

HOMEWORK ASSIGNMENT N4, MATH 4567, SPRING 2018  
Due on April 16 (Monday) – problems and solutions

**Problem 1.** a) Given parameters  $c > 0$  and  $\beta > 0$ , show that the Sturm-Liouville boundary value problem

$$\begin{aligned}y'' + \lambda y &= 0, & y &= y(x), \quad 0 \leq x \leq c, \\y'(0) &= \beta y(0), \\y'(c) &= \beta y(c),\end{aligned}$$

has exactly one negative eigenvalue  $\lambda_0$  and that this eigenvalue is independent on  $c > 0$ . Find  $\lambda_0$  and an associated eigenfunction  $y_0(x)$ .

b) Determine whether or not  $\lambda = 0$  is an eigenvalue. If yes, find an associated eigenfunction.

**Solution.** We know that any Sturm-Liouville problem

$$\begin{aligned}y''(x) + \lambda y(x) &= 0, & 0 \leq x \leq c, \\ \alpha_1 y(0) + \alpha_2 y'(0) &= 0, \\ \beta_1 y(c) + \beta_2 y'(c) &= 0\end{aligned}$$

has only non-negative eigenvalues  $\lambda$ , when  $\alpha_1 \alpha_2 \leq 0$  and  $\beta_1 \beta_2 \geq 0$ . In our case,  $\alpha_1 = 1$ ,  $\alpha_2 = -\beta$ ,  $\beta_1 = 1$ ,  $\beta_2 = -\beta$ , so this condition is not fulfilled, and one should consider separately the three cases to find all eigenvalues. Actually, in this problem, we are asked about non-positive  $\lambda$ .

a) Case 1:  $\lambda = -\alpha^2$  ( $\alpha > 0$ ). The general solution to the Sturm-Liouville equation  $y''(x) + \lambda y(x) = 0$  may be described in this case by the formula

$$y = C_1 \cosh(\alpha x) + C_2 \sinh(\alpha x)$$

with arbitrary coefficients  $C_1$  and  $C_2$ , where we use the hyperbolic functions  $\cosh(t) = \frac{e^t + e^{-t}}{2}$  and  $\sinh(t) = \frac{e^t - e^{-t}}{2}$ . Hence

$$y(0) = C_1, \quad y' = \alpha C_1 \sinh(\alpha x) + \alpha C_2 \cosh(\alpha x), \quad y'(0) = \alpha C_2,$$

and the first boundary condition becomes

$$\alpha C_2 = \beta C_1.$$

This gives a necessary relationship between the coefficients. Since we are looking for a non-trivial (not identically zero) solution  $y = y(x)$  up to a multiplicative factor, one may take any fixed value of  $C_1 \neq 0$ , and then the unique choice for the other coefficient will be  $C_2 = \beta C_1 / \alpha$ . To work with simpler expressions, let us choose  $C_1 = \alpha$ , so that  $C_2 = \beta$ . The solution  $y(x)$  and its derivative  $y'(x)$  are thus simplified to

$$y = \alpha \cosh(\alpha x) + \beta \sinh(\alpha x), \quad y' = \alpha^2 \sinh(\alpha x) + \beta \alpha \cosh(\alpha x).$$

In particular,

$$y(c) = \alpha \cosh(\alpha c) + \beta \sinh(\alpha c), \quad y'(c) = \alpha^2 \sinh(\alpha c) + \beta \alpha \cosh(\alpha c),$$

and the second boundary condition becomes

$$\alpha^2 \sinh(\alpha c) + \beta \alpha \cosh(\alpha c) = \beta \alpha \cosh(\alpha c) + \beta^2 \sinh(\alpha c).$$

Since  $\sinh(\alpha c) > 0$ , this is equivalent to  $\alpha^2 = \beta^2$ , which is uniquely solved as  $\alpha = \beta$  (in positive numbers). Therefore,

$$\lambda = -\alpha^2 = -\beta^2, \quad y = \beta \cosh(\beta x) + \beta \sinh(\beta x) = \beta e^{\beta x}.$$

Since we are looking for a non-trivial solution  $y = y(x)$  up to a multiplicative factor, the constant  $\beta$  may be removed.

b) Case 2:  $\lambda = 0$ . Then the general solution is  $y = C_1 + C_2 x$ , and the boundary conditions become

$$\begin{aligned} C_2 &= \beta C_1, \\ C_2 &= \beta (C_1 + C_2 c). \end{aligned}$$

This linear system is solved as  $C_1 = C_2 = 0$ , which means that this case is impossible.

**Answer.**

- a)  $\lambda_0 = -\beta^2$  is the unique negative eigenvalue with eigenfunction  $y_0 = e^{\beta x}$ .
- b)  $\lambda = 0$  is not an eigenvalue of the given Sturm-Liouville problem.

**Problem 2.** Solve the temperature problem:

$$\begin{aligned} u_t &= k u_{xx}, & u &= u(x, t), \quad 0 \leq x \leq \pi, \quad t > 0 \quad (k > 0) \\ u_x(0, t) &= \beta u(0, t), \\ u_x(\pi, t) &= \beta u(\pi, t), \\ u(x, 0) &= f(x), \end{aligned}$$

where  $f(x)$  is a given continuous function on  $[0, \pi]$  and  $\beta$  is a positive parameter. Write your answer in the form of an infinite series

$$u(x, t) = \sum_{n=0}^{\infty} c_n y_n(x) T_n(t). \tag{1}$$

Describe the functions  $y_n(x)$  and  $T_n(t)$  that are involved and indicate how to compute the coefficients  $c_n$  in terms of  $f$ .

**Solution.** As we know, the solution to the given temperature problem is described by the functional series (1) in which  $y_n$  represent the eigenfunctions of the associated Sturm-Liouville boundary value problem, namely

$$\begin{aligned} y'' + \lambda y &= 0, & y &= y(x), \quad 0 \leq x \leq \pi, \\ y'(0) &= \beta y(0), \\ y'(\pi) &= \beta y(\pi), \end{aligned}$$

and  $T_n(t) = e^{-\lambda_n kt}$ , where  $\lambda_n$  is the eigenvalue to which the eigenfunction  $y_n$  belongs. Moreover, the coefficients  $c_n$  in (1) are determined by the last non-homogeneous boundary condition via the orthogonal series expansion

$$f(x) = \sum_{n=0}^{\infty} c_n y_n(x).$$

Here, necessarily

$$c_n = \frac{\langle f, y_n \rangle}{\|y_n\|^2}, \quad \langle f, y_n \rangle = \int_0^{\pi} f(x) y_n(x) dx, \quad \|y_n\|^2 = \int_0^{\pi} y_n(x)^2 dx.$$

As one can notice, the associated Sturm-Liouville boundary value problem is exactly the same as the one in Problem 1 with  $c = \pi$ . So, we already know that this problem has a unique negative eigenvalue  $\lambda_0 = -\beta^2$  with its eigenfunction  $y_0(x) = e^{\beta x}$ . Hence,  $T_0(t) = e^{\beta^2 kt}$ . We also know that all remaining eigenvalues are positive. To find these eigenvalues and the eigenfunctions, we need to consider:

*Case 3:*  $\lambda = \alpha^2$  ( $\alpha > 0$ ). In this case, the general solution to the Sturm-Liouville equation  $y''(x) + \lambda y(x) = 0$  is given by

$$y = C_1 \cos(\alpha x) + C_2 \sin(\alpha x), \quad y(0) = C_1,$$

for which we have

$$y' = \alpha C_2 \cos(\alpha x) - \alpha C_1 \sin(\alpha x), \quad y'(0) = \alpha C_2.$$

Hence  $\alpha C_2 = \beta C_1$ , by the first boundary condition. Without loss of generality, put  $C_1 = \alpha$ ,  $C_2 = \beta$ , so that

$$y = \alpha \cos(\alpha x) + \beta \sin(\alpha x), \quad y' = \beta \alpha \cos(\alpha x) - \alpha^2 \sin(\alpha x).$$

Involving the second boundary condition, we obtain that  $\alpha$  should solve the equation

$$\beta \alpha \cos(\alpha \pi) - \alpha^2 \sin(\alpha \pi) = \beta \alpha \cos(\alpha \pi) + \beta^2 \sin(\alpha \pi)$$

which is equivalent to  $\sin(\alpha \pi) = 0$ . It is solved as  $\alpha = \alpha_n = n$ ,  $n \geq 1$ , in which case

$$y = y_n(x) = n \cos(nx) + \beta \sin(nx).$$

Let us also specify the  $L^2$ -norms of the eigenfunctions: For  $n = 0$  we have

$$\|y_0\|^2 = \int_0^{\pi} y_0(x)^2 dx = \int_0^{\pi} e^{2\beta x} dx = \frac{e^{2\beta\pi} - 1}{2\beta},$$

and for integers  $n \geq 1$ ,

$$\begin{aligned}
 \|y_n\|^2 &= \int_0^\pi (n \cos(nx) + \beta \sin(nx))^2 dx \\
 &= n^2 \int_0^\pi \cos^2(nx) dx + 2\beta n \int_0^\pi \cos(nx) \sin(nx) dx + \beta^2 \int_0^\pi \sin^2(nx) dx \\
 &= n^2 \int_0^\pi \frac{1 + \cos(2nx)}{2} dx + \beta n \int_0^\pi \sin(2nx) dx + \beta^2 \int_0^\pi \frac{1 - \cos(2nx)}{2} dx \\
 &= \frac{\pi}{2} (n^2 + \beta^2).
 \end{aligned}$$

Here we used the obvious identities

$$\int_0^\pi \cos(2nx) dx = \int_0^\pi \sin(2nx) dx = 0.$$

**Answer.** The solution to the temperature problem is given by the functional series (1), where

$$\begin{aligned}
 y_0(x) &= e^{\beta x}, & y_n(x) &= n \cos(nx) + \beta \sin(nx) \quad (n \geq 1), \\
 T_0(t) &= e^{\beta^2 kt}, & T_n(t) &= e^{-nkt} \quad (n \geq 1),
 \end{aligned}$$

with coefficients

$$\begin{aligned}
 c_0 &= \frac{2\beta}{e^{2\beta\pi} - 1} \int_0^\pi f(x) e^{\beta x} dx, \\
 c_n &= \frac{2}{\pi(n^2 + \beta^2)} \int_0^\pi f(x) (n \cos(nx) + \beta \sin(nx)) dx \quad (n \geq 1).
 \end{aligned}$$

**Problem 3.** Assume that the initial temperatures are constant: Let  $f(x) = 1$ . Determine the first three terms in the representation (1), that is, eventually, evaluate  $c_0$ ,  $c_1$ , and  $c_2$ .

**Solution.** Using the answer to Problem 2, we find that

$$\begin{aligned}
 c_0 &= \frac{2\beta}{e^{2\beta\pi} - 1} \int_0^\pi e^{\beta x} dx = 2 \frac{e^{\beta\pi} - 1}{e^{2\beta\pi} - 1}, \\
 c_1 &= \frac{2}{\pi(1 + \beta^2)} \int_0^\pi (\cos x + \beta \sin x) dx = \frac{4\beta}{\pi(1 + \beta^2)}, \\
 c_2 &= \frac{2}{\pi(4 + \beta^2)} \int_0^\pi (2 \cos(2x) + \beta \sin(2x)) dx = 0.
 \end{aligned}$$

**Answer.** The solution to the temperature problem with the first three terms is given by

$$u(x, t) = 2 \frac{e^{\beta\pi} - 1}{e^{2\beta\pi} - 1} e^{\beta x} e^{\beta^2 kt} + \frac{4\beta}{\pi(1 + \beta^2)} (\cos x + \beta \sin x) e^{-kt} + 0 + \dots$$

**Problem 4.** Given parameters  $A, B, C$  and  $\beta > 0$ , consider the temperature problem with non-homogeneous boundary conditions:

$$\begin{aligned} u_t &= k u_{xx}, & u &= u(x, t), \quad 0 \leq x \leq \pi, \quad t > 0 \quad (k > 0) \\ u_x(0, t) &= \beta u(0, t) + A, \\ u_x(\pi, t) &= \beta u(\pi, t) + B, \\ u(x, 0) &= Cx. \end{aligned}$$

Reduce it to Problem 2 by virtue of a suitable substitution  $u(x, t) = U(x, t) + \Phi(x)$ . Indicate new initial temperatures  $F(x)$  in the homogeneous problem about  $U(x, t)$ .

**Solution.** One may look for a suitable substitution  $u(x, t) = U(x, t) + \Phi(x)$  with a linear function  $\Phi(x) = C_1 + C_2 x$ , in which case the problem becomes

$$\begin{aligned} U_t &= k U_{xx}, & U &= U(x, t), \quad 0 \leq x \leq \pi, \quad t > 0 \quad (k > 0) \\ U_x(0, t) + C_2 &= \beta (U(0, t) + C_1) + A, \\ U_x(\pi, t) + C_2 &= \beta (U(\pi, t) + C_1 + C_2 \pi) + B, \\ U(x, 0) + (C_1 + C_2 x) &= Cx. \end{aligned}$$

In order to reach homogeneous vertical boundary conditions, the unknown parameters should satisfy

$$\begin{aligned} C_2 &= \beta C_1 + A, \\ C_2 &= \beta (C_1 + C_2 \pi) + B. \end{aligned}$$

This is a linear system in two unknowns which is easily solved. Equalizing the two equalities, we get  $A = C_2 \beta \pi + B$ , so

$$C_2 = \frac{A - B}{\beta \pi}.$$

Inserting this into the second equality, we also find

$$C_1 = \frac{1 - \beta \pi}{\beta^2 \pi} A - \frac{1}{\beta^2 \pi} B.$$

**Answer.** The new problem with homogeneous vertical boundary conditions is

$$\begin{aligned} U_t &= k U_{xx}, & U &= U(x, t), \quad 0 \leq x \leq \pi, \quad t > 0 \quad (k > 0) \\ U_x(0, t) &= \beta U(0, t), \\ U_x(\pi, t) &= \beta U(\pi, t), \\ U(x, 0) &= F(x) \end{aligned}$$

with new initial temperatures

$$F(x) = Cx - \left( \frac{1 - \beta \pi}{\beta^2 \pi} A - \frac{1}{\beta^2 \pi} B \right) - \frac{A - B}{\beta \pi} x.$$