21) and using (7.0.22) to estimate the second one, we arrive at the following bound

$$
\mu f^{2} \log \frac{f^{2}}{\mu f^{2}} \leq\left(\frac{2^{q}}{q^{q}} C+2 m^{-\frac{q}{2}}\right) \sum_{j}\left(\mu\left|\nabla_{j} f\right|^{2}\right)^{\frac{q}{2}}\left(\mu(f-\mu f)^{2}\right)^{\frac{2-q}{2}}
$$

## CHAPTER 8

## $\mathbf{L S}_{q}$-inequalities and Markov semigroups

In this section we discuss consequences of $\mathrm{LS}_{q}$ for the semigroups preserving the related measure. The first result provides a relation to the ultracontractivity property of a semigroup generated by a Dirichlet operator satisfying

$$
\mu(f \mathbf{L} f)=-\mu|\nabla f|_{2}^{2}
$$

The result holds in a general setting including both discrete as well as continuous spaces.

Theorem 8.1. Suppose the dimension of the underlying space is finite and assume that $r \in(2, \infty)$ and $\frac{1}{r}+\frac{1}{q}=1$. If $\mathrm{LS}_{q}$ is satisfied, then the following inequality is true for any $p \in(2, \infty)$ with generator $\mathbf{L}$ corresponding to the $L_{2}$ Dirichlet form
$\left(\mathrm{GLS}_{q}\right)$

$$
\mu f^{p} \log \frac{f^{p}}{\mu f^{p}} \leq-p \varepsilon(p) \mu\left(f^{p-1} \mathbf{L} f\right)+p \eta(p) \mu f^{p}
$$

where

$$
\varepsilon(p) \equiv \frac{a}{(p-1)} p^{-\delta} \quad \text { and } \quad \eta(p) \equiv b p^{-2+\delta q /(2-q)}
$$

with some constants $a, b \in(0, \infty)$ and $\delta \in\left(0, \frac{2-q}{q}\right)$, for all nonnegative functions $f$ for which the right hand side is finite.

Therefore the semigroup $P_{t} \equiv e^{t \mathbf{L}}$ is ultracontractive, that is for any $p \in[1, \infty)$ and $t>0$ the operator $P_{t}: L^{p} \longrightarrow L^{\infty}$ is bounded.

Proof. If the dimension of the space is finite, by similar arguments as in the proof of Theorem 7.4 , one gets

$$
\mu f^{2} \log \frac{f^{2}}{\mu f^{2}} \leq C^{\prime}\left(\mu|\nabla f|_{2}^{2}\right)^{\frac{q}{2}}\left(\mu f^{2}\right)^{\frac{2-q}{2}}
$$

and hence for any constant $\alpha \in(0, \infty)$ we have

$$
\mu f^{2} \log \frac{f^{2}}{\mu f^{2}} \leq C^{\prime} \frac{q}{2} \alpha \mu|\nabla f|_{2}^{2}+C^{\prime} \frac{2-q}{q} \alpha^{-\frac{q}{2-q}} \mu f^{2}
$$

Substituting $f^{\frac{p}{2}}$, with $p \in(2, \infty)$, and using the well known inequality

$$
\mu\left|\nabla f^{\frac{p}{2}}\right|_{2}^{2} \leq-\frac{p^{2}}{4(p-1)} \mu\left(f^{p-1} \mathbf{L} f\right)
$$

for the Dirichlet generator $\mathbf{L}$, we get, with appropriate constants $a, b \in(0, \infty)$ independent of $p$, the following bound

$$
\mu f^{p} \log \frac{f^{p}}{\mu f^{p}} \leq-a \frac{p^{2}}{4(p-1)} \alpha \mu\left(f^{p-1} \mathbf{L} f\right)+b \alpha^{-\frac{q}{2-q}} \mu f^{p}
$$

We will choose

$$
\alpha \equiv \alpha(p) \equiv p^{-\delta}
$$

with some $\delta \in\left(0, \frac{2-q}{q}\right)$ and with this choice we have the following inequality

$$
\mu f^{p} \log \frac{f^{p}}{\mu f^{p}} \leq-p \varepsilon(p) \mu\left(f^{p-1} \mathbf{L} f\right)+p \eta(p) \mu f^{p}
$$

where

$$
\varepsilon(p) \equiv a \frac{p^{1-\delta}}{4(p-1)} \quad \text { and } \quad \eta(p) \equiv b p^{-1+\frac{\delta q}{2-q}}
$$

It is not difficult to check that

$$
\int_{2}^{\infty} p^{-1} \varepsilon(p) d p<\infty \quad \text { and } \quad \int_{2}^{\infty} p^{-1} \eta(p) d p<\infty
$$

which following [G1] implies the ultracontractivity of the semigroup $P_{t}=e^{t \mathbf{L}}$.
Remark 8.2. We recall following [D-S] (see also $[\mathbf{D}]$ and $[$ Ros $]$ ) that in finite dimensions the distributions with super-Gaussian tails define Dirichlet generators of ultracontractive semigroups.

### 8.1. Ergodicity for nonlinear semigroups

Let $T_{t} \equiv e^{t \mathcal{L}}$ be a semigroup with generator $\mathcal{L}$ defined by

$$
\mathcal{L} \equiv \sum_{j} \mathcal{L}_{j}
$$

where $\mathcal{L}_{j}$ satisfies

$$
\mu\left(g \mathcal{L}_{j}(f)\right)=-\mu\left(\nabla_{j} g \cdot \mathbf{v}_{j}\left(\left|\nabla_{j} f\right|\right) \nabla_{j} f\right)
$$

with a nonnegative function $\mathbf{v}_{j}$.
Theorem 8.3. ( $\mathbf{L}_{2}$ ergodicity) If the following Spectral Gap inequality is satisfied

$$
\begin{equation*}
m \cdot \mu|g-\mu g|^{q} \leq \sum_{j} \mu\left(\mathbf{v}_{j}\left(\left|\nabla_{j} g\right|\right) \cdot\left|\nabla_{j} g\right|^{2}\right) \tag{SG}
\end{equation*}
$$

with some constants $q \in(1, \infty)$ and $m \in(0, \infty)$, then following estimate is true

$$
\begin{equation*}
\mu\left(T_{t} f-\mu f\right)^{2} \leq D \cdot t^{-\beta} \tag{8.1.1}
\end{equation*}
$$

with some constants $D, \beta \in(0, \infty)$ for all nonnegative functions $f$ in. $\circ$
Remark 8.4. Recall that by Theorem 2.1, $\mathrm{SG}_{q}$ follows from $\mathrm{LS}_{q}$, (but if the dimension of the underlying space is infinite $\mathrm{SG}_{q}$ does not imply $\mathrm{SG}_{2}$ ).

Proof. Let $f_{t} \equiv T_{t} f$, for a smooth bounded function $f$ such that $\mu f=0$. We have formally

$$
\begin{equation*}
\frac{d}{d t} \mu f_{t}^{2}=2 \mu f_{t} \mathcal{L} f_{t}=-2 \sum_{j} \mu\left(\mathbf{v}_{j}\left(\left|\nabla_{j} f_{t}\right|\right) \cdot\left|\nabla_{j} f_{t}\right|^{2}\right) \tag{8.1.2}
\end{equation*}
$$

(Since a'priori the question whether $f_{t}$ is in the domain of $\mathcal{L}$ may be problematic, one may need to use here appropriate approximation arguments.) If the following spectral gap inequality is satisfied

$$
m \cdot \mu|g-\mu g|^{q} \leq \sum_{j} \mu\left(\mathbf{v}_{j}\left(\left|\nabla_{j} g\right|\right) \cdot\left|\nabla_{j} g\right|^{2}\right)
$$

then one has

$$
\begin{equation*}
\frac{d}{d t} \mu f_{t}^{2} \leq-2 m \cdot \mu\left|f_{t}\right|^{q} \tag{8.1.3}
\end{equation*}
$$

Suppose $q \leq 2$, then we use Hölder inequality to get the following bound

$$
\mu\left|f_{t}\right|^{2} \leq\left(\mu f_{t}^{q}\right)^{\frac{1}{u}}\left(\mu f_{t}^{(2-s) w}\right)^{\frac{1}{w}} \leq\left(\mu f_{t}^{q}\right)^{\frac{1}{u}}\left(\mu f^{(2-s) w}\right)^{\frac{1}{w}}
$$

with arbitrary $s \in(0,2)$ and $u, w \in(1, \infty)$ satisfying $s u=q, \frac{1}{u}+\frac{1}{w}=1$. Hence we have

$$
\begin{equation*}
\left(\frac{\mu\left|f_{t}\right|^{2}}{\left(\mu f^{(2-s) w}\right)^{\frac{1}{w}}}\right)^{u} \leq \mu f_{t}^{q} \tag{8.1.4}
\end{equation*}
$$

We use this to bound the right hand side of (8.1.3) as follows

$$
\begin{equation*}
\frac{d}{d t} \mu f_{t}^{2} \leq-2 m \cdot\left(\frac{\mu\left|f_{t}\right|^{2}}{\left(\mu f^{(2-s) w}\right)^{\frac{1}{w}}}\right)^{u} \tag{8.1.5}
\end{equation*}
$$

Solving this differential inequality, we arrive at

$$
\begin{align*}
\mu f_{t}^{2} & \leq\left(\left(\mu f^{2}\right)^{1-u}+2 m(u-1) \cdot\left(\frac{1}{\left(\mu f^{(2-s) w}\right)^{\frac{1}{w}}}\right)^{u} t\right)^{-\frac{1}{u-1}}  \tag{8.1.6}\\
& \leq\left(\frac{\left(\mu f^{(2-s) w}\right)^{\frac{u}{w}}}{2 m(u-1)}\right)^{\frac{1}{u-1}} \cdot t^{-\frac{1}{u-1}}
\end{align*}
$$

whence the bound follows.

Theorem 8.5. (Ergodicity in the Entropy Sense) Suppose

$$
\mu(f \mathcal{L}(f)) \equiv-\sum_{j \in \Lambda} \mu\left(\nabla_{j} f \cdot \mathbf{v}_{j}\left(\left|\nabla_{j} f\right|\right) \nabla_{j} f\right)
$$

with $v_{j}(y) \geq|y|^{\kappa-2}$ for some $\kappa \geq 2$. If the following $\mathrm{LS}_{q}$ inequality is satisfied

$$
\begin{equation*}
\operatorname{Ent}_{\mu}(f) \leq q^{q} C \cdot \mu\left|\nabla f^{\frac{1}{q}}\right|_{q}^{q} \tag{q}
\end{equation*}
$$

with $q \in(1,2), \frac{1}{q}+\frac{1}{\kappa}=1$, and some constant $C \in(0, \infty)$ and $|\Lambda|<\infty$, then following estimate is true

$$
\begin{equation*}
\operatorname{Ent}_{\mu}\left(f_{t}\right) \leq|\Lambda| C^{-\frac{\kappa}{\kappa-q}}\left(\frac{\kappa-q}{q}\right)^{\frac{q}{\kappa-q}} \cdot t^{-\frac{q}{\kappa-q}} \tag{8.1.7}
\end{equation*}
$$

for all nonnegative functions $f$.

Proof. Let $f_{t} \equiv T_{t} f \equiv e^{t \mathcal{L}}$, for a smooth bounded function $f \geq 0$. Setting $S_{t} \equiv \operatorname{Ent}_{\mu}\left(f_{t}\right)$, we have

$$
\begin{equation*}
\frac{d}{d t} S_{t}=\mu \mathcal{L} f_{t} \log \frac{f_{t}}{\mu f_{t}}=-\sum_{j} \mu\left(f_{t}^{-1} \mathbf{v}_{j}\left(\left|\nabla_{j} f_{t}\right|\right) \cdot\left|\nabla_{j} f_{t}\right|^{2}\right) \tag{8.1.8}
\end{equation*}
$$

If $\mathbf{v}_{j}(y) \geq|y|^{\kappa-2}$, then we have

$$
\begin{equation*}
\mu\left(f_{t}^{-1} \mathbf{v}_{j}\left(\left|\nabla_{j} f_{t}\right|\right) \cdot\left|\nabla_{j} f_{t}\right|^{2}\right) \geq\left(1-\frac{1}{\kappa}\right)^{-\kappa} \cdot \mu\left|\nabla_{j} f_{t}^{1-\frac{1}{\kappa}}\right|^{\kappa} \tag{8.1.9}
\end{equation*}
$$

Thus, if $1-\frac{1}{\kappa}=\frac{1}{q}$, with $q \in(1,2)$, we have

$$
\begin{equation*}
\frac{d}{d t} S_{t} \leq-q^{\kappa} \cdot \sum_{j} \mu\left|\nabla_{j} f_{t}^{\frac{1}{q}}\right|^{\kappa} \tag{8.1.10}
\end{equation*}
$$

Since $1<q \leq \kappa$, we can use Hölder inequality for $l_{p}(\Lambda)$ norms to get

$$
\begin{equation*}
\frac{d}{d t} S_{t} \leq-q^{\kappa}|\Lambda|^{1-\frac{\kappa}{q}} \cdot\left(\sum_{j} \mu\left|\nabla_{j} f_{t}^{\frac{1}{q}}\right|^{q}\right)^{\frac{\kappa}{q}} \tag{8.1.11}
\end{equation*}
$$

If now the following $\mathrm{LS}_{q}$ inequality is satisfied

$$
\operatorname{Ent}_{\mu}(g) \leq q^{q} C \sum_{j} \mu\left|\nabla_{j} g^{\frac{1}{q}}\right|^{q}
$$

then using (8.1.11), we obtain

$$
\begin{equation*}
\frac{d}{d t} S_{t} \leq-|\Lambda|^{1-\frac{\kappa}{q}} C^{-\frac{\kappa}{q}} \cdot S_{t}^{\frac{\kappa}{q}} \tag{8.1.12}
\end{equation*}
$$

Solving this differential inequality one arrives at the following bound

$$
\begin{equation*}
S_{t} \leq\left(S_{0}^{-\frac{\kappa-q}{q}}+|\Lambda|^{1-\frac{\kappa}{q}} C^{-\frac{\kappa}{q}}\left(\frac{\kappa}{q}-1\right) \cdot t\right)^{-\frac{q}{\kappa-q}} \tag{8.1.13}
\end{equation*}
$$

which implies (8.1.7).
Remark 8.6. We remark that as $q \longrightarrow 2$ the right hand side of (8.1.13) converges to the known exponential decay of entropy which is independent of $\Lambda$.

## CHAPTER 9

## Isoperimetry

Given a convex body $K$ in $\mathbf{R}^{n}$ (a convex, compact, $n$-dimensional set with non-empty interior), denote by $\mu$ the corresponding normalised Lebesgue measure, that is, the uniform distribution on $K$. Thus, it has density

$$
\frac{d \mu(x)}{d x}=\frac{1_{K}(x)}{\operatorname{vol}_{n}(K)}, \quad x \in \mathbf{R}^{n},
$$

with respect to Lebesgue measure on $\mathbf{R}^{n}$, where we use $1_{K}$ to denote the indicator function of a set $K$. The canonical relative isoperimetric problem for $K$ is to minimise the quantity $\mu\left(A+h B_{2}\right)$ provided that $\mu(A)$ and $h>0$ are fixed ( $A$ is a Borel subset of $\mathbf{R}^{n}$ and $B_{2}$ is the unit Euclidean ball with centre at the origin). Another, formally weaker problem asks how to minimise the relative $\mu$-perimeter

$$
\begin{equation*}
\mu_{2}^{+}(A)=\liminf _{h \downarrow 0} \frac{\mu\left(A+h B_{2}\right)-\mu(A)}{h}, \tag{9.0.1}
\end{equation*}
$$

that is, how to find the so-called isoperimetric function for $\mu$,

$$
I_{\mu}(t)=\inf _{\mu(A)=t} \mu^{+}(A), \quad 0<t<1
$$

The problem makes sense for any probability measure $\mu$ on $\mathbf{R}^{n}$ and is of a large interest in Analysis and Probability.

Furthermore, we often obtain a different interesting problem when considering a different enlargement of sets by replacing the ball $B_{2}$ with other convex bodies $B$ in $\mathbf{R}^{n}$ such as for example $B_{p}$, the unit balls with respect to $\ell^{p}$ metric, $(1 \leq p \leq$ $+\infty)$. Appropriately, one should understand the notion of the perimeter and of the isoperimetric function.

Even in the traditional situation (the case of a convex body and of the Euclidean distance), the isoperimetric problem seems to be very far from being solved. One usually tries to study suitable families of convex bodies or some isoperimetric-type inequalities such as of the Cheeger-type

$$
\begin{equation*}
\mu_{2}^{+}(A) \geq c \min \{\mu(A), 1-\mu(A)\}, \quad A \subset \mathbf{R}^{n} \text { Borel. } \tag{9.0.2}
\end{equation*}
$$

The optimal value in such an inequality,

$$
c(K)=\inf _{0<t \leq \frac{1}{2}} \frac{I_{\mu}(t)}{t}
$$

was introduced in 1969 by J. Cheeger [C] and since then it is called Cheeger's isoperimetric constant. Thus, it has to be estimated from below in terms of $K$ or in terms of $\mu$ in the general measure case. In 1995, R. Kannan, L. Lovász, and
M. Simonovits [K-L-S] developed a localization technique which led to a number of interesting estimates for this quantity. In particular, they proved the bound

$$
\begin{equation*}
c(K) \geq \frac{c_{0}}{\int|x| d \mu(x)} \tag{9.0.3}
\end{equation*}
$$

with some universal constant $c_{0}$. A different approach, based on the application of Prékopa-Leindler's theorem, was later suggested in [B3]. A weak point of this bound is that the dimension $n$ is essentially involved in the right hand side of (1.3) and therefore one may not hope, for example, to reach on this way the concentration of measure phenomenon: in the simplest case of the unit Euclidean ball $K=B_{2}$, the Cheeger constant is of order $\sqrt{n}$, while the right hand side of (9.0.3) is close to a constant, as $n$ tends to infinity. The same holds for all other convex bodies $K$ which are brought in isotropic position.

Apparently this negative observation stimulated the authors of [K-L-S] to look for other better feeling the geometry of a convex body quantities. At the end of their paper, they introduced a quantity

$$
\chi(K)=\int \chi_{K}(x) d \mu(x)
$$

where $\chi_{K}(x)$ denotes the length of the longest interval in $K$ with centre at $x \in K$, that is,

$$
\begin{equation*}
\chi_{K}(x)=\operatorname{diam}((K-x) \cap(x-K)) . \tag{9.0.4}
\end{equation*}
$$

Using it they derived the following isoperimetric inequality

$$
\begin{equation*}
\mu(A)(1-\mu(A)) \leq \chi(K) \mu^{+}(A) \tag{9.0.5}
\end{equation*}
$$

which is equivalent, up to a universal factor, to the bound

$$
\begin{equation*}
c(K) \geq \frac{1}{2 \chi(K)} \tag{9.0.6}
\end{equation*}
$$

For the Euclidean balls, both sides of (9.0.6) are now of the same order. One should however emphasize that (9.0.6) remains inappropriate for some canonical cases. According to [K-L-S], (9.0.6) is not good for the regular simplex (and clearly for $\ell^{1}$-ball $K=B_{1}$, as well). Another "negative" example is given by the unit $\ell^{\infty}$ _ball $K=[-1,1]^{n}$ when the Cheeger constant $c(K)$ is of order 1 . In this case, $\chi_{K}^{2}(x)=\sum_{i}\left(1-\left|x_{i}\right|\right)^{2}$, so $\chi(K)$ is of order $\frac{1}{\sqrt{n}}$.

The question of what is really essential to effectively estimate the isoperimetric function $I_{\mu}$ and, in particular, the Cheeger constant $c(K)$ is still open. There is however a striking hypothesis suggested in the same paper [K-L-S] which says the following: if we consider the inequality (9.0.2) in the class of all half-spaces $A$ of $\mathbf{R}^{n}$, the corresponding optimal constant $c$ should be equivalent to $c(K)$ up to universal factors. The same can be conjectured for general log-concave probability measures $\mu$ on $\mathbf{R}^{n}$. For example, the hypothesis is true in the class of product log-concave measures; this was shown in [B-H1] with the help of a suitable induction argument going back to [B2].

Although, too little is known about sharp bounds in the general case, the K-L-S theorem (9.0.5)-(9.0.6) seems to be well adapted to study uniformly convex bodies and uniformly convex probability measures. The isoperimetric problem for such measures was first investigated in the 1987 paper by M. Gromov and V. D. Milman [G-M]. They introduced a special kind of localization allowing to work on
the sphere of a uniformly convex Banach space; see also [A]. For various localization arguments, we can refer the reader to the recent works $[\mathrm{F}-\mathrm{G}]$ and $[\mathrm{N}-\mathrm{S}-\mathrm{V}]$.

Let us return to the isoperimetric inequality (9.0.2) and write it as the relation

$$
I_{\mu}(t) \geq c \min \{t, 1-t\}, \quad 0<t<1
$$

How sharp is it in general ? It is well known that, in the body case, the isoperimetric function $I_{\mu}(t)$ for small $t$ resembles very much the function $c t^{(n-1) / n}$. Therefore one also considers (especially in Riemannian Geometry, cf. [Cr], [Li]) the so-called Sobolev constant

$$
s_{n}(K)=\inf _{0<t \leq \frac{1}{2}} \frac{I_{\mu}(t)}{t^{\frac{n-1}{n}}} .
$$

It turns out that this more sensitive quantity (in comparison with $c(K)$ ) can also be related to the geometric characteristic $\chi(K)$, and we have the following refinement of (9.0.6):

Theorem 9.1. For any convex body $K$ in $\mathbf{R}^{n}$,

$$
\begin{equation*}
s_{n}(K) \geq \frac{1}{C \chi(K)}, \tag{9.0.7}
\end{equation*}
$$

where $C>0$ is a numerical constant.
The proof uses a localization argument of [K-L-S] which we discuss in the next section. To be more precise, assume $n \geq 2$ and define another quantity

$$
\chi_{n}(K)=\left(\int \chi_{K}(x)^{\frac{n}{n-1}} d \mu(x)\right)^{\frac{n-1}{n}}=\left\|\chi_{K}\right\|_{L^{\frac{n}{n-1}}(K, \mu)} .
$$

With this quantity we have the following result.
Theorem 9.2. Given a convex body $K$ in $\mathbf{R}^{n}$ with the normalised Lebesgue measure $\mu$ on it, for all Borel $A \subset \mathbf{R}^{n}$,

$$
\begin{equation*}
\mu(A)^{\frac{n-1}{n}}(1-\mu(A)) \leq 4 \chi_{n}(K) \mu^{+}(A) \tag{9.0.8}
\end{equation*}
$$

In particular,

$$
\begin{equation*}
s_{n}(K) \geq \frac{1}{8 \chi_{n}(K)} \tag{9.0.9}
\end{equation*}
$$

These estimates remain to hold with respect to any norm on $\mathbf{R}^{n}$ with the corresponding notion of the diameter in (9.0.4) and of the perimeter in (9.0.1).

As we will see, in the Euclidean case, $\chi_{n}(K) \leq C \chi(K)$, for some universal constant $C$, so Theorem 9.2 is more general. We will illustrate the bound (9.0.9) on the example of uniformly convex bodies in $\mathbf{R}^{n}$.

Let us also mention a well known relationship of the Sobolev constants with Sobolev-type inequalities (which is due to V. G. Maz'ja [M1], cf. also [M2], [B-Z]). The combination of the coarea formula with (9.0.8) allows one easily to deduce:

Corollary 9.3. For any convex body $K$ in $\mathbf{R}^{n}$, and every smooth function $f$ on $\mathbf{R}^{n}$ with $\int_{K} f(x) d x=0$, with respect to any norm $\|\cdot\|$ on $\mathbf{R}^{n}$,

$$
\begin{equation*}
\left(\int_{K}|f(x)|^{\frac{n}{n-1}} d \mu(x)\right)^{\frac{n-1}{n}} \leq C \chi_{n}(K) \int_{K}\|\nabla f(x)\|_{*} d \mu(x) \tag{9.0.10}
\end{equation*}
$$

where $\|\cdot\|_{*}$ is the dual norm and $C$ is a numerical constant.

## The localization argument

The localization argument of [K-L-S] can be applied to a variety of multidimensional geometric inequalities. In particular, it can be applied to (9.0.8). With the following lemma due to L. Lovász and M. Simonovits [L-S] it can be then reduced to a certain analytic problem in dimension one.

Lemma 10.1. Let $u$ and $v$ be two lower semi-continuous, integrable functions on $\mathbf{R}^{n}$ such that

$$
\int_{\mathbf{R}^{n}} u(x) d x>0, \quad \int_{\mathbf{R}^{n}} v(x) d x>0 .
$$

Then for some points $a, b \in \mathbf{R}^{n}$ and some affine function $\ell:(0,1) \rightarrow(0,+\infty)$,

$$
\int_{0}^{1} u((1-t) a+t b) \ell(t)^{n-1} d t>0, \quad \int_{0}^{1} v((1-t) a+t b) \ell(t)^{n-1} d t>0
$$

Recall that a function is lower semi-continuous, if it can be represented as a pointwise limit of an increasing sequence of continuous functions. For example, the indicator function of an open set is lower semi-continuous. Lemma 10.1 implies:

Lemma 10.2. Let $\alpha, \beta>0$ and let $u_{1}, u_{2}, u_{3}, u_{4}$ be continuous, non-negative functions defined on a convex body $K$ in $\mathbf{R}^{n}$ and such that, for all $a, b \in K$ and for any affine function $\ell:(0,1) \rightarrow(0,+\infty)$,

$$
\left(\int_{0}^{1} u_{1}((1-t) a+t b) \ell(t)^{n-1} d t\right)^{\alpha}\left(\int_{0}^{1} u_{2}((1-t) a+t b) \ell(t)^{n-1} d t\right)^{\beta}
$$

$$
\begin{equation*}
\leq\left(\int_{0}^{1} u_{3}((1-t) a+t b) \ell(t)^{n-1} d t\right)^{\alpha}\left(\int_{0}^{1} u_{4}((1-t) a+t b) \ell(t)^{n-1} d t\right)^{\beta} \tag{10.0.1}
\end{equation*}
$$

Then,

$$
\begin{equation*}
\left(\int_{K} u_{1}(x) d x\right)^{\alpha}\left(\int_{K} u_{2}(x) d x\right)^{\beta} \leq\left(\int_{K} u_{3}(x)\right)^{\alpha}\left(\int_{K} u_{4}(x) d x\right)^{\beta} \tag{10.0.2}
\end{equation*}
$$

Lemma 10.2 is formulated in [K-L-S] in a slightly different way in terms of the so-called exponential needles. The corresponding arguments are as follows.

Without loss of generality, let the functions $u_{j}$ be strictly positive so that all the considered integrals (including one dimensional) do not vanish. By considering an $\varepsilon$-interior of $K$ and then letting $\varepsilon \downarrow 0$, we may also assume that $K$ is open and bounded.

Now, assuming the contrary to (10.0.2), one can take a number $A>0$ such that

$$
\left(\frac{\int_{K} u_{1}}{\int_{K} u_{3}}\right)^{\alpha}>A>\left(\frac{\int_{K} u_{4}}{\int_{K} u_{2}}\right)^{\beta}
$$

or equivalently,

$$
\int_{\mathbf{R}^{n}}\left(u_{1}-A^{1 / \alpha} u_{3}\right) 1_{K}>0, \quad \int_{\mathbf{R}^{n}}\left(A^{1 / \beta} u_{2}-u_{4}\right) 1_{K}>0 .
$$

The functions under the integral sign are lower semi-continuous, so, by Lemma 10.1, one can find a segment $(a, b) \subset \mathbf{R}^{n}$ and a positive affine function $\ell$ on $(a, b)$ such that

$$
\int_{(a, b)}\left(u_{1}-A^{1 / \alpha} u_{3}\right) \ell^{n-1} 1_{K}>0, \quad \int_{(a, b)}\left(A^{1 / \beta} u_{2}-u_{4}\right) \ell^{n-1} 1_{K}>0
$$

where the integrals are understood with respect to Lebesgue measure on $(a, b)$. In terms of the interval $\left(a^{\prime}, b^{\prime}\right)=(a, b) \cap K \subset K$, we have

$$
\int_{\left(a^{\prime}, b^{\prime}\right)}\left(u_{1}-A^{1 / \alpha} u_{3}\right) \ell^{n-1}>0, \quad \int_{\left(a^{\prime}, b^{\prime}\right)}\left(A^{1 / \beta} u_{2}-u_{4}\right) \ell^{n-1}>0
$$

Since again these two inequalities may be written as

$$
\left(\frac{\int_{\left(a^{\prime}, b^{\prime}\right)} u_{1} \ell^{n-1}}{\int_{\left(a^{\prime}, b^{\prime}\right)} u_{3} \ell^{n-1}}\right)^{\alpha}>A>\left(\frac{\int_{\left(a^{\prime}, b^{\prime}\right)} u_{4} \ell^{n-1}}{\int_{\left(a^{\prime}, b^{\prime}\right)} u_{2} \ell^{n-1}}\right)^{\beta}
$$

we obtain a contradiction with (10.0.1).
Clearly (using suitable approximation), Lemma 10.2 remains to hold for many discontinuous functions involving, for example, indicator functions of "regular" subsets of $\mathbf{R}^{n}$. To derive (9.0.5), Kannan, Lovász, and Simonovits considered the inequality of the form

$$
\begin{equation*}
\mu(A) \mu(B) \leq \int_{K} \chi_{K}(x) d \mu(x) \frac{\mu(C)}{h}, \quad h>0 \tag{10.0.3}
\end{equation*}
$$

where $A, B, C$ is an arbitrary partition of $K$ into non-empty "regular" subsets such that $\operatorname{dist}(A, B)=h$. Since it is of the form (10.0.2) with

$$
\alpha=\beta=1, \quad u_{1}=1_{A}, \quad u_{2}=1_{B}, \quad u_{3}=\chi_{K}, \quad u_{4}=\frac{1_{C}}{h}
$$

we are reduced to the corresponding inequality (10.0.1) for dimension one. On the other hand, fixing a set $A$, taking $B=K \backslash\left(A+h B_{2}\right), C=K \backslash(A \cup B)$, and letting $h$ tend to zero, the inequality (10.0.3) will turn unto the bound (9.0.5) for the Cheeger constant.

Similarly and better for our purposes, one may study an inequality of the form

$$
\begin{equation*}
c \mu(A)^{\frac{n-1}{n}} \mu(B) \leq\left(\int_{K} \chi_{K}(x)^{\frac{n}{n-1}} d \mu(x)\right)^{\frac{n-1}{n}} \frac{\mu(C)}{h} \tag{10.0.4}
\end{equation*}
$$

which for small $h$ becomes the desired inequality (9.0.8) of Theorem 9.2 (the constant $c>0$ has to be precised later on the basis of a certain one dimensional problem). Again, we are in position to apply Lemma 10.2 with

$$
\alpha=\frac{n-1}{n}, \beta=1, \quad u_{1}=1_{A}, \quad u_{2}=c 1_{B}, \quad u_{3}=\chi_{K}^{\frac{n}{n-1}}, \quad u_{4}=\frac{1_{C}}{h}
$$

so we are reduced to the particular case of (10.0.1) which can be written as follows

$$
\begin{equation*}
c \mu_{\ell}(A)^{\frac{n-1}{n}} \mu_{\ell}(B) \leq\left(\int_{(a, b)} \chi_{K}(x)^{\frac{n}{n-1}} d \mu_{\ell}(x)\right)^{\frac{n-1}{n}} \frac{\mu_{\ell}(C)}{h} . \tag{10.0.5}
\end{equation*}
$$

Here, $\ell$ is a positive, affine function defined on a segment $(a, b) \subset K$, and $\mu_{\ell}$ denotes the measure on $(a, b)$ with density $\ell^{n-1}$ with respect to Lebesgue measure on $(a, b)$. Recall that, (because the distance is determined by the norm $\|\cdot\|$ in $\mathbf{R}^{n}$ ), we have

$$
\begin{gathered}
h=\operatorname{dist}(A, B)=\inf \{\|a-b\|: a \in A, b \in B\} \\
\chi_{K}(x)=2 \max \left\{\|u\|: x \pm u \in K, u \in \mathbf{R}^{n}\right\}
\end{gathered}
$$

The inequality (10.0.5) will formally become stronger, if the sets $A, B, C$ are replaced respectively by $A \cap(a, b), B \cap(a, b), C \cap(a, b)$ : in this case the distance $h$ will only become larger. Therefore, in (10.0.5) it suffices to consider partitions of $(a, b)$. Moreover, the inequality will formally be strengthened if we replace $\chi_{K}$ by the smaller function

$$
\chi_{(a, b)}(x)=2 \sup \left\{\|u\|: x \pm u \in(a, b), u \in \mathbf{R}^{n}\right\}, \quad x \in(a, b)
$$

Thus, (10.0.5) is reduced to the form

$$
\begin{equation*}
c \mu_{\ell}(A)^{\frac{n-1}{n}} \mu_{\ell}(B) \leq\left(\int_{(a, b)} \chi_{(a, b)}(x)^{\frac{n}{n-1}} d \mu_{\ell}(x)\right)^{\frac{n-1}{n}} \frac{\mu_{\ell}(C)}{h} \tag{10.0.6}
\end{equation*}
$$

where $A, B, C$ is an arbitrary "regular" partition of the segment $(a, b)$ such that $h=\operatorname{dist}(A, B)$.

This problem, where the set $K$ has been eliminated, is already a purely one dimensional problem. For convenience, it may be formulated on the segment $(0,1)$ of the real line using the parametrisation of $(a, b)$ by the map $x_{t}=a+t(b-a)$, $0<t<1$. Then, $\left\|x_{t}-x_{s}\right\|=\|b-a\||t-s|$, so

$$
h=\|b-a\| \operatorname{dist}\left(A^{\prime}, B^{\prime}\right)
$$

where $A^{\prime}=\left\{t \in(0,1): x_{t} \in A\right\}, B^{\prime}=\left\{t \in(0,1): x_{t} \in B\right\}$, and where we use the usual distance on the real line. On the hand, the vector $u$ appearing in the definition of $\chi_{(a, b)}$ has to be of the form $s(b-a)$, for some $s \in \mathbf{R}$. Therefore, as easy to see,

$$
x_{t}+u \in(a, b) \text { and } x_{t}-u \in(a, b) \Longleftrightarrow|s|<\min \{t, 1-t\} .
$$

Consequently, $\chi_{(a, b)}\left(x_{t}\right)=\|b-a\| \min \{t, 1-t\}$. Introducing $C^{\prime}=\left\{t \in(0,1): x_{t} \in\right.$ $C\}$ and the measure $\mu_{\ell}^{\prime}$ on $(0,1)$ - pre-image of $\mu_{\ell}$ under the parametrisation map, we can thus rewrite (10.0.6) as

$$
c \mu_{\ell}^{\prime}\left(A^{\prime}\right)^{\frac{n-1}{n}} \mu_{\ell}^{\prime}\left(B^{\prime}\right) \leq\left(\int_{0}^{1}(\min \{t, 1-t\})^{\frac{n}{n-1}} d \mu_{\ell}^{\prime}(t)\right)^{\frac{n-1}{n}} \frac{\mu_{\ell}\left(C^{\prime}\right)}{\operatorname{dist}\left(A^{\prime}, B^{\prime}\right)}
$$

Finally, note that, up to a numerical constant, the measure $\mu_{\ell}^{\prime}$ has density $\ell\left(x_{t}\right)^{n-1}$ and that $\ell\left(x_{t}\right)$ is an affine and positive function in $t \in(0,1)$. Thus we conclude with:

Corollary 10.3. Given $c>0$ and $h>0$, the multidimensional inequality (10.0.4) is equivalent to the property that

$$
\begin{equation*}
c \mu_{\ell}(A)^{\frac{n-1}{n}} \mu_{\ell}(B) \leq\left(\int_{0}^{1}(\min \{t, 1-t\})^{\frac{n}{n-1}} d \mu_{\ell}(t)\right)^{\frac{n-1}{n}} \frac{\mu_{\ell}(C)}{h} \tag{10.0.7}
\end{equation*}
$$

where $A, B, C$ is a partition of $(0,1)$ with $\operatorname{dist}(A, B)=h$, where $\ell$ is a affine, positive function on $(0,1)$, and $\mu_{\ell}$ is a measure on $(0,1)$ with density $\ell^{n-1}$.

The above argument showed that (10.0.7) implies (10.0.4). In turn, choosing infinitesimally small truncated cones $K$ in (10.0.4), one obtains in the limit (10.0.7).

Note that the inequality (10.0.7) is homogeneous with respect to $\ell$, so we may always assume that $\mu_{\ell}$ is a probability measure, that is, $\int_{0}^{1} \ell(t)^{n-1} d t=1$.

## CHAPTER 11

## Infinitesimal version

To prove Theorem 9.2, we are now faced with the one dimensional problem (10.0.7) where finding an optimal (or some) constant $c$ does seem a much simpler, although not yet immediate, task. Recall that the inequality (10.0.4) yields the desired result by letting $h \rightarrow 0$. Similarly, we would obtain a simpler one dimensional inequality by letting $h \rightarrow 0$ in (10.0.7). This would potentially simplify our task, but then, can we return back? That is, will we obtain an equivalent inequality? An affirmative answer is given in the following:

Proposition 11.1. In the setting of an arbitrary probability metric space ( $M, \rho, \mu$ ), for any $\alpha \in\left[\frac{1}{2}, 1\right]$ and $c>0$, the inequality of the form

$$
\begin{equation*}
c \mu(A)^{\alpha} \mu(B) \leq \frac{\mu(C)}{h} \tag{11.0.1}
\end{equation*}
$$

where $A, B, C$ is an arbitrary partition of $M$ into Borel measurable subsets such that $h=\operatorname{dist}(A, B)$, is equivalent to its infinitesimal version obtained by letting $h \rightarrow 0$,

$$
\begin{equation*}
c \mu(A)^{\alpha}(1-\mu(A)) \leq \mu^{+}(A), \quad A \subset M \text { Borel. } \tag{11.0.2}
\end{equation*}
$$

In the general case, the perimeter is defined similarly to (9.0.1) as

$$
\mu^{+}(A)=\liminf _{h \downarrow 0} \frac{\mu\left(A^{h}\right)-\mu(A)}{h},
$$

where $A^{h}=\{x \in M: \rho(x, a)<h$, for some $a \in A\}$ is the open $h$-neighbourhood of $A$ with respect to $\rho$. With this notations, and since the right hand side of (11.0.1) is continuous in $h$, one can put in (11.0.1) $B=M \backslash A^{h}, C=A^{h} \backslash A$, and then the inequality itself will take the form

$$
\mu\left(A^{h}\right) \geq \frac{\mu(A)+c \cdot h \mu(A)^{\alpha}}{1+c \cdot h \mu(A)^{\alpha}}, \quad A \subset M \text { Borel. }
$$

In the sequel (throughout this section), since the parameter $c$ may be absorbed by the variable $h$, we assume $c=1$, so that (11.0.1) has the form

$$
\begin{equation*}
\mu\left(A^{h}\right) \geq L_{h}(\mu(A)), \quad h>0, \tag{11.0.3}
\end{equation*}
$$

for the family

$$
L_{h}(p)=\frac{p+h p^{\alpha}}{1+h p^{\alpha}}, \quad 0<p<1
$$

In turn, write (11.0.2) as

$$
\begin{equation*}
\mu^{+}(A) \geq I(\mu(A)), \quad 0<\mu(A)<1, \tag{11.0.4}
\end{equation*}
$$

for

$$
I(p)=p^{\alpha} q, \quad 0<p<1, q=1-p
$$

It is well known, (cf. e.g. [B-H2], [L]), that given a continuous non-negative function $I$ on ( 0,1 ), the isoperimetric inequality (11.0.4) can be "integrated over $h$ " and equivalently be written as

$$
\begin{equation*}
\mu\left(A^{h}\right) \geq R_{h}(\mu(A)), \quad h>0, \mu(A)>0 \tag{11.0.5}
\end{equation*}
$$

for a suitable semi-group $\left\{R_{h}\right\}_{h \geq 0}$. Namely, one can associate with $I$ a distribution function on the real line, $F=F(x), x \in \mathbf{R}$, defined via its inverse by

$$
F^{-1}(p)=\int_{1 / 2}^{p} \frac{d t}{I(t)}, \quad 0<p<1
$$

It has median at zero and, as a measure, $F$ is concentrated on the interval $(a, b) \subset \mathbf{R}$, finite or not, with

$$
a=-\int_{0}^{1 / 2} \frac{d t}{I(t)}, \quad b=\int_{1 / 2}^{1} \frac{d t}{I(t)}
$$

In our concrete case, $a$ is finite (if $\alpha<1$ ), while $b=+\infty$; although this will not matter later. What is important for us is that one can take in (11.0.5)

$$
\begin{equation*}
R_{h}(p)=F\left(F^{-1}(p)+h\right), \quad 0 \leq p \leq 1 \tag{11.0.6}
\end{equation*}
$$

Here, by continuity, $R_{h}(0)=F(a+h), R_{h}(1)=1$.
Recalling (11.0.3), and in view of the equivalence of (11.0.4) and (11.0.5), Proposition 11.1 will be implied by:

Proposition 11.2. For $I(p)=p^{\alpha} q$ with $\frac{1}{2} \leq \alpha \leq 1$, we have $R_{h}(p) \geq L_{h}(p)$, for all $h>0$ and $p \in[0,1]$.

For the convenience, we divide the proof into several steps. First we state one crucial property of the family $L_{h}$. Clearly, these functions bijectively act from $[0,1]$ onto itself.

Lemma 11.3. For all $h_{1}, h_{2} \geq 0$ and $p \in[0,1]$,

$$
\begin{equation*}
L_{h_{1}+h_{2}}(p) \leq L_{h_{1}}\left(L_{h_{2}}(p)\right) \tag{11.0.7}
\end{equation*}
$$

Proof. Indeed, using $L_{h}(p)=1-\frac{q}{1+h p^{\alpha}}$, the property (11.0.7) is just

$$
\frac{1-L_{h_{2}}(p)}{1+h_{1} L_{h_{2}}(p)^{\alpha}} \leq \frac{q}{1+\left(h_{1}+h_{2}\right) p^{\alpha}}
$$

that is,

$$
\left(1-L_{h_{2}}(p)\right)\left(1+\left(h_{1}+h_{2}\right) p^{\alpha}\right) \leq q\left(1+h_{1} L_{h_{2}}(p)^{\alpha}\right)
$$

This inequality is affine with respect to $h_{1}$, so the cases $h_{1}=0$ and $h_{1}=+\infty$ have to be considered, only. When $h_{1}=0$, we obtain equality (by the definition of $L_{h}$ ). Letting $h_{1} \rightarrow+\infty$, we are reduced to
$\left(1-L_{h_{2}}(p)\right) p^{\alpha} \leq q L_{h_{2}}(p)^{\alpha} \Longleftrightarrow \frac{p^{\alpha} q}{1+h_{2} p^{\alpha}} \leq q L_{h_{2}}(p)^{\alpha} \Longleftrightarrow \frac{p}{\left(1+h_{2} p^{\alpha}\right)^{1 / \alpha}} \leq L_{h_{2}}(p)$.
Replacing $h_{2}$ by $h$ and using the definition of $L_{h}$, we arrive at

$$
\begin{equation*}
p\left(1+h p^{\alpha}\right)^{\frac{1}{\alpha}-1} \leq p+h p^{\alpha} \tag{11.0.8}
\end{equation*}
$$

Note that, due to the assumption $\alpha \in\left[\frac{1}{2}, 1\right]$, the left hand side represents a nondecreasing, concave function in $h \geq 0$, while the right hand side is affine. Hence, it suffices to consider the behaviour of both sides near zero. When $h=0$, (11.0.8) turns into equality, while the corresponding derivatives at $h=0$ turn the inequality into $p\left(\frac{1}{\alpha}-1\right) \leq 1$ which is true. Thus, the property (11.0.7) is verified.

In contrast with (11.0.7), it follows immediately from the definition (11.0.6) that:

Lemma 11.4. For all $h_{1}, h_{2} \geq 0$ and $p \in[0,1], R_{h_{1}+h_{2}}(p)=R_{h_{1}}\left(R_{h_{2}}(p)\right)$.
Lemma 11.5. For $I(p)=p^{\alpha} q$ with $\frac{1}{2} \leq \alpha \leq 1$, there exist $0<p_{0}<p_{1}<1$ and $h_{0}>0$ such that

$$
\begin{equation*}
R_{h}(p) \geq L_{h}(p) \tag{11.0.9}
\end{equation*}
$$

whenever $h \in\left[0, h_{0}\right]$ and $p \in\left[0, p_{0}\right] \cup\left[p_{1}, 1\right]$.
Proof. We may assume $\frac{1}{2}<\alpha<1$ to simplify the argument in some details. The function $I(p)=p^{\alpha} q$ is concave on ( 0,1 ), so $F$ has log-concave density $f(x)=$ $F^{\prime}(x)$ on the supporting interval $(a, b)$. Recall that, according to the definition of $F$,

$$
f\left(F^{-1}(p)\right)=I(p)=p^{\alpha} q, \quad 0<p<1 .
$$

Clearly, the density $f$ is smooth on $(a, b)$. Formally differentiating the above equality, we get

$$
\begin{equation*}
f^{\prime}\left(F^{-1}(p)\right)=I(p) I^{\prime}(p)=p^{2 \alpha-1} q(\alpha q-p), \quad 0<p<1 . \tag{11.0.10}
\end{equation*}
$$

Further differentiation of $J(p)=f^{\prime}\left(F^{-1}(p)\right)$ gives

$$
J^{\prime}(p)=p^{2 \alpha-2}\left((\alpha+1)(2 \alpha+1) p^{2}-2 \alpha(2 \alpha+1) p+\alpha(2 \alpha-1)\right) .
$$

Thus, for small $p, J^{\prime}(p)$ behaves like $\alpha(2 \alpha-1) p^{2 \alpha-2}>0$, so

$$
J^{\prime}(0+)=+\infty, \quad J^{\prime}(1-)=1
$$

Hence, $J(p)$ increases near $p=0$ and near $p=1$, and the same is true for the function $f^{\prime}(x)$ near $x=a$ and $x=+\infty$.

For a fixed $p \in(0,1)$, write Taylor's expansion for $R_{h}(p)=F\left(F^{-1}(p)+h\right)$ as a function of $h$ up to the quadratic term:

$$
R_{h}(p)=p+I(p) h+f^{\prime}\left(F^{-1}(p)+w\right) \frac{h^{2}}{2}, \quad 0<w<h
$$

If $p$ is sufficiently close to 1 , or if $p$ and $h$ are sufficiently close to 0 , by the above mentioned monotonicity,

$$
R_{h}(p) \geq p+I(p) h+f^{\prime}\left(F^{-1}(p)\right) \frac{h^{2}}{2}
$$

Therefore, by (11.0.10) and since $L_{h}(p)=1-\frac{q}{1+h p^{\alpha}},(11.0 .9)$ would be implied by

$$
p+I(p) h+I(p) I^{\prime}(p) \frac{h^{2}}{2} \geq 1-\frac{q}{1+h p^{\alpha}} .
$$

Subtracting $p$ and dividing by $I(p) h$, this inequality becomes $1+I^{\prime}(p) \frac{h}{2} \geq \frac{1}{1+h p^{\alpha}}$, or, $-\frac{1}{2} I^{\prime}(p) \leq \frac{p^{\alpha}}{1+h p^{\alpha}}$. For small $p$, this bound is immediate, since $I$ increases near zero. To treat the other case, rewrite the bound as $\frac{1}{2}(p-\alpha q) \leq \frac{p}{1+h p^{\alpha}}$. This
inequality is readily fulfilled for sufficiently small $q$ and $h$ (independently of each other). Hence Lemma 11.5 follows.

Proof of Proposition 11.2. Given $t \in(0,1)$, define

$$
\Delta=\left\{h \geq 0: R_{h}(p) \geq L_{t h}(p), \text { for all } p \in[0,1]\right\}
$$

If we show that $\Delta=[0,+\infty)$, then letting $t \rightarrow 1$ will give the desired inequality (11.0.9). For any $h_{1}, h_{2} \in \Delta$ and $p \in[0,1]$, by Lemmas 11.3 and 11.4,

$$
R_{h_{1}+h_{2}}(p)=R_{h_{1}}\left(R_{h_{2}}(p)\right) \geq L_{t h_{1}}\left(L_{t h_{2}}(p)\right) \geq L_{t\left(h_{1}+h_{2}\right)}(p)
$$

Hence, $h_{1}+h_{2} \in \Delta$, so $\Delta+\Delta=\Delta$. Therefore, it suffices to show that $\Delta$ contains all sufficiently small $h>0$. By Lemma 11.5, in the inequality $R_{h}(p) \geq L_{t h}(p)$, it remains to consider the values $p_{0} \leq p \leq p_{1}$, only. Writing the Taylor's expansion over $h \downarrow 0$ for the functions $R_{h}(p)$ and $L_{t h}(p)$ we have:

$$
\begin{equation*}
R_{h}(p)=p+I(p) h+O\left(h^{2}\right), \quad L_{h}(p)=p+t I(p) h+O\left(h^{2}\right) \tag{11.0.11}
\end{equation*}
$$

Here both constants in $O\left(h^{2}\right)$ are uniform in $p$, since the corresponding second derivatives are bounded. Indeed, by (11.0.10), $f^{\prime}$ is bounded, so is $\frac{d^{2} R_{h}(p)}{d h^{2}}=$ $f^{\prime}\left(F^{-1}(p)+h\right)$. In addition,

$$
\frac{d^{2} L_{h}(p)}{d h^{2}}=-\frac{2 p^{2 \alpha} q}{\left(1+h p^{\alpha}\right)^{3}}
$$

is evidently bounded. But $I(p)$ is separated from zero on $\left[p_{0} . p_{1}\right]$, so, for small $h>0$, the inequality $R_{h}(p) \geq L_{t h}(p)$ immediately follows from (11.0.11).

## CHAPTER 12

## Proof of Theorem 9.2

Now we can return to the one dimensional inequality (10.0.7) and apply Proposition 11.1 to rewrite (10.0.7) in the infinitesimal form as

$$
\begin{equation*}
c \mu_{\ell}(A)^{\alpha}\left(1-\mu_{\ell}(A)\right) \leq\left(\int_{0}^{1}(\min \{t, 1-t\})^{1 / \alpha} d \mu_{\ell}(t)\right)^{\alpha} \mu_{\ell}^{+}(A) . \tag{12.0.1}
\end{equation*}
$$

Here $A$ is an arbitrary Borel subset of $(0,1)$ and $\alpha=\frac{n-1}{n}$. Note that $\frac{1}{2} \leq \alpha<1$, for any integer $n \geq 2$. Since $L^{1 / \alpha}$-norm majorizes $L^{1}$-norm, (12.1) can be strengthened as

$$
\begin{equation*}
c \mu_{\ell}(A)^{\alpha}\left(1-\mu_{\ell}(A)\right) \leq \int_{0}^{1} \min \{t, 1-t\} d \mu_{\ell}(t) \mu_{\ell}^{+}(A) \tag{12.0.2}
\end{equation*}
$$

(remaining equivalent up to a numerical factor). Note that, in terms of the distribution function $F(t)=\mu_{\ell}[0, t], 0 \leq t \leq 1$, the above integral has a simple representation

$$
\begin{equation*}
\int_{0}^{1} \min \{t, 1-t\} d \mu_{\ell}(t)=\int_{0}^{1 / 2}\left(F\left(t+\frac{1}{2}\right)-F(t)\right) d t \tag{12.0.3}
\end{equation*}
$$

Thus, we are dealing in (12.0.2) with an inequality on the real line of the form

$$
\begin{equation*}
\bar{c} \lambda(A)^{\alpha}(1-\lambda(A)) \leq \lambda^{+}(A) \tag{12.0.4}
\end{equation*}
$$

for special probability measures $\lambda=\mu_{\ell}$ (the new constant $\bar{c}$ temporarily absorbs the other factor in (12.0.2) independent of $A$ ). Since the non-negative functions of the form $\ell^{n-1}$ on $[0,1]$ extended to the whole line by zero outside $[0,1]$ are log-concave, the measures $\mu_{\ell}$ are log-concave in the sense of Prékopa (cf. [Pr], [Bor2]). The isoperimetric problem for any log-concave measure $\lambda$ on $R$ has a simple solution: the perimeter $\lambda^{+}(A)$ attains minimum within all Borel measurable sets $A$ with prescribed value of $\lambda(A)$ either for a half-axis $A=(-\infty, x]$, or for $A=[x,+\infty)$; see [B1] for details. Equivalently, in terms of the distribution function $F(x)=\lambda(-\infty, x]$ and its density $f$, we have a simple representation for the isoperimetric function of such a measure,

$$
\begin{equation*}
I_{\lambda}(p)=\min \left\{f\left(F^{-1}(p)\right), f\left(F^{-1}(q)\right)\right\}, \quad 0<p<1, q=1-p, \tag{12.0.5}
\end{equation*}
$$

where $F^{-1}$ is the corresponding inverse function. Therefore, the inequality (12.0.4) is needed to verify on half-axes, only, and the optimal value of $\bar{c}$ comes from

$$
\begin{equation*}
\bar{c} p^{\alpha} q \leq I_{\lambda}(p) \tag{12.0.6}
\end{equation*}
$$

with the right hand side described in (12.0.5).
So, let us find an analytic expression for the isoperimetric function in case $\lambda=\mu_{\ell}$. Since $I_{\lambda}(p)$ is always symmetric about the point $p=\frac{1}{2}$, and since the
integral in (12.0.2) is invariant under the transformation $t \rightarrow 1-t$, we may assume that $\ell(t)$ is non-decreasing in $t$ and thus has the form

$$
\ell(t)=\operatorname{const}(r+t)
$$

with necessarily $r \geq 0$. In view of the condition $\mu_{\ell}[0,1]=1$, the distribution and density functions of this measure are given by

$$
\begin{equation*}
F_{r}(t)=\frac{(r+t)^{n}-r^{n}}{(r+1)^{n}-r^{n}}, \quad f_{r}(t)=\frac{n}{(r+1)^{n}-r^{n}}(r+t)^{n-1} \quad(0<t<1), \tag{12.0.7}
\end{equation*}
$$

where $r$ is a non-negative parameter. From this, as it is easy to see,

$$
f_{r}\left(F_{r}^{-1}(p)\right)=\frac{n}{(r+1)^{n}-r^{n}}\left(p(r+1)^{n}+q r^{n}\right)^{\alpha} .
$$

Given $x, y \geq 0$, the quantity $\frac{(x p+y q)^{\alpha}}{p^{\alpha} q}=\left(x+\frac{y}{q}\right)^{\alpha} \frac{1}{q}$ attains minimum for $q=1$, when it equals $x+y$. In particular,

$$
\inf _{0<p<1} \frac{\left(p(r+1)^{n}+q r^{n}\right)^{\alpha}}{p^{\alpha} q}=\inf _{0<p<1} \frac{\left(q(r+1)^{n}+p r^{n}\right)^{\alpha}}{p^{\alpha} q}=\left((r+1)^{n}+r^{n}\right)^{\alpha}
$$

Hence, by (12.0.5), the optimal constant in (12.0.6) for the isoperimetric function $I_{\lambda}$ of $\mu_{\ell}$ satisfies

$$
\begin{equation*}
\bar{c}=n \frac{\left((r+1)^{n}+r^{n}\right)^{\alpha}}{(r+1)^{n}-r^{n}} \geq n \frac{(r+1)^{n-1}}{(r+1)^{n}-r^{n}} . \tag{12.0.8}
\end{equation*}
$$

Now, according to (12.0.7),

$$
\begin{equation*}
\int_{0}^{1 / 2}\left(F_{r}\left(t+\frac{1}{2}\right)-F_{r}(t)\right) d t=\frac{1}{n+1} \frac{(r+1)^{n+1}-2\left(r+\frac{1}{2}\right)^{n+1}+r^{n+1}}{(r+1)^{n}-r^{n}} \tag{12.0.9}
\end{equation*}
$$

Recalling (12.0.2)-(12.0.3) and applying (12.0.8)-(12.0.9), we may summarize:
Lemma 12.1. The inequality (12.0.1) for the measure $\mu_{\ell}$ defined in (12.0.7) with parameter $r \geq 0$ holds true with

$$
\begin{equation*}
c=\frac{n}{n+1} \frac{(r+1)^{n-1}\left((r+1)^{n+1}-2\left(r+\frac{1}{2}\right)^{n+1}+r^{n+1}\right)}{\left((r+1)^{n}-r^{n}\right)^{2}} . \tag{12.0.10}
\end{equation*}
$$

At last, we are ready to complete the proof of Theorem 9.2.
Proof of Theorem 9.2. It remains to estimate from below by a numerical constant the right hand side of (12.0.10). So, let us look at the inequality of the form

$$
\begin{equation*}
(r+1)^{n-1}\left((r+1)^{n+1}-2\left(r+\frac{1}{2}\right)^{n+1}+r^{n+1}\right) \geq D_{n}\left((r+1)^{n}-r^{n}\right)^{2} \tag{12.0.11}
\end{equation*}
$$

Note that both sides represent certain polynomials of degree $2(n-1)$. More precisely, according to Newton's binomial formula,

$$
\left((r+1)^{n}-r^{n}\right)^{2}=\sum_{i=1}^{n} \sum_{j=1}^{n} C_{n}^{i} C_{n}^{j} r^{2 n-(i+j)},
$$

where $C_{n}^{i}=\frac{n!}{i!(n-i)!}$ are usual binomial coefficients. Writing $(r+1)^{n-1}=\sum_{i=1}^{n} C_{n-1}^{i-1} r^{n-i}$ and noting that
$(r+1)^{n+1}-2\left(r+\frac{1}{2}\right)^{n+1}+r^{n+1}=\sum_{j=1}^{n} C_{n+1}^{j+1}\left(1-\frac{1}{2^{j}}\right) r^{n-j} \geq \frac{1}{2} \sum_{j=1}^{n} C_{n+1}^{j+1} r^{n-j}$.
we see that (12.0.11) is implied by

$$
\begin{equation*}
\sum_{i=1}^{n} \sum_{j=1}^{n} C_{n-1}^{i-1} C_{n+1}^{j+1} r^{2 n-(i+j)} \geq 2 D_{n}\left(\sum_{i=1}^{n} \sum_{j=1}^{n} C_{n}^{i} C_{n}^{j} r^{2 n-(i+j)}\right) \tag{12.0.12}
\end{equation*}
$$

Let us compare here the coefficients $A_{i j}=C_{n-1}^{i-1} C_{n+1}^{j+1}$ with $B_{i j}=C_{n}^{i} C_{n}^{j}$. We have in general $A_{i j}=\frac{i}{j+1} \frac{n+1}{n} B_{i j}$, so

$$
A_{i i}=\frac{i}{i+1} \frac{n+1}{n} B_{i i} \geq \frac{1}{2} \frac{n+1}{n} B_{i i}, \quad 1 \leq i \leq n
$$

In view of $B_{i j}=B_{j i}$, for the case $1 \leq i \neq j \leq n$ we also have

$$
A_{i j}+A_{j i}=\left(\frac{i}{j+1}+\frac{j}{i+1}\right) \frac{n+1}{n} \frac{B_{i j}+B_{j i}}{2} \geq \frac{1}{2} \frac{n+1}{n}\left(B_{i j}+B_{j i}\right)
$$

because one of the fractions $\frac{i}{j+1}, \frac{j}{i+1}$ is greater or equal to 1 . This shows that (12.0.12) holds true with $2 D_{n}=\frac{1}{2} \frac{n+1}{n}$. Therefore, by Lemma 12.1, we may take $c=\frac{1}{4}$ in (12.0.1) and Theorem 9.2 is now proved.

## CHAPTER 13

## Euclidean distance (proof of Theorem 9.1)

Since $\frac{n}{n-1} \leq 2$, in order to derive Theorem 9.1 from Theorem 9.2 , it will be sufficient to establish a Khichine-type inequality for the functions $\chi_{K}$. As we will see, it holds true with respect to any log-concave measure.

Proposition 13.1. For any convex body $K$ in $\mathbf{R}^{n}$ and every log-concave probability measure $\lambda$ concentrated on $K$,

$$
\begin{equation*}
\left(\int \chi_{K}^{2} d \lambda\right)^{1 / 2} \leq C \int \chi_{K} d \lambda \tag{13.0.1}
\end{equation*}
$$

where $C$ is a numerical constant.
Proof. Log-concavity of a measure $\lambda$ on $\mathbf{R}^{n}$ means (cf. [Bor2]) that, for all $t, s>0$ with $t+s=1$ and for all non-empty Borel subsets $A, B, C$ of $R^{n}$ with $t A+s B \subset C$,

$$
\lambda(C) \geq \lambda(A)^{t} \lambda(B)^{s}
$$

The uniform distribution $\lambda=\mu$ on $K$ possesses this property due to BrunnMinkowski inequality (cf. e.g. [B-Z]).

By (9.0.4) for the distance given by the Euclidean norm $|\bullet|$ in $\mathbf{R}^{n}$, for any $x \in K$ one has

$$
\chi_{K}(x)=2 \max \left\{|u|: x \pm u \in K, u \in \mathbf{R}^{n}\right\} .
$$

For given $x, y \in K$, let the vectors $u, v \in \mathbf{R}^{n}$ be the maximizers in this definition. Since for all $t, s>0$ with $t+s=1$, necessarily $t x+s y \pm(t u \pm s v) \in K$, we get

$$
\chi_{K}^{2}(t x+s y) \geq 4|t u \pm s v|^{2}=4\left(t^{2}|u|^{2} \pm 2 t s\langle u, v\rangle+s^{2}|v|^{2}\right) .
$$

Choosing appropriately the sign on the right hand side, we obtain

$$
\begin{equation*}
\chi_{K}^{2}(t x+s y) \geq t^{2} \chi_{K}^{2}(x)+s^{2} \chi_{K}^{2}(y), \quad x, y \in K, t, s>0, t+s=1 . \tag{13.0.2}
\end{equation*}
$$

Let $D=\max _{x \in K} \chi_{K}^{2}(x)$ (clearly, the function $\chi_{K}$ is continuous, so the maximum is attained at some point), and let $\xi=\chi_{K}^{2}$ be considered as a random variable on the probability space $(K, \lambda)$ so that

$$
\begin{aligned}
\int \chi_{K}^{2} d \lambda & =\mathbf{E} \xi=\int_{0}^{D} \lambda\{\xi \geq a\} d a \\
\int \chi_{K} d \lambda & =\mathbf{E} \xi^{1 / 2}=\int_{0}^{D} \lambda\{\xi \geq a\} d a^{1 / 2}
\end{aligned}
$$

For any $a \in[0, D]$, define the set $A(a)=\{x \in K: \xi(x) \geq a\}$. By (13.0.2), for all $a, b \in[0, D]$, we have $t A(a)+s A(b) \subset A\left(t^{2} a+s^{2} b\right)$. Hence, by the log-concavity of $\lambda$, the function under the integral sign, $f(a)=\lambda(A(a))$, possesses the property

$$
\begin{equation*}
f\left(t^{2} a+s^{2} b\right) \geq f(a)^{t} f(b)^{s}, \quad 0 \leq a, b \leq D, t, s>0, t+s=1 \tag{13.0.3}
\end{equation*}
$$

The inequality remains to hold if we extend $f$ by zero to the interval $(D,+\infty)$. Note also that $f(0)=1$ and that $f$ will be non-increasing and left continuous on $[0,+\infty)$. The last mentioned properties ensure that there exists a maximal number $m \in[0 . D]$ such that $f(m) \geq \frac{1}{e}$. Hence,

$$
\mathbf{E} \xi^{1 / 2} \geq \int_{0}^{m} f(a) d a^{1 / 2} \geq \frac{1}{e} m^{1 / 2}
$$

and

$$
\begin{equation*}
m \leq e^{2}\left(\mathbf{E} \xi^{1 / 2}\right)^{2} \tag{13.0.4}
\end{equation*}
$$

On the other hand, let us apply (13.0.3) with $b=0, a=\frac{a_{0}}{t^{2}}, r=\frac{1}{t^{2}}$ where $a_{0}>m$ is a fixed parameter. Then we get

$$
f\left(a_{0} r\right) \leq f\left(a_{0}\right)^{\sqrt{r}} \leq e^{-\sqrt{r}}, \quad r>1,
$$

so

$$
\mathbf{E} \xi=\int_{0}^{a_{0}} f(a) d a+a_{0} \int_{1}^{+\infty} f\left(a_{0} r\right) d r \leq a_{0}+a_{0} \int_{1}^{+\infty} e^{-\sqrt{r}} d r=\left(1+\frac{4}{e}\right) a_{0}
$$

Letting $a_{0} \downarrow m$, we get $\mathbf{E} \xi \leq\left(1+\frac{4}{e}\right) m$ which together with (13.0.4) yields

$$
\mathbf{E} \xi \leq\left(e^{2}+4 e\right)\left(\mathbf{E} \xi^{1 / 2}\right)^{2}
$$

This is just the desired estimate (13.0.1) with $C=\sqrt{e^{2}+4 e}$.
Remark 13.2. The property (13.0.2) is very close to the description of $\chi_{K}^{2}$ as a concave function. For example, in case of the unit Euclidean ball $K=B_{2}$, the function $\chi_{K}^{2}(x)=1-|x|^{2}$ is indeed concave. Khinchine-type inequalities for such functions were investigated by L. Berwald who showed (cf. [Ber] and [Bor1] for another proof), that for all $p>q>0$, within the class of all convex bodies $K$ in $\mathbf{R}^{n}$ and all concave functions $g$ on $K$, the ratio

$$
\frac{\|g\|_{L^{p}(K, \mu)}}{\|g\|_{L^{q}(K, \mu)}}
$$

is maximized in case of the simplex $K=\left\{x \in \mathbf{R}_{+}^{n}: x_{1}+\cdots+x_{n} \leq 1\right\}$ and the linear function $g(x)=x_{1}$. Here, $\mu$ denotes the uniform measure on $K$. In particular, this result implies a Khinchine-type inequality

$$
\begin{equation*}
\|g\|_{L^{2}(K, \mu)} \leq \sqrt{2}\|g\|_{L^{1}(K, \mu)} \tag{13.0.5}
\end{equation*}
$$

## CHAPTER 14

## Uniformly convex bodies

Let $K$ be a symmetric about the origin, convex body in $\mathbf{R}^{n}$. We are going to examine the bound (9.0.9) for the Sobolev constant $s_{n}(K)$ with respect to the canonical metric generated by the norm

$$
\|x\|=\min \{\lambda \geq 0: x \in \lambda K\}, \quad x \in \mathbf{R}^{n}
$$

Being restricted to $K$, this metric is called the inner metric for $K$.
Note that in general, for any $x \in K$, the assumptions $x \pm u \in K$ imply $u \in K$, so $\chi_{K}(x) \leq 1$. Hence, by Theorem 9.2,

Proposition 14.1. The Sobolev constant $s_{n}(K)$ defined with respect to the inner metric on $K$ satisfies

$$
\begin{equation*}
s_{n}(K) \geq c . \tag{14.0.1}
\end{equation*}
$$

with some universal constant $c \geq \frac{1}{2}$ independent of $K$.
Actually, the bound (9.0.9) yields $c=\frac{1}{8}$, and the improvement can be reached using a different argument. Recall the classical Brunn-Minkowski inequality (cf. e.g. [B-Z]): for all $t \in(0,1)$ and all non-empty Borel sets $A, B \subset \mathbf{R}^{n}$,

$$
\operatorname{vol}_{n}(t A+(1-t) B) \geq\left(t \operatorname{vol}_{n}(A)^{1 / n}+(1-t) \operatorname{vol}_{n}(B)^{1 / n}\right)^{n}
$$

Hence, the uniform probability measure $\mu$ on $K$ possesses the same property:

$$
\mu(t A+(1-t) B) \geq\left(t \mu(A)^{1 / n}+(1-t) \mu(B)^{1 / n}\right)^{n}
$$

Applying it to the couple $(A, K)$ and to the couple $(K \backslash A, K)$, we get

$$
\begin{gathered}
\mu(t A+(1-t) K) \geq\left(t p^{1 / n}+1-t\right)^{n} \\
\mu(t(K \backslash A)+(1-t) K) \geq\left(t q^{1 / n}+1-t\right)^{n}
\end{gathered}
$$

where $p=\mu(A), q=1-p$. Adding the last two inequalities and letting $t \rightarrow 1$, in the limit we obtain:

Proposition 14.2. For every Borel subset $A$ of $K$ of measure $\mu(A)=p$,

$$
\begin{equation*}
\mu^{+}(A) \geq \frac{n}{2}\left(p^{\frac{n-1}{n}}+q^{\frac{n-1}{n}}-1\right), \quad q=1-p \tag{14.0.2}
\end{equation*}
$$

Here, the perimeter is understood as in (9.0.1) with $B_{2}$ replaced by $K$. To be more precise, in the above argument one has to assume that $A$ is smooth enough, and then one may easily extend (14.0.2) from smooth sets to all Borel measurable. In a slightly different form, the argument appeared in [B3]; the present formulation with similar proof is given by F. Barthe [Bar].

Letting $n \rightarrow \infty$ in (14.0.2), we obtain a dimension free bound

$$
\mu^{+}(A) \geq \frac{1}{2}\left(p \log \frac{1}{p}+q \log \frac{1}{q}\right), \quad q=1-p
$$

Actually, this inequality remains to hold for any log-concave probability measure concentrated on $K([\mathrm{~B} 3])$. Clearly, it implies that the Cheeger constant $c(K)$ is separated from zero; nevertheless, (14.0.2) contains more information concerning the behaviour of the isoperimetric function $I_{\mu}(p)$ near zero. For example, we have:

Corollary 14.3. For any $n \in \mathbf{N}$

$$
\frac{n}{2} \leq \liminf _{p \rightarrow 0} \frac{I_{\mu}(p)}{p^{(n-1) / n}} \leq \limsup _{p \rightarrow 0} \frac{I_{\mu}(p)}{p^{(n-1) / n}} \leq n
$$

Here, the first inequality immediately follows from (14.0.2). In order to see the third one, it suffices to consider the set $A=p^{1 / n} K$ in which case $\mu^{+}(A)=$ $n p^{(n-1) / n}$. Hence, $I_{\mu}(p) \leq n p^{(n-1) / n}$, for all $p \in(0,1)$.

Proof of Proposition 14.1. We need to check that (14.0.2) implies the inequality (14.0.1) with constant $c=\frac{1}{2}$. Let $\alpha=\frac{n-1}{n}$ so that $n=\frac{1}{1-\alpha}$. Rewrite the (desired) inequality $\frac{1}{2(1-\alpha)}\left(p^{\alpha}+q^{\alpha}-1\right) \geq c p^{\alpha}$ as

$$
\begin{equation*}
\frac{1}{2(1-\alpha)}\left(1-\frac{1-q^{\alpha}}{p^{\alpha}}\right) \geq c, \quad 0<p \leq \frac{1}{2} \tag{14.0.3}
\end{equation*}
$$

By direct differentiation, one sees that the function $p \rightarrow \frac{1-q^{\alpha}}{p^{\alpha}}$ is increasing on $(0,1)$. Hence, the left hand side of (14.0.3) is minimized as $p=\frac{1}{2}$, and we are reduced to $a_{n} \equiv n\left(1-2^{-1 / n}\right) \geq c$. The sequence $a_{n}$ is increasing, so $c=a_{1}=\frac{1}{2}$.

Proposition 14.1 is proved.
As turns out, in spaces of large dimension, the Sobolev constant $s_{n}(K)$ can be much larger, and the inequality (14.0.1) can be considerably sharpened for "sufficiently" convex bodies $K$.

A convex, symmetric about the origin body $K$ in $\mathbf{R}^{n}$ is called uniformly convex, if the corresponding modulus of convexity,

$$
\delta_{K}(\varepsilon)=\inf \left\{1-\left\|\frac{x+y}{2}\right\|:\|x\|=\|y\|=1,\|x-y\| \leq \varepsilon\right\}
$$

is positive for each $\varepsilon>0$. Here as before, the norm in $\mathbf{R}^{n}$ is inner for $K$. The Banach space ( $\mathbf{R}^{n},\|\cdot\|$ ) is said to have a modulus of convexity of power $p \geq 2$ with constant $C>0$, if $\delta_{K}(\varepsilon) \geq C \varepsilon^{p}$, for every $\varepsilon \in(0,2]$. Equivalently (cf. e.g. [P]), for all $x, y \in \mathbf{R}^{n}$,

$$
\begin{equation*}
\frac{\|x\|^{p}+\|y\|^{p}}{2}-\left\|\frac{x+y}{2}\right\|^{p} \geq \kappa^{p}\|x-y\|^{p} \tag{14.0.4}
\end{equation*}
$$

for some $\kappa>0$. For example (cf. $[\mathrm{L}-\mathrm{T}]$ ), the unit $\ell^{p}$-ball $K=B_{p}$ with $p \geq 2$ satisfies (14.0.4) with optimal $\kappa=\frac{1}{2}$.

Under the assumption (14.0.4), the characteristic $\chi_{n}(K)$ may easily be estimated in terms of the constant $\kappa$. Indeed, for $x \in K$, let $x \pm u \in K$ with maximal possible value of $\|u\|$ (equal to $\frac{1}{2} \chi_{K}(x)$ ). Since $\|x \pm u\| \leq 1$, application of (14.0.4)
to $x+u$ and $x-u$ (in the place of $x$ and $y$, respectively) gives $1-\|x\|^{p} \geq \kappa^{p} 2^{p}\|u\|^{p}$. Consequently, we get:

Lemma 14.4. Under (14.0.4), for all $x \in K$,

$$
\begin{equation*}
\chi_{K}(x) \leq \frac{1}{\kappa}\left(1-\|x\|^{p}\right)^{1 / p} \tag{14.0.5}
\end{equation*}
$$

To apply (9.0.9), we may use the above bound which leads to computing the norm $\|g\|_{L^{\frac{n}{n-1}}(K, \mu)}$ for the function $g(x)=\left(1-\|x\|^{p}\right)^{1 / p}$. To a little simplify the task, let us note that $g$ is concave on $K$ and apply Berwald's inequality (13.0.5). With (14.0.5), we then get

$$
\begin{equation*}
\chi_{n}(K) \leq \frac{1}{\kappa}\|g\|_{L^{\frac{n}{n-1}}(K, \mu)} \leq \frac{\sqrt{2}}{\kappa}\|g\|_{L^{1}(K, \mu)} . \tag{14.0.6}
\end{equation*}
$$

Now,

$$
\int_{K}\left(1-\|x\|^{p}\right)^{1 / p} d \mu(x)=\int_{0}^{1}\left(1-t^{p}\right)^{1 / p} d t^{n}=\frac{\Gamma\left(\frac{n}{p}+1\right) \Gamma\left(\frac{1}{p}+1\right)}{\Gamma\left(\frac{n+1}{p}+1\right)}
$$

Therefore, by (14.0.6) and Theorem 9.2,

$$
\begin{equation*}
s_{n}(K) \geq \frac{\kappa}{8 \sqrt{2}} \frac{\Gamma\left(\frac{n+1}{p}+1\right)}{\Gamma\left(\frac{n}{p}+1\right) \Gamma\left(\frac{1}{p}+1\right)} . \tag{14.0.7}
\end{equation*}
$$

One may also apply Stirling's formula, to develop the right hand side of (14.0.7):
Proposition 14.6. Let $K$ be a symmetric convex body in $\mathbf{R}^{n}$ satisfying the property (14.0.4) with parameters $p \geq 2$ and $\kappa>0$. Then, with respect to the inner metric in $K$, the Sobolev constant satisfies

$$
\begin{equation*}
s_{n}(K) \geq c_{n, p} n^{1 / p} \kappa, \tag{14.0.8}
\end{equation*}
$$

for some double indexed sequence such that $c_{n, p} \rightarrow c$, as $n \rightarrow \infty$ and $\frac{p}{n} \rightarrow 0$ (where $c$ is a positive universal constant).

In particular, for unit $\ell^{p}$-balls $K=B_{p}$ with $p \geq 2$, we have $s_{n}\left(B_{p}\right) \geq c_{n, p} n^{1 / p}$.

## CHAPTER 15

## From Isoperimetry to $\mathrm{LS}_{q}$-inequalities

To illustrate what can happen for $l^{\infty}$ ball (when Dirichlet boundary are involved), we utilise an idea of $[\mathbf{Z 3}]$ as follows. First we recall that for the $n$th Cartesian product of the unit intervals one has the following Classical Sobolev inequality with Dirichlet boundary conditions for the product measure $\nu_{n}$ on $[0,1]^{n}$.

$$
\begin{equation*}
\left(\int|f|^{\frac{n}{n-1}} d \nu_{n}\right)^{\frac{n-1}{n}} \leq \frac{1}{n} \sum_{i=1}^{n} \int\left|\nabla_{i} f\right| d \nu_{n} \tag{15.0.1}
\end{equation*}
$$

Using (15.0.1), for any $q \in(1, \infty)$ with the help of Holder inequality, we obtain

$$
\begin{equation*}
\left(\int|f|^{\frac{n}{n-1}} d \nu_{n}\right)^{\frac{n-1}{n}} \leq \frac{1}{n^{1 / q}} \int|\nabla f|_{q} d \nu_{n} \tag{15.0.2}
\end{equation*}
$$

where $|\nabla f|_{q}$ denotes the $l_{q}$ norm of the gradient. We note that if $f$ is nonnegative, with $p \in(1, \infty)$ such that $\frac{1}{p}+\frac{1}{q}=1$, using the differentialtion rules and the Young inequality we get

$$
\frac{1}{n^{1 / q}}|\nabla f|_{q}=p^{\frac{1}{p}} f^{\frac{1}{p}} \cdot \frac{1}{n^{1 / q}} q p^{-\frac{1}{p}}\left|\nabla f^{\frac{1}{q}}\right|_{q} \leq f+\frac{1}{n}(q-1)^{q-1}|\nabla f|_{q}^{q}
$$

Inserting this into (15.0.2) and performing simple algebraic transformations, we obtain the following relation

$$
\frac{\|f\|_{\frac{n}{n-1}}-\|f\|_{1}}{\frac{1}{n}} \leq(q-1)^{q-1} \int\left|\nabla f^{\frac{1}{q}}\right|_{q}^{q} d \nu_{n}
$$

Hence passing to the limit with $n \longrightarrow \infty$ and replacing $f$ by $f^{q}$, we arrive at the following $\mathrm{LS}_{q}$-inequality

$$
\operatorname{Ent}_{\nu}\left(f^{q}\right) \leq(q-1)^{q-1} \int|\nabla f|_{q}^{q} d \nu
$$

with the infinte dimensional product measure $\nu$ on $[0,1]^{\mathbf{N}}$. (Note that the limiting procedure involves in fact also a sequence of functions, as we need to preserve the Dirichlet condition all time on the way.)

## CHAPTER 16

## Isoperimetric Functional Inequalities

The following functional inequality was first introduced in $[\mathbf{B 2}]$ for the two point measure and via central limit theorem extended to the standard Gaussian measure
$\left(\mathrm{IFI}_{2}\right)$

$$
\mathcal{I}(\mu f) \leq \mu\left(\mathcal{I}(f)^{2}+C|\nabla f|_{2}^{2}\right)^{\frac{1}{2}}
$$

with the Gaussian isoperimetric function $\mathcal{I} \equiv \gamma\left(\Phi^{-1}(x)\right)$, given by $\Phi^{\prime}(x)=\gamma(x) \equiv$ $\exp \left(-x^{2} / 2\right) / \sqrt{2 \pi}$ and some constant $C \in(0, \infty)$ for any function $0 \leq f \leq 1$ for which the right hand side is finite. One of the key properties of $\left(\mathrm{IFI}_{2}\right)$ is the product property which makes it suitable for the infinite dimensional analysis.

The result of [B2] was later extended in a number of directions including $[\mathbf{B a - L}],[\mathbf{B}-\mathbf{G} 2],[\mathbf{F o}],[\mathbf{Z} 1], \ldots$ and others, expanding greatly the set of measures for which it holds true.

To recover the isoperimetric information from $\left(\mathrm{IFI}_{2}\right)$, we note that the function $\mathcal{I}$ is (concave, symmetric around $1 / 2$ and) vanishing at 0 and 1 . Therefore one can use a suitable approximation of the characteristic function of a set $A$ to obtain the following relation

$$
\begin{equation*}
\mathcal{I}(\mu(A)) \leq C \mu_{2}^{+}(A) \tag{2}
\end{equation*}
$$

where the surface measure $\mu_{2}^{+}$corresponds formally to a suitable metric for which the natural length of the gradient of a Lipschitz function is $|\nabla f|_{2}$.

The inequality $\left(\mathrm{IFI}_{2}\right)$ is in general stronger than $\left(\mathrm{ISO}_{2}\right)$. However, following ([B2]) it is interesting to see by the following direct computation that using $\left(\mathrm{ISO}_{2}\right)$ with some extra information one can derive $\left(\mathrm{IFI}_{2}\right)$ in the Gaussian case. Namely if $\left(\mathrm{ISO}_{2}\right)$ is satisfied for the $\gamma^{\otimes 2}$, with $\gamma \equiv \gamma d x$ and $C=1$, then for $A=\{x \leq$ $\left.\Phi^{-1}(f(y))\right\}$, defined with a smooth function $0 \leq f \leq 1$, one gets

$$
\begin{equation*}
\mu(A)=\iint 1_{\left\{x \leq \Phi^{-1}(f(y))\right\}} \gamma(d x) \gamma(d y)=\int f(y) \gamma(d y) \tag{16.0.1}
\end{equation*}
$$

and choosing natural parametrisation $\left(x, \Phi^{-1}(f(x))\right)$ of the boundary $\partial A$ and using the definition of $\mu_{2}^{+}(\partial A)$ together with elementary calculus we obtain

$$
\begin{align*}
\mu_{2}^{+}(\partial A) & =\int \gamma\left(\Phi^{-1}(f)\right) \sqrt{1+\left|\nabla \Phi^{-1}(f)\right|^{2}} \gamma(d x) \\
16.0 .2) & =\int \mathcal{I}(f) \sqrt{1+\frac{1}{\mathcal{I}(f)^{2}}|\nabla f|^{2}(x)} \gamma(d x)=\int \sqrt{\mathcal{I}(f)^{2}+|\nabla f|^{2}} \gamma(d x) \tag{16.0.2}
\end{align*}
$$

Thus the relations (16.0.1) and (16.0.2) together with (ISO) imply ( $\mathrm{IFI}_{2}$ ).

One of the key features of $\left(\mathrm{IFI}_{2}\right)$ is the product property which makes it suitable for infinite dimensional analysis. As shown in $[\mathbf{Z 1}]$ the following natural generalisation of this inequality has also the product property

$$
\begin{equation*}
\mathcal{U}(\mu f) \leq \mu\left(\mathcal{U}(f)^{q}+C|\nabla f|_{q}^{q}\right)^{\frac{1}{q}} \tag{q}
\end{equation*}
$$

with a (concave) function $\mathcal{U}$ defined on the unit interval and vanishing at 0 and 1 , where $C \in(0, \infty)$ is independent of a function $0 \leq f \leq 1$ for which the right hand side make sense.

In this section we consider a function $\mathcal{U}(x) \equiv \varphi\left(F^{-1}(x)\right)$ where $F^{-1}$ is the inverse of the distribution function associated to the density $\varphi(z)=\alpha_{r}^{-1} e^{-|z|^{r}} d z \equiv$ $\varphi_{r}(z)$ with $\alpha_{r} \equiv \int_{\mathbb{R}} e^{-|z|^{r}} d z$ being a normalisation factor and $r \in(1, \infty)$. The choice of $\mathcal{U}$ is motivated by the fact that in dimension one this is the right isoperimetric function for the measure $\nu_{r}(d z) \equiv \varphi_{r}(z) d z,[\mathbf{B}-\mathbf{H} 2]$. Here, for $q \in(1,2]$ we choose $r \in[2, \infty)$ such that $\frac{1}{r}+\frac{1}{q}=1$ and show that in this case $\mathrm{IFI}_{q}$ implies $\mathrm{LS}_{q}$, extending the known result for $q=2$ (see e.g. [Ba-L]; we remark also that for that index $q=2$ the converse implication was shown recently by the semigroup technique in [Fo]).

We begin from the following lemma in which (as well as later on) we use a simplified notation $\nu(d z) \equiv \varphi(z) d z \equiv \nu_{r}(d z)$.

## Lemma 16.1.

(i) For $y<0$, we have

$$
\begin{equation*}
r^{-1}|y|^{-r+1}\left(1-q^{-1}|y|^{-r}\right) \varphi(y) \leq F(y) \leq r^{-1}|y|^{-r+1} \varphi(y) \tag{16.0.3}
\end{equation*}
$$

where $q^{-1}+r^{-1}=1$.
(ii) For $x \neq \frac{1}{2}$, we have

$$
\mathcal{U}^{\prime}(x)=-r\left|F^{-1}(x)\right|^{r-1} \operatorname{sign}\left(F^{-1}(x)\right)
$$

and

$$
\begin{equation*}
\mathcal{U}(x) \mathcal{U}^{\prime \prime}(x)=-\frac{r^{2}}{q}\left|F^{-1}(x)\right|^{r-2} \tag{16.0.4}
\end{equation*}
$$

(iii)

$$
\begin{equation*}
\lim _{x \longrightarrow 0^{+}} \frac{\mathcal{U}(x)}{x\left(\log \frac{1}{x}\right)^{1 / q}}=r \tag{16.0.5}
\end{equation*}
$$

Proof of (iii). We note that

$$
\begin{equation*}
\lim _{x \longrightarrow 0^{+}} \frac{\mathcal{U}(x)}{x\left(\log \frac{1}{x}\right)^{1 / q}}=\lim _{y \longrightarrow-\infty} \frac{\varphi(y)}{F(y)\left(\log \frac{1}{F(y)}\right)^{1 / q}} \tag{16.0.6}
\end{equation*}
$$

Using (16.0.3) and the fact that the function $x\left(\log \frac{1}{x}\right)^{1 / q}$ is increasing for small $x$, we have

$$
F(y)\left(\log \frac{1}{F(y)}\right)^{1 / q} \leq r^{-1}|y|^{-r+1} \varphi(y)\left(\log \frac{1}{r^{-1}|y|^{-r+1} \varphi(y)}\right)^{1 / q}
$$

$$
\begin{aligned}
& \leq r^{-1}|y|^{-r+1} \varphi(y)\left(|y|^{r}+\left|\log \frac{1}{r^{-1}|y|^{-r+1} \alpha_{r}^{-1}}\right|\right)^{1 / q} \\
& \leq r^{-1} \varphi(y)\left(|y|^{-r+1+r / q}+|y|^{-r+1}\left(\left|\log \frac{1}{r^{-1}|y|^{-r+1} \alpha_{r}^{-1}}\right|\right)^{1 / q}\right)
\end{aligned}
$$

Hence using the fact that $\frac{1}{r}+\frac{1}{q}=1$, we get

$$
F(y)\left(\log \frac{1}{F(y)}\right)^{1 / q} \leq r^{-1} \varphi(y)\left(1+|y|^{-r+1}\left(\left|\log \frac{1}{r^{-1}|y|^{-r+1} \alpha_{r}^{-1}}\right|\right)^{1 / q}\right)
$$

Thus

$$
\frac{\varphi(y)}{F(y)\left(\log \frac{1}{F(y)}\right)^{1 / q}} \geq r\left(1+|y|^{-r+1}\left(\left|\log \frac{1}{r^{-1}|y|^{-r+1} \alpha_{r}^{-1}}\right|\right)^{1 / q}\right)^{-1} \longrightarrow y \longrightarrow \infty r
$$

Similar arguments using the lower bound (16.0.3) for the distribution $F$ allows us to get the upper bound converging to $r$ as $y$ goes to $-\infty$. This ends the proof of (iii) .

Recall that given a probability measure $\mu$ and a nonnegative measurable function $f$ we define the relative entropy as follows

$$
\operatorname{Ent}_{\mu}(f) \equiv \mu\left(f \log \frac{f}{\mu f}\right)
$$

Proposition 16.2. Suppose $r \in[2, \infty)$ and $\frac{1}{r}+\frac{1}{q}=1$. Then there is a constant $A \in[r, \infty)$ such that the function

$$
\xi(x) \equiv \mathcal{U}^{q}(x) / x^{q-1}-A x \log \frac{1}{x}
$$

is convex on $\left[0, \frac{1}{2}\right]$ and therefore for any probability measure $\mu$ and any measurable function $0 \leq f \leq \frac{1}{2}$ we have

$$
\begin{equation*}
\mathcal{U}^{q}(\mu f) /(\mu f)^{q-1}-\mu\left(\mathcal{U}^{q}(f) / f^{q-1}\right) \leq A \operatorname{Ent}_{\mu}(f) \tag{16.0.7}
\end{equation*}
$$

Moreover, for any function $0 \leq f \leq 1$, we have

$$
\begin{equation*}
\lim _{\varepsilon \longrightarrow 0^{+}} \frac{1}{\varepsilon}\left(\mathcal{U}^{q}(\varepsilon \mu f) /(\varepsilon \mu f)^{q-1}-\mu\left(\mathcal{U}^{q}(\varepsilon f) /(\varepsilon f)^{q-1}\right)\right)=r \operatorname{Ent}_{\mu}(f) \tag{16.0.8}
\end{equation*}
$$

Proof. The inequality (16.0.7) follows by Jensen inequality from the convexity of the function $\xi$. To prove the convexity of $\xi$ we note that

$$
\xi^{\prime}(x) \equiv q \mathcal{U}^{q-1}(x) \mathcal{U}^{\prime}(x) / x^{q-1}-(q-1) \mathcal{U}^{q}(x) / x^{q}-A \log \frac{1}{x}+A
$$

and hence

$$
\begin{aligned}
\xi^{\prime \prime}(x) & \equiv q(q-1) \mathcal{U}^{q-2}(x)\left(\mathcal{U}^{\prime}(x)\right)^{2} / x^{q-1}-2 q(q-1) \mathcal{U}^{q-1}(x) \mathcal{U}^{\prime}(x) / x^{q} \\
& -q(q-1) \mathcal{U}^{q}(x) / x^{q+1}+q \mathcal{U}^{q-1}(x) \mathcal{U}^{\prime \prime}(x) / x^{q-1}+\frac{A}{x}
\end{aligned}
$$

and therefore

$$
\begin{align*}
\xi^{\prime \prime}(x) & \equiv q(q-1) \mathcal{U}^{q-2} / x^{q-1}\left(\mathcal{U}^{\prime}(x)-\mathcal{U}(x) / x\right)^{2}  \tag{16.0.9}\\
& +\left[q \mathcal{U}^{q-1}(x) \mathcal{U}^{\prime \prime}(x) / x^{q-1}+\frac{A}{x}\right]
\end{align*}
$$

Thus we need to show that the second term on the right hand side of (16.0.9) is nonnegative. To this end we note that, choosing $x=F(y)$ and using (16.0.4) we have

$$
\begin{gather*}
q \mathcal{U}^{q-1}(x) \mathcal{U}^{\prime \prime}(x) / x^{q-1}+\frac{A}{x}=q \mathcal{U}^{q-2}(x)\left(-r(r-1)\left|F^{-1}(x)\right|^{r-2}\right) / x^{q-1}+\frac{A}{x}=  \tag{16.0.10}\\
q \varphi(y)^{q-2}\left(-\frac{r^{2}}{q}|y|^{r-2}\right) /\left(F(y)^{q-1}\right)+\frac{A}{F(y)}=\frac{1}{F(y)}\left[A-r^{2}|y|^{r-2} \varphi(y)^{q-2} /\left(F(y)^{q-2}\right)\right]
\end{gather*}
$$

Next, for $y \leq-1$ we use the lower bound from (16.0.3) to see that

$$
\begin{aligned}
r^{2}|y|^{r-2} /\left(\left(r^{-1}|y|^{-r+1}\left(1-q^{-1}|y|^{-r}\right)\right)^{q-2}\right) & \left.=r^{q}|y|^{r-2+(r-1)(q-2)} /\left(1-q^{-1}|y|^{-r}\right)^{q-2}\right) \\
& \left.=r^{q} /\left(1-q^{-1}|y|^{-r}\right)^{q-2}\right) \leq r^{2 q r}
\end{aligned}
$$

where in the last line we have taken into the account that $(r-1)(q-1)=1$ (because $r$ and $q$ are dual to each other). On the other hand if $r \geq 2$, one has that

$$
\sup _{y \in[-1,0]}|y|^{r-2} \varphi(y)^{q-2} /\left(F(y)^{q-2}\right)<\infty
$$

Hence choosing $A \equiv \max \left(r^{2 q r}, \sup _{y \in[-1,0]}|y|^{r-2} \varphi(y)^{q-2} /\left(F(y)^{q-2}\right)\right)$ we get $\xi(x) \geq$ 0 for all $x \in\left[1, \frac{1}{2}\right]$.

The relation (16.0.8) follows from (16.0.7) and (16.0.5) .
Assuming we are given a natural gradient $\nabla f \equiv\left(\nabla_{j} f\right)_{j \in \mathcal{R}}$, with coordinates $\nabla_{j} f$ indexed by a countable set $\mathcal{R}$, we define its $l_{q}$ norm $|\nabla f|_{q}$ as follows

$$
|\nabla f|_{q}^{q} \equiv \sum_{j}\left|\nabla_{j} f\right|^{q}
$$

Theorem 16.3. Suppose $r \in[2, \infty)$ and $\frac{1}{r}+\frac{1}{q}=1$ and $\operatorname{let} \mathcal{U}(x) \equiv \varphi\left(F^{-1}(x)\right)$ with $\varphi=\varphi_{r}=F^{\prime}$. If there is a constant $C \in(0, \infty)$ such that

$$
\begin{equation*}
\mathcal{U}(\mu f) \leq \mu\left(\mathcal{U}(f)^{q}+C|\nabla f|_{q}^{q}\right)^{\frac{1}{q}} \tag{q}
\end{equation*}
$$

for any function $0 \leq f \leq \frac{1}{2}$ for which $|\nabla f|_{q}$ is integrable, then there is a constant $C^{\prime} \in(0, \infty)$ such that

$$
\begin{equation*}
\mu f^{q} \log \frac{f^{q}}{\mu f^{q}} \leq C^{\prime} \mu|\nabla f|_{q}^{q} \tag{q}
\end{equation*}
$$

for all nonnegative functions $f$ for which the right hand side is finite.
Proof. Given a function $f$ for which $|\nabla f|_{q}^{q}$ is $\mu$ integrable, we define a function $G(\lambda, z)$ of nonnengative arguments $\lambda, z \in[0, \infty)$, by

$$
G(\lambda, z) \equiv \mu\left[\lambda^{q} f^{q}+z^{q}\left(\mathcal{U}(f)^{q}+C|\nabla f|_{q}^{q}\right)\right]^{\frac{1}{q}}-\left[\lambda^{q}(\mu f)^{q}+z^{q} \mathcal{U}(\mu f)^{q}\right]^{\frac{1}{q}}
$$

with some $C \in(0, \infty)$. For later purposes we note that $G(\lambda, z)$ is continuous in its domain and uniformly continuous in an open neighbourhood of the origin. Its derivative with respect to $z^{q}$ for $\lambda>0$ is given by

$$
\frac{\partial}{\partial z^{q}} G(\lambda, z)=\mu \frac{\left(\mathcal{U}(f)^{q}+C|\nabla f|_{q}^{q}\right)}{\left[\lambda^{q} f^{q}+z^{q}\left(\mathcal{U}(f)^{q}+C|\nabla f|_{q}^{q}\right)\right]^{1-\frac{1}{q}}}-\frac{\mathcal{U}(\mu f)^{q}}{\left[\lambda^{q}(\mu f)^{q}+z^{q} \mathcal{U}(\mu f)^{q}\right]^{1-\frac{1}{q}}}
$$

In particular at $z=0$ we have

$$
\begin{equation*}
\lambda^{\frac{q}{r}} \frac{\partial}{\partial z^{q}} G(\lambda, z=0)=\left(\mu\left(\frac{\mathcal{U}(f)^{q}}{f^{q-1}}\right)-\frac{\mathcal{U}(\mu f)^{q}}{(\mu f)^{q-1}}\right)+q^{q} C \mu\left|\nabla f^{\frac{1}{q}}\right|_{q}^{q} \tag{16.0.11}
\end{equation*}
$$

To show implication $\mathrm{IFI}_{q} \Longrightarrow \mathrm{LS}_{q}$, we note that by Minkowski inequality one gets (16.0.12)

$$
G(\lambda, z) \geq\left[\lambda^{q}(\mu f)^{q}+z^{q}\left(\mu\left(\mathcal{U}(f)^{q}+C|\nabla f|_{q}^{q}\right)^{\frac{1}{q}}\right)^{q}\right]^{\frac{1}{q}}-\left[\lambda^{q}(\mu f)^{q}+z^{q} \mathcal{U}(\mu f)^{q}\right]^{\frac{1}{q}}
$$

Thus if $\mathrm{IFI}_{q}$ is satisfied, the function $G(\lambda, z)$ is nonnegative and equal to zero for $z=0$. This implies that its derivative with respect to $z^{q}$ at $z=0$ and $\lambda>0$, is also nonnegative and using the formula (16.0.11) we have

$$
\begin{equation*}
\frac{\mathcal{U}(\mu f)^{q}}{(\mu f)^{q-1}}-\mu\left(\frac{\mathcal{U}(f)^{q}}{f^{q-1}}\right) \leq q^{q} C \mu\left|\nabla f^{\frac{1}{q}}\right|_{q}^{q} \tag{16.0.13}
\end{equation*}
$$

We replace $f$ by $\varepsilon f$ and divide both sides by $\varepsilon$. Passing with $\varepsilon$ to 0 and using (16.0.8) we conclude that the following inequality is true

$$
\begin{equation*}
\operatorname{Ent}_{\mu}(f) \leq r^{-1} q^{q} C \mu\left|\nabla f^{\frac{1}{q}}\right|_{q}^{q} \tag{16.0.14}
\end{equation*}
$$

For nonnegative functions $f$ this is equivalent to $\mathrm{LS}_{q}$. This ends the proof.

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